

# