
New materials for holographic hydrodynamics

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Motivation and Overview

- Collaboration between string theorists and condensed matter theorists
- Proposing new materials to test predictions from gauge/gravity duality
- Outline:
 - Hydrodynamics for electrons in solids
 - Brief intro to gauge/gravity duality
 - Predictions of gauge/gravity duality for electron hydrodynamics

- New paper:

- Turbulent hydrodynamics in strongly correlated Kagome metals

- Domenico Di Sante, J. E., Martin Greiter, Ioannis Matthaiakakis, René Meyer, David Rodriguez Fernandez, Ronny Thomale, Erik van Loon, Tim Wehling

- arXiv:cond-mat/1911.06810

- Proposal for a new Dirac material with stronger electronic coupling than in graphene: Scandium-Herbertsmithite

- in view of enhanced hydrodynamic behaviour of the electrons

- Reaching smaller η/s (ratio of shear viscosity over entropy density)

Motivation and Overview

- Strongly coupled electron fluids in the Poiseuille regime

J.E., I. Matthaiakakis, R. Meyer, D. Rodriguez, Phys. Rev. B 98 (2018) 195143

- Functional dependence of the Hall viscosity-induced transverse voltage in two-dimensional Fermi liquids

J.E., E. Hankiewicz, I. Matthaiakakis, R. Meyer, D. Rodriguez, C. Tutschku, Phys. Rev. B 101 (2020) 045423

Hydrodynamics for electrons in solids

When phonon and impurity interactions are suppressed,

Electron-electron interactions may lead to a hydrodynamic electron flow

(Small parameter window)

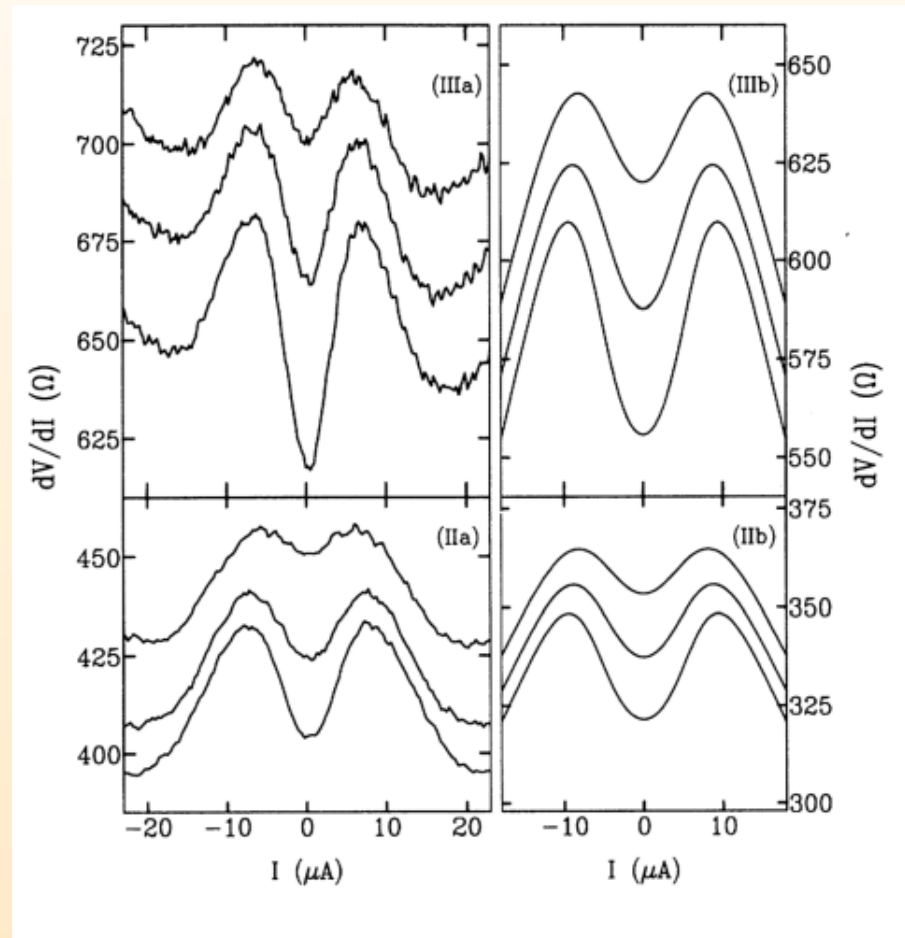
Some Implications:

- Decrease of differential resistance dV/dI with increasing current I

Weak coupling: High mobility wires

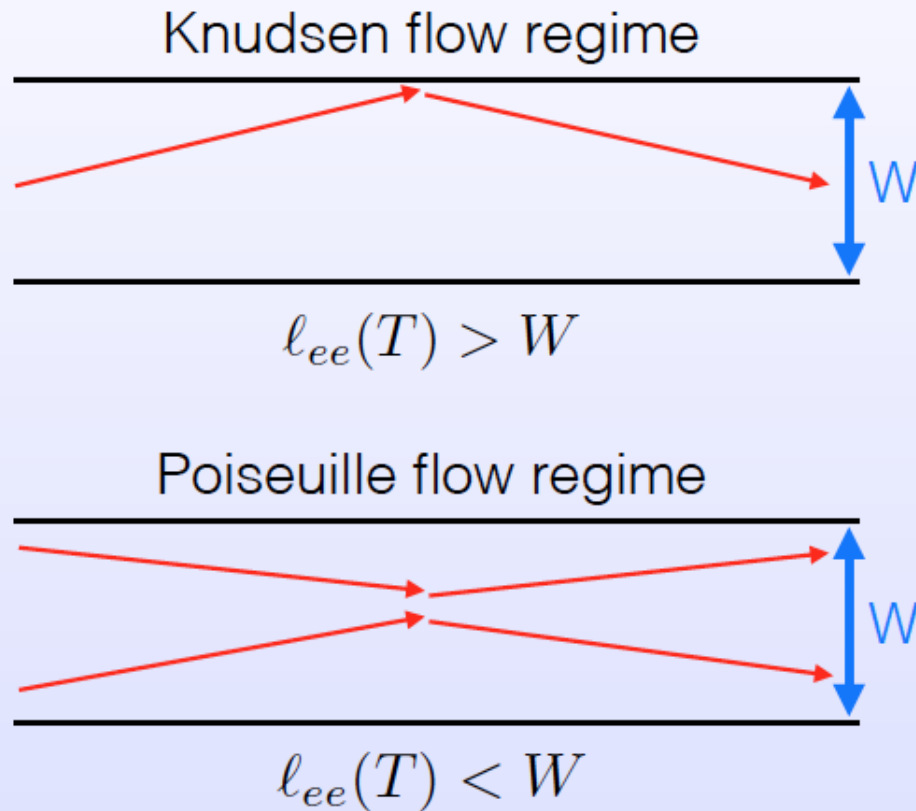
Transition: Knudsen flow \Rightarrow Poiseuille flow Gurzhi effect

Molenkamp, de Jong Phys. Rev. B 51 (1995) 13389 for GaAs in 2+1 dimensions

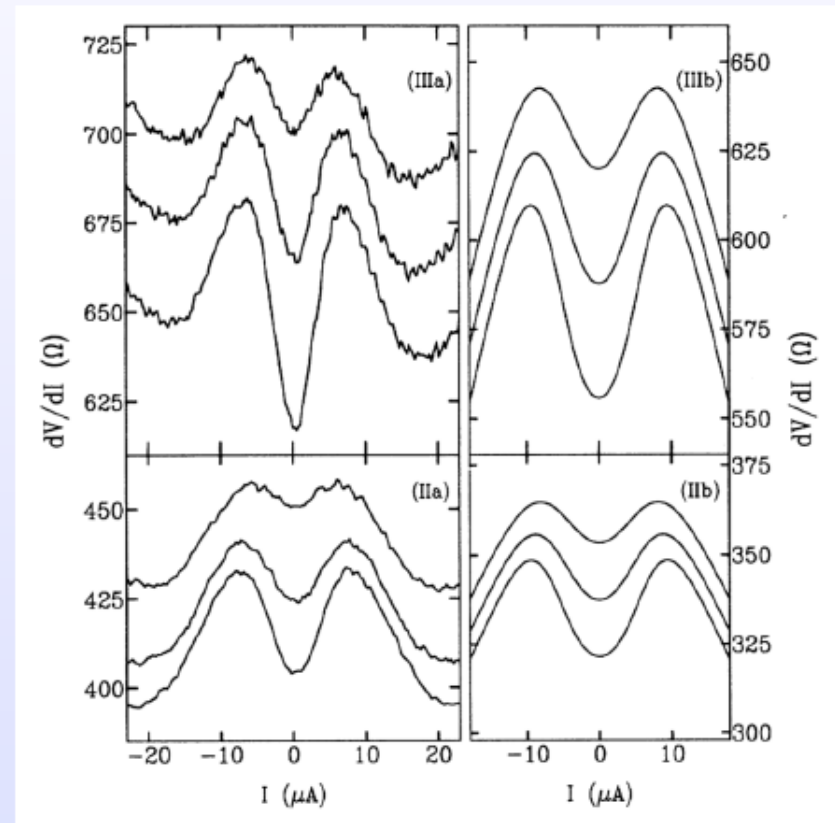


Transition from ballistic to hydrodynamic regime

- 2D Electrons in (Al)GaAs Heterostructures



[Gurzhi 1968]



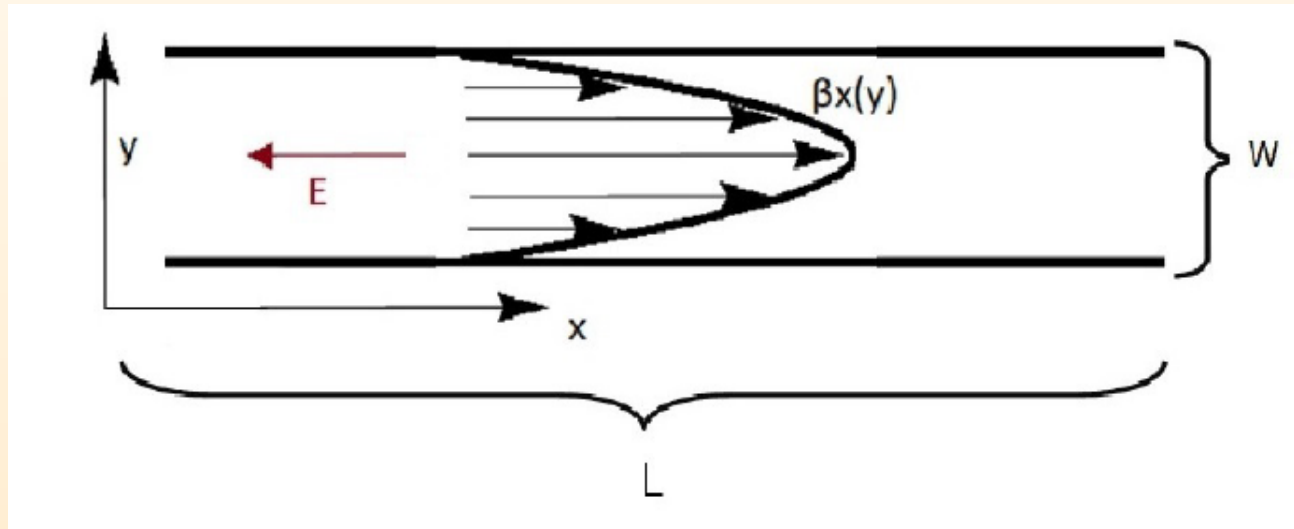
[Molenkamp+de Jong 1994,95]

Conditions for hydrodynamic behaviour of electrons in solids

$$l_{ee} < l_{\text{imp}}, l_{\text{phonon}}, W$$

l_{ee} : Typical scale for electron-electron scattering

Flow profile in wire



Effective electron-electron coupling strength

$$\alpha_{\text{eff}} = \frac{e^2}{\epsilon_0 \epsilon_r \hbar v_F}$$

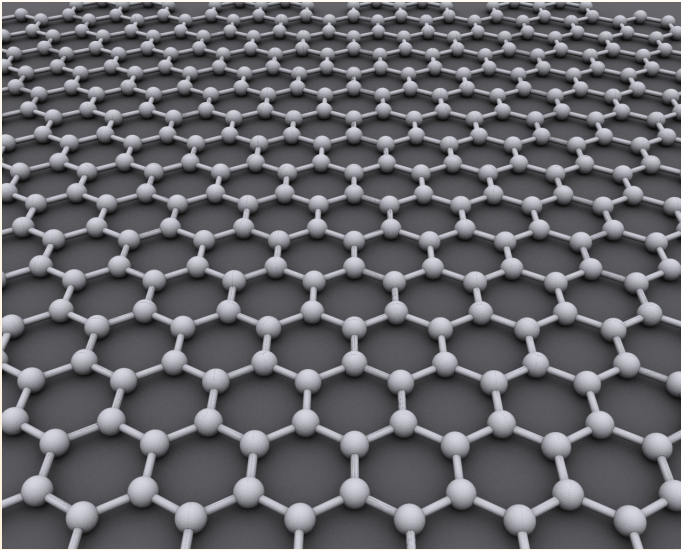
Electron-electron scattering length:

$$l_{\text{ee}} \propto \frac{1}{\alpha_{\text{eff}}^2}$$

Larger electronic coupling \Rightarrow More robust hydrodynamic behaviour

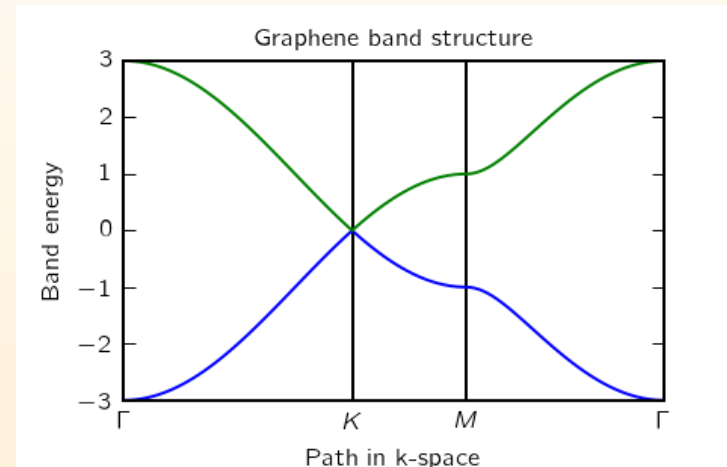
Hydrodynamics in Dirac materials: Graphene

Hexagonal carbon lattice



Source: Wikipedia

Dirac material: Linear dispersion relation



Considerable theoretical and experimental effort for viscous fluids

Review: Polini + Geim, [arXiv:1909.10615](https://arxiv.org/abs/1909.10615)

Relativistic hydrodynamics

Relativistic hydrodynamics: Expansion in four-velocity derivatives

$$T_{\mu\nu} = (\epsilon + P)u^\mu u^\nu + P\eta^{\mu\nu} - \sigma^{\mu\nu} + \dots$$

$$\sigma^{\mu\nu} = P^{\mu\alpha}P^{\nu\beta} \left(\eta(\nabla_\alpha u_\beta + \nabla_\beta u_\alpha - \frac{2}{3}\nabla_\gamma u^\gamma \eta_{\alpha\beta}) + \zeta \nabla_\gamma u^\gamma \eta_{\alpha\beta} \right)$$

Shear viscosity η , bulk viscosity ζ

$$P^{\mu\nu} = \eta^{\mu\nu} + u^\mu u^\nu$$

Relativistic hydrodynamics: Expansion in four-velocity derivatives

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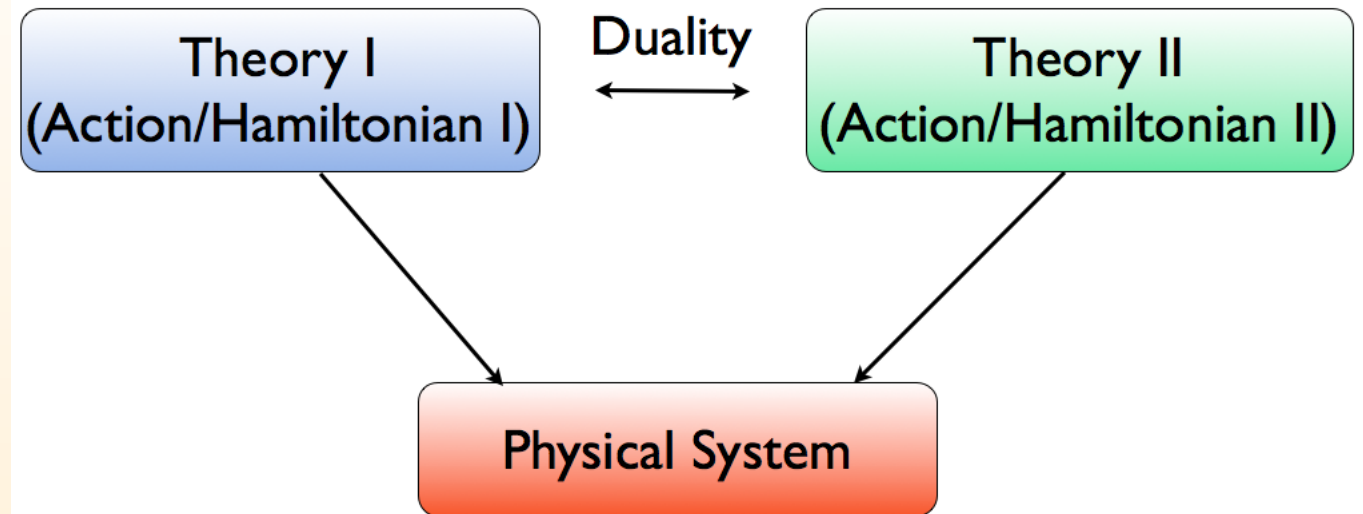
Shear viscosity η , bulk viscosity ζ

$$P^{\mu\nu} = \eta^{\mu\nu} + u^\mu u^\nu$$

Shear viscosity for strongly correlated systems may be calculated from gauge/gravity duality!

Short Intro to Gauge/Gravity Duality

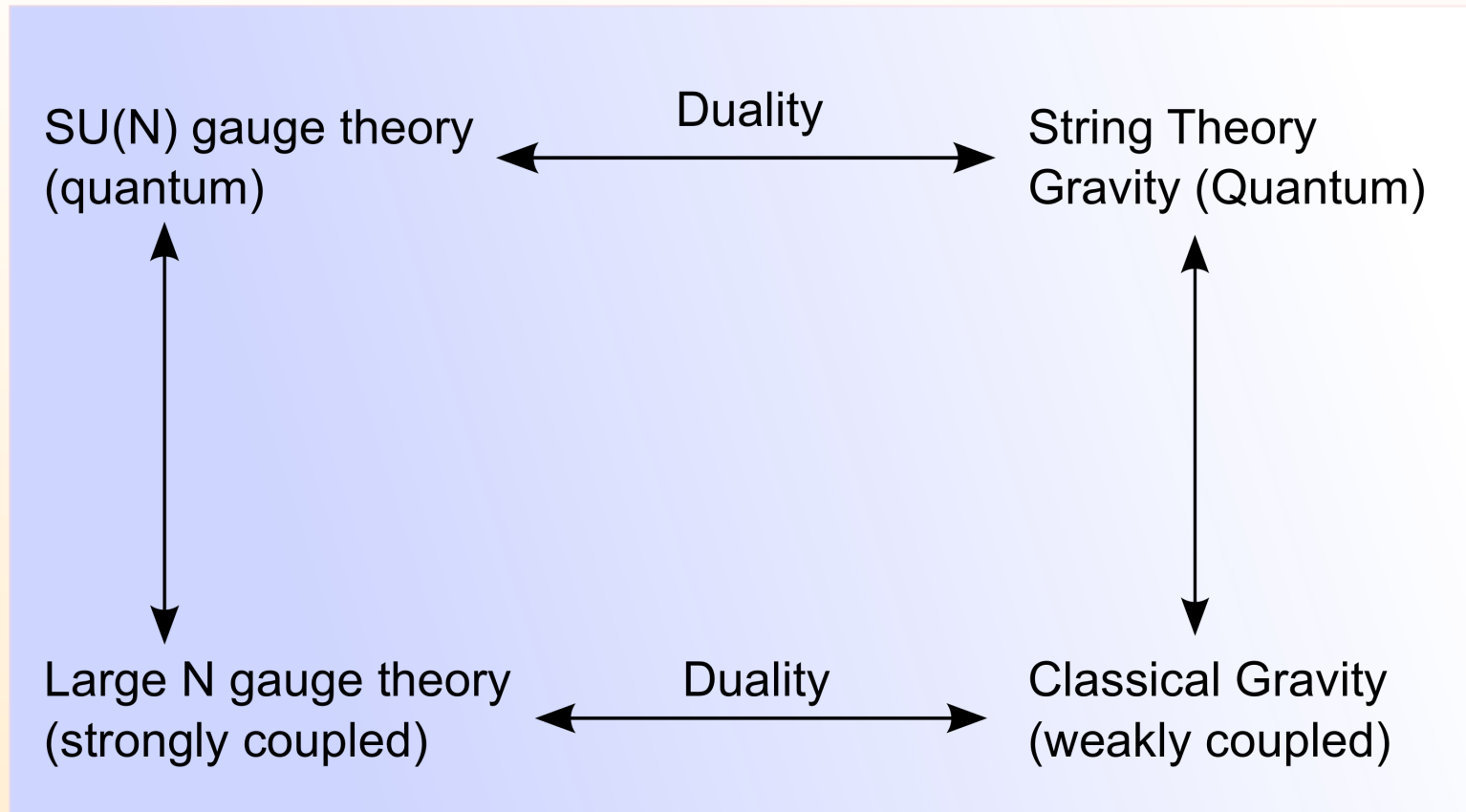
Duality:

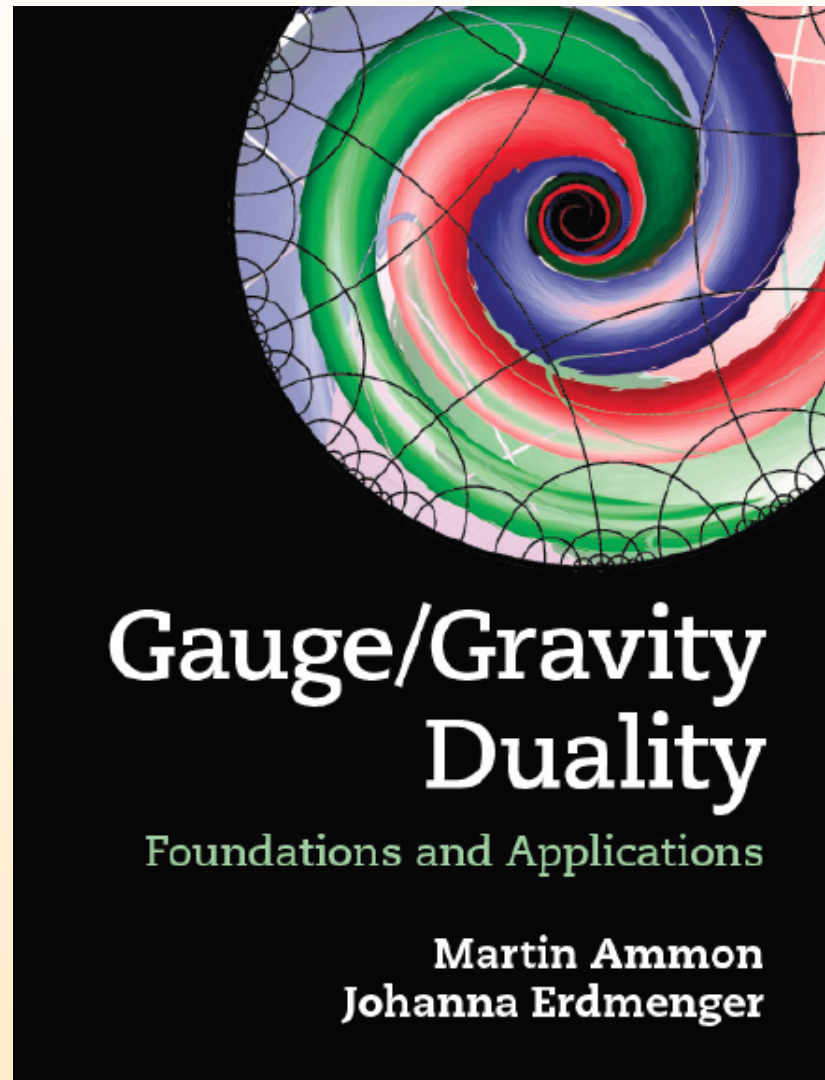


Gauge/Gravity Duality

- Conjecture which follows from a low-energy limit of string theory
- **Duality:**
Quantum field theory at **strong** coupling
 \Leftrightarrow **Theory of gravitation** at **weak** coupling
- **Holography:**
Quantum field theory in d dimensions
 \Leftrightarrow **Gravitational theory** in $d + 1$ dimensions
Quantum field theory defined on the boundary of the $d + 1$ -dimensional space

Gauge/Gravity Duality: String Theory Origin





AdS/CFT correspondence

AdS/CFT correspondence:

Example of gauge/gravity duality with huge amount of symmetry

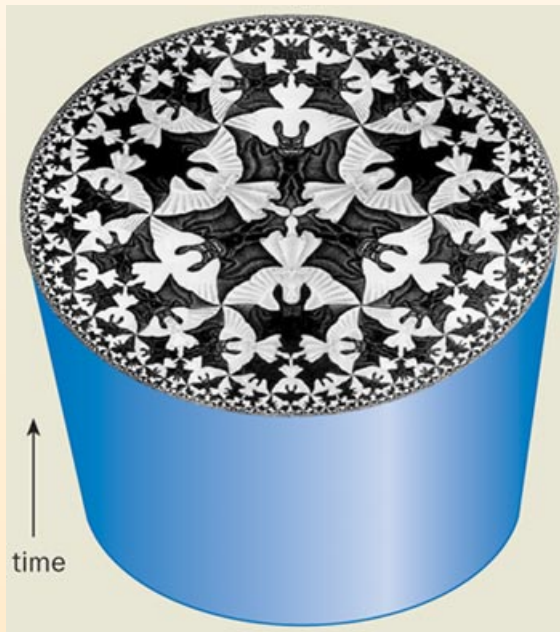
AdS: Anti-de Sitter space:

Hyperbolic space

with constant negative curvature

CFT: Conformal field theory

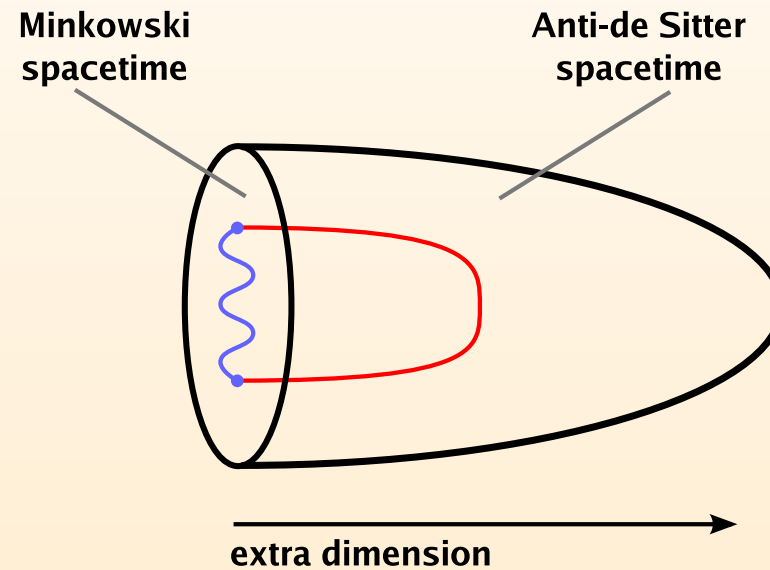
Example: QFT at RG fixed point



Quelle: Institute of Physics, Copyright: C. Escher

Gauge/Gravity Duality: Bulk-boundary correspondence

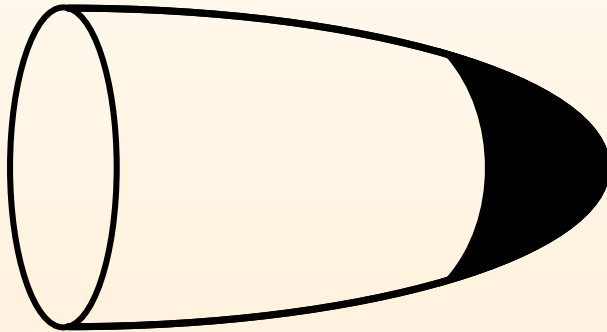
Quantum observables at the boundary of the curved space
may be calculated from propagation through curved space



Gauge/Gravity Duality: Bulk-boundary correspondence

Quantum theory at finite temperature:

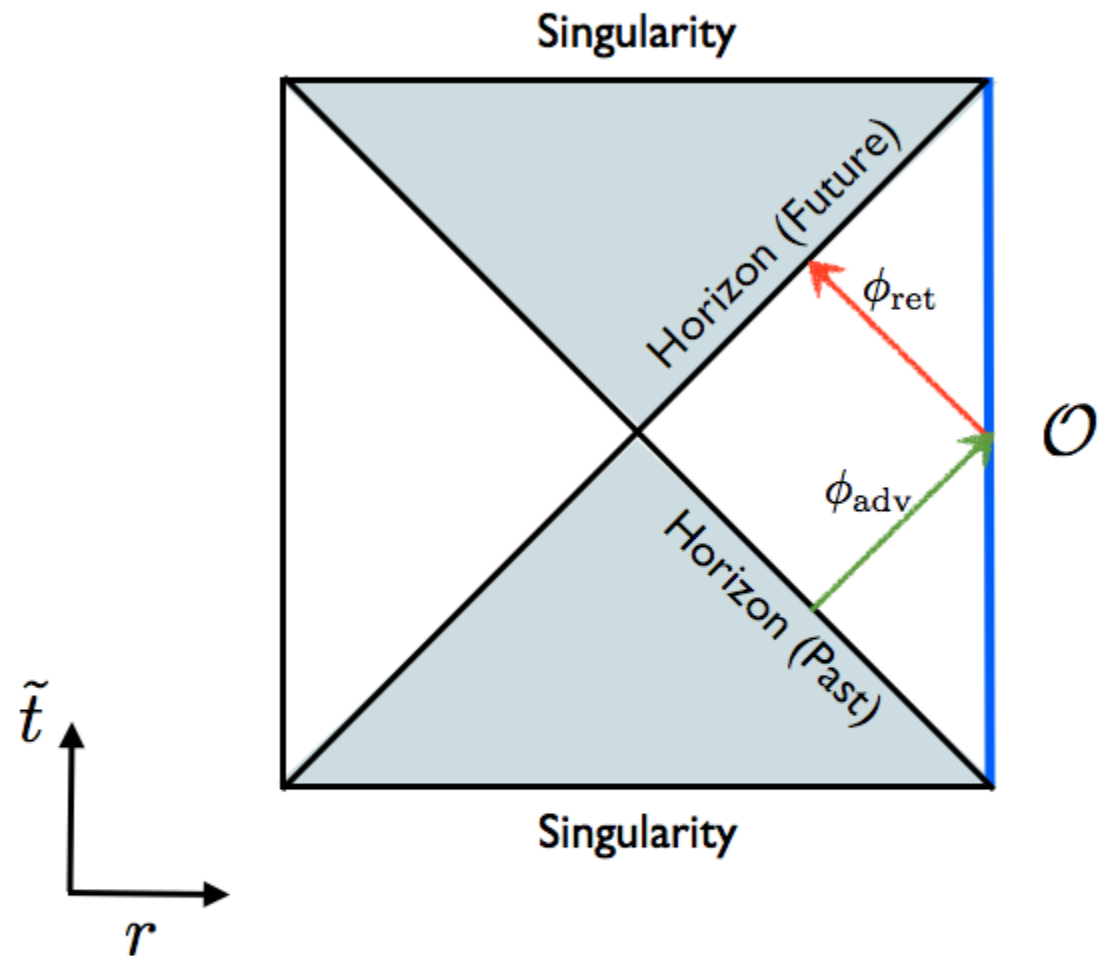
Dual to gravity theory with black hole (in Anti-de Sitter space)



Hawking temperature identified with temperature in the dual field theory

Retarded Green's Functions in Strongly Coupled Systems

Anti-de Sitter
black hole



Retarded Green's function:
$$G_{\mathcal{O}_A \mathcal{O}_B}^R = \left. \frac{\delta \langle \mathcal{O}_A \rangle}{\delta \phi_{B(0)}} \right|_{\delta \phi=0} = \frac{\delta \phi_{A(1)}}{\delta \phi_{B(0)}}$$

subject to **infalling** boundary condition at horizon

- Energy-momentum tensor $T_{\mu\nu}$ dual to graviton $g^{\mu\nu}$
- Calculate correlation function $\langle T_{xy}(x_1)T_{xy}(x_2) \rangle$ from propagation through black hole space
- Shear viscosity is obtained from **Kubo formula**:

$$\eta = -\lim_{\omega \rightarrow 0} \frac{1}{\omega} \text{Im} G_{xy,xy}^R(\omega)$$

- Shear viscosity $\eta = \pi N^2 T^3 / 8$, entropy density $s = \pi^2 N^2 T^3 / 2$

$$\frac{\eta}{s} = \frac{1}{4\pi} \frac{\hbar}{k_B T}$$

(Note: Quantum critical system: $\tau = \hbar / (k_B T)$)

Holographic hydrodynamics

Holography: From propagation of graviton in dual gravity subject to

$$S_{E-H} = \int d^{d+1}x \sqrt{-g} (R - 2\Lambda)$$

For $SU(N)$ gauge theory at infinite coupling, $N \rightarrow \infty$, $\lambda = g^2 N \rightarrow \infty$:

$$\frac{\eta}{s} = \frac{1}{4\pi} \frac{\hbar}{k_B}$$

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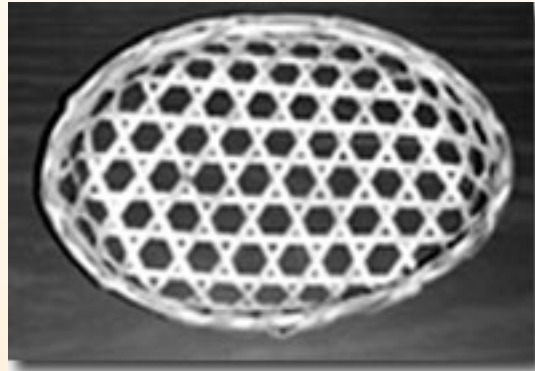
$$\frac{\eta}{s} = \frac{1}{4\pi} \frac{\hbar}{k_B}$$

Leading correction in the inverse 't Hooft coupling $\propto \lambda^{-3/2}$

From R^4 terms contributing to the gravity action

Kagome materials

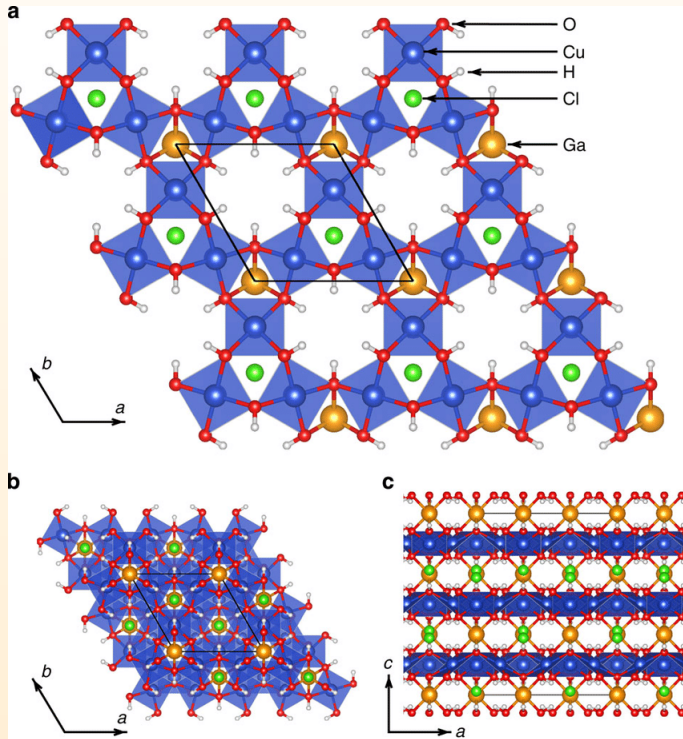
Kagome: Japanese basket weaving pattern



Source: Wikipedia

Kagome materials

Hexagonal lattice



Source: Nature

Herbertsmithite: ZnCu₃(OH)₆Cl₂



Source: Wikipedia

Scandium-Herbertsmithite

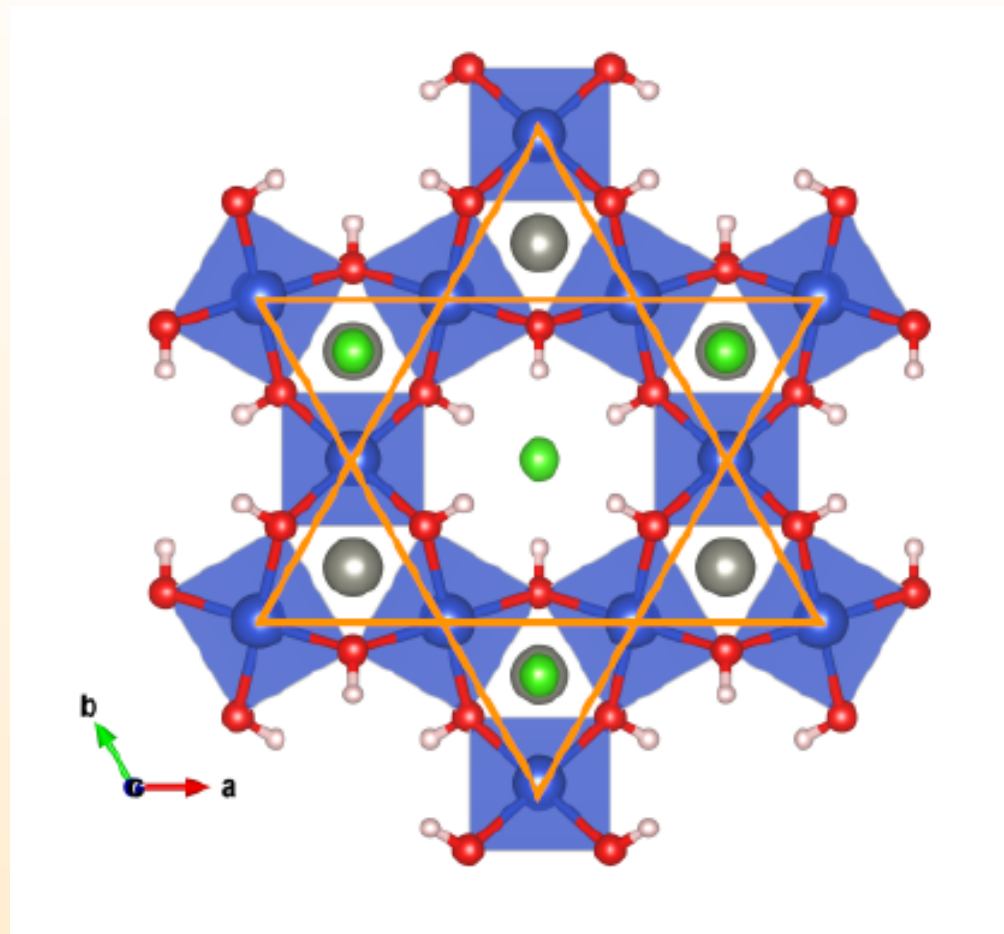
Original Herbertsmithite has Zn^{2+}

Fermi surface below Dirac point

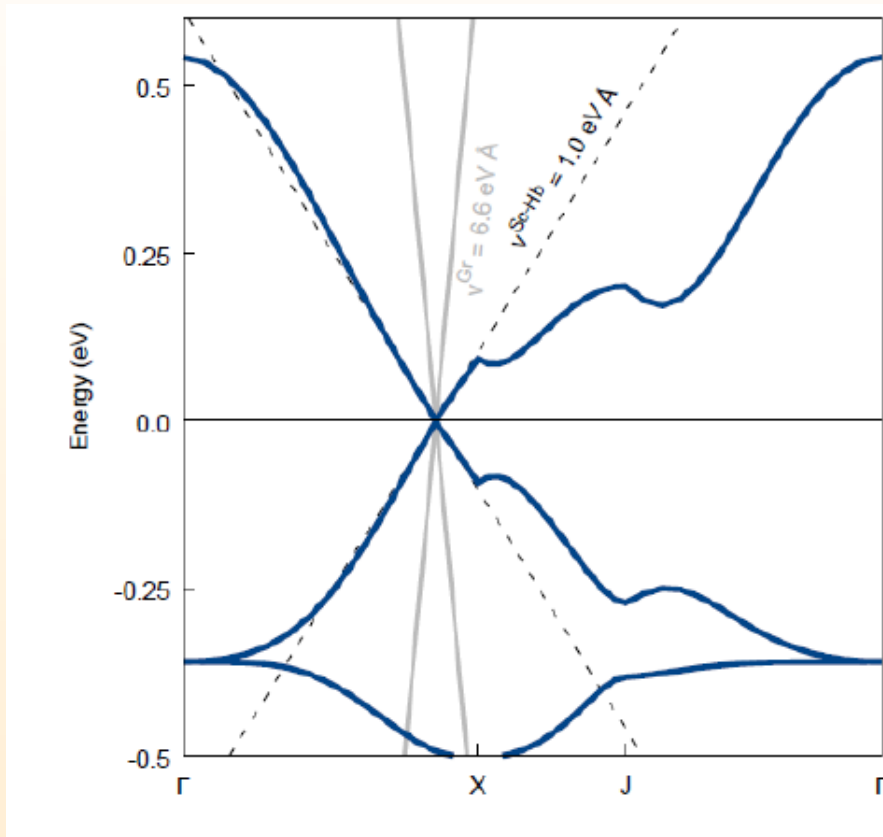
Idea: Replace Zinc by Scandium, Sc^{3+}

Places Fermi surface exactly at Dirac point

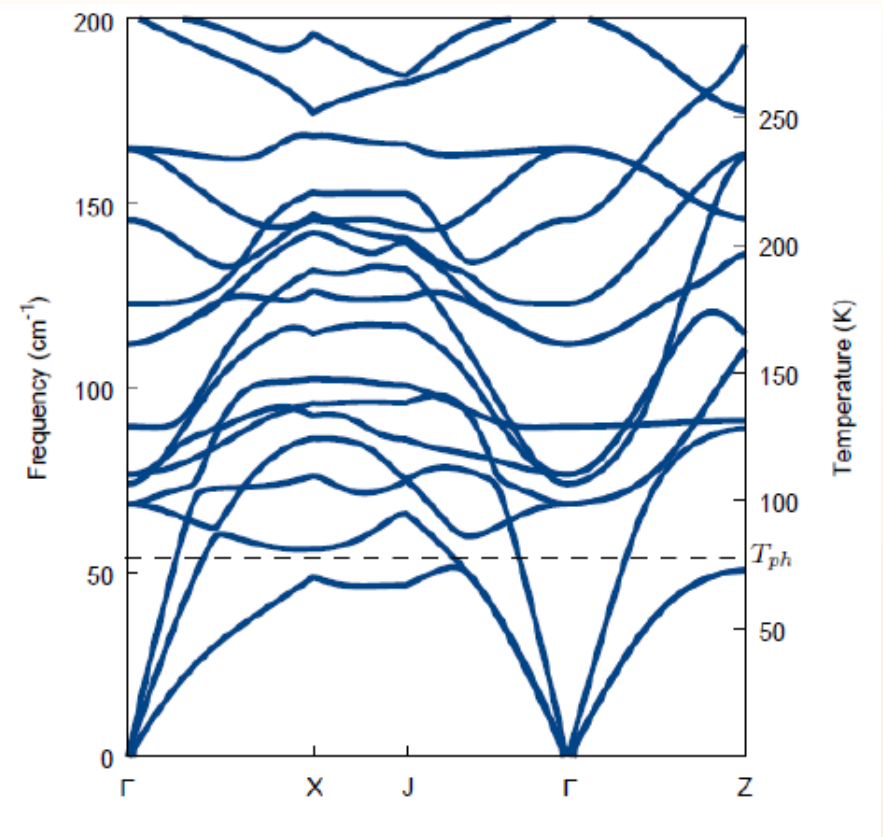
Scandium-Herbertsmithite



Scandium-Herbertsmithite



Band structure



Phonon dispersion

Scandium-Herbertsmithite

- CuO_4 plaquettes form Kagome lattice
- Low-energy physics captured by $d_{x^2-y^2}$ orbital at each Cu site
- Fermi level is at Dirac point (filling fraction $n = 4/3$)
- Orbital hybridization allows for larger Coulomb interaction (confirmed by cRPA calculation)
- Prediction: $\alpha^{\text{Sc-Hb}} = 2.9$ versus $\alpha^{\text{Graphene}} = 0.9$
- Optical phonons are thermally activated only for temperatures above $T = 80\text{K}$
- Enhanced hydrodynamic behaviour: $\ell_{ee}^{\text{Sc-Hb}} = \frac{1}{6}\ell_{ee}^{\text{graphene}}$
- Candidate to test universal predictions from holography

Estimate of the Shear viscosity

Weak coupling : Kinetic theory

$$\frac{\eta}{s} \propto \frac{1}{\alpha^2}$$

Strong coupling: Holography

Take correction

$$\frac{\eta}{s} = \frac{\hbar}{4\pi k_B} \left(1 + \frac{\mathcal{C}}{\alpha^{3/2}} \right)$$

Vary \mathcal{C} from 0.0005 to 2

AdS gravity computation: Corrections of higher order in the curvature

$$S = S_{E-H} + \int \sqrt{-g} (\gamma_2 R^2 + \gamma_3 R^3 + \gamma_4 R^4 + \dots)$$

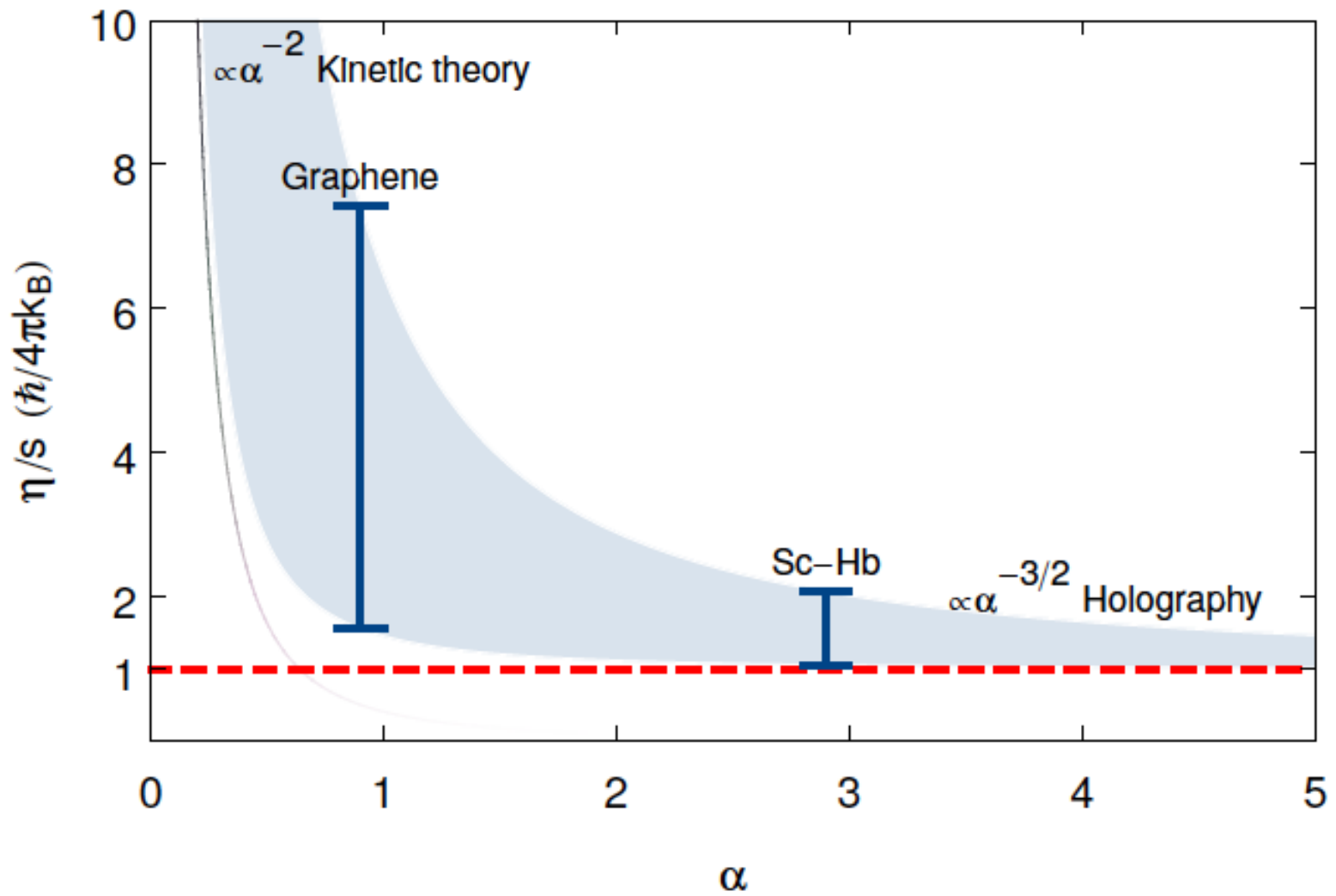
- R^2 term is topological for bulk theory in $d = 4$
- R^3 terms absent in type II supergravity parent theories
- R^4 term: Coefficient $\mathcal{O}(\lambda^{-3/2})$

$$\frac{\eta}{s} = \frac{\hbar}{4\pi k_B} \left(1 + \frac{\mathcal{C}}{\alpha^{3/2}} \right)$$

- R^4 correction is model-dependent.

We parametrize this by varying the coefficient \mathcal{C}

Estimate of the Shear viscosity



Estimate of the Reynolds number

$$\text{Re} = \left(\frac{\eta k_B}{s \hbar} \right)^{-1} \frac{k_B T u_{\text{typ}} (\eta/s)}{\hbar v_F v_F} W$$

u_{typ} typical velocity, enhanced at strong coupling

Navier-Stokes equation:

$$\frac{d\bar{v}}{dt} = -\nabla P + \frac{1}{\text{Re}} \nabla^2 \bar{v} + f$$

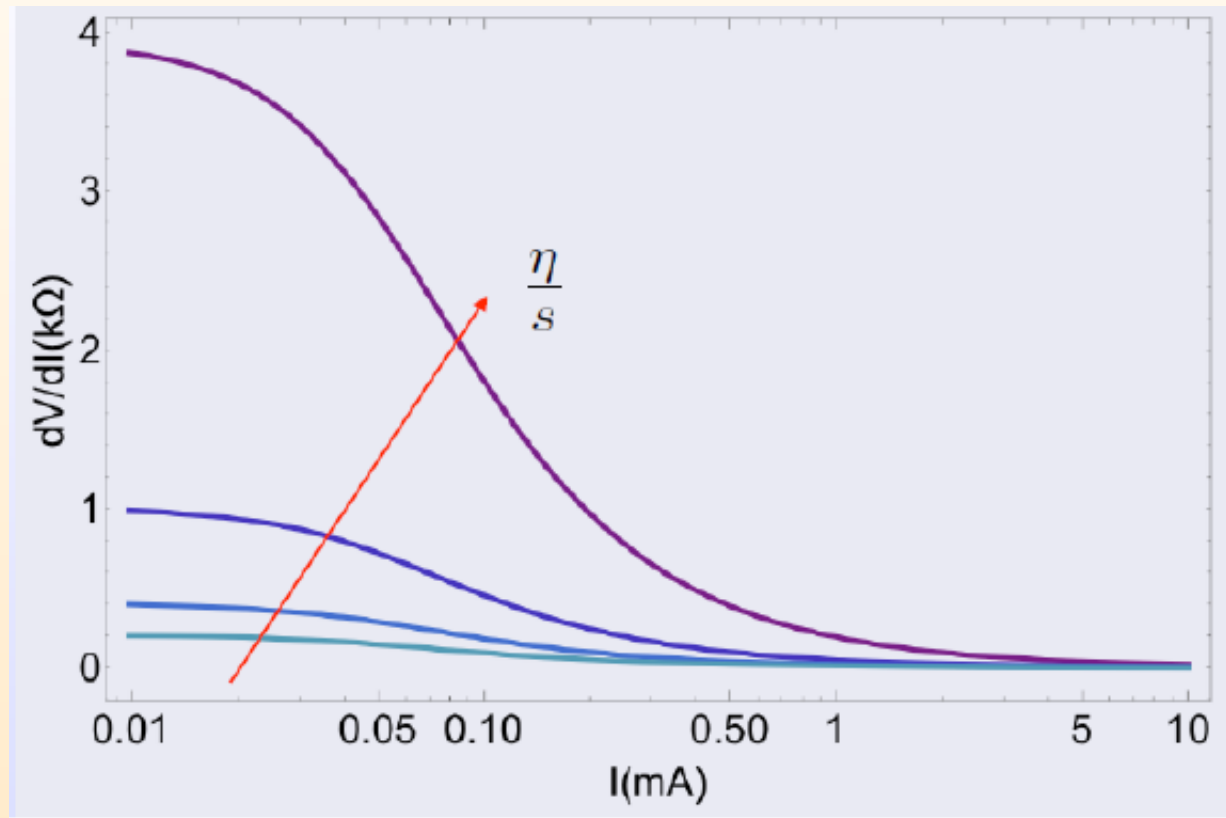
Turbulence: Reynolds number must be $\mathcal{O}(1000)$

In Sc-Hb, factor 100 larger than in graphene

Differential resistance in Poiseuille flow

J.E., Matthiakakis, Meyer, Rodriguez Fernandez PRB 2018

dV/dI increases as η/s increases



Velocity profile at varying η/s

More strongly coupled fluids flow faster

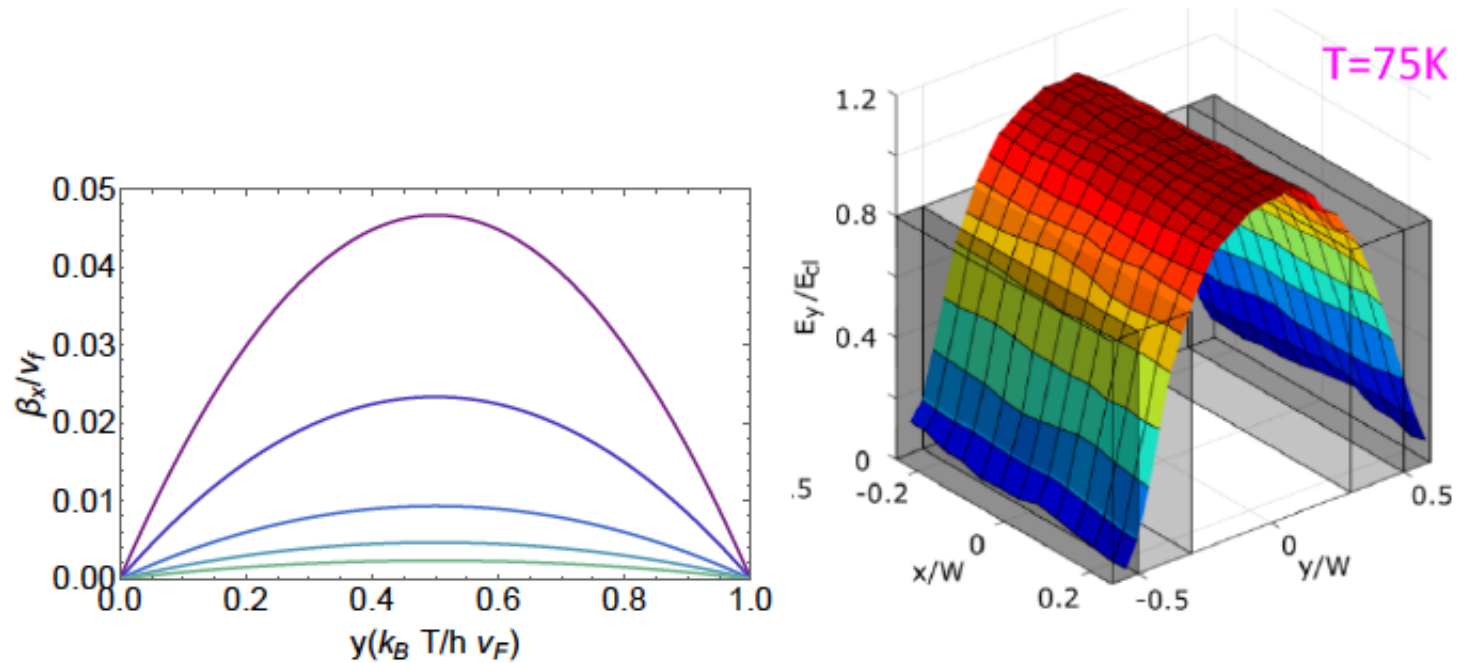


Figure: Left figure: Top curve, $\eta/s = \hbar/4\pi k_B$ (Holography). Right figure: Experimental observation of the Poiseuille flow in graphene (fig. taken from J. Sulpizio *et al* [1905.11662])

Strongly coupled fluids

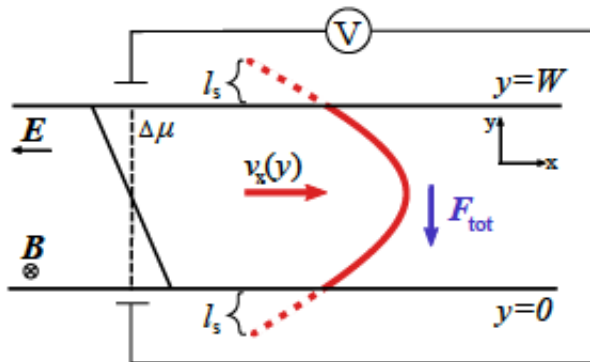
- **Strongly coupled** fluids (low η/s) **flow faster**. A promising realistic material to realize this experimentally is **Sc-Hb**
- $R(I)$ highly sensitive to the Coulomb coupling strength α_{eff} (through shear viscosity) in the hydrodynamic regime
- **Strongly coupled electron** fluids show the **smallest wire resistance** and smallest **Joule heating effect** $J \sim \sigma_Q E_x^2$

Parity breaking hydrodynamics: Hall viscosity

Functional dependence of the Hall viscosity-induced transverse voltage in two-dimensional Fermi liquids

Ioannis Matthaiakakis,^{1,*} David Rodríguez Fernández,^{1,*} Christian Tutschku,^{1,*}
Ewelina M. Hankiewicz,¹ Johanna Erdmenger,¹ and René Meyer¹

¹*Institute for Theoretical Physics and Astrophysics and Würzburg-Dresden Cluster of Excellence ct.qmat,
Julius-Maximilians-Universität Würzburg, 97074 Würzburg, Germany*



$$(\partial_t + \mathbf{v} \cdot \nabla) \rho = -\rho \nabla \cdot \mathbf{v}, \quad (\text{S1})$$

$$m_{\text{eff}} \rho (\partial_t + \mathbf{v} \cdot \nabla) \mathbf{v} = -\nabla p + \eta \nabla^2 \mathbf{v} + \eta_H \nabla^2 (\mathbf{v} \times \mathbf{e}_z) \\ + e\rho(\mathbf{E} + \mathbf{v} \times \mathbf{B}) - \frac{\rho_0 v_F m_{\text{eff}}}{l_{\text{imp}}} \mathbf{v}. \quad (\text{S2})$$

Parity breaking hydrodynamics: Hall viscosity

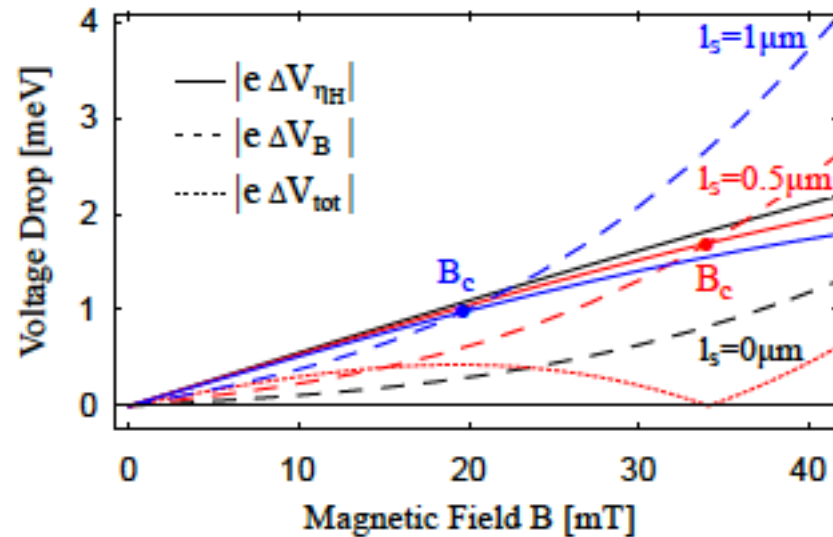


FIG. 4. Absolute values of the Lorentz ΔV_B and Hall viscous contribution $\Delta V_{\eta H}$ to the total Hall voltage ΔV_{tot} in GaAs are shown as functions of the magnetic field B for $l_s = 0, 0.5, 1.0 \mu\text{m}$. Parameters for this calculation are given in the caption of Fig. 3. For $B < B_c$, we find $|\Delta V_{\eta H}|/|\Delta V_B| > 1$, whereas otherwise $|\Delta V_{\eta H}|/|\Delta V_B| < 1$. At $B = B_c$, the ratio $\Delta V_{\eta H}/\Delta V_B = -1$ implying a vanishing Hall voltage $\Delta V_{\text{tot}} = 0$.

Conclusion and outlook

- Scandium-substituted Herbertsmithite has predicted coupling $\alpha_{\text{eff}} = 2.9$
- Factor 3.2 larger than Graphene
- May reach region of robust hydrodynamics in solids
- Smaller ratio of η/s - parameter region where gauge/gravity duality applies
- Strongly coupled fluids flow faster
- Poiseuille flow
- Cancellation of Hall viscosity induced voltage with standard Hall voltage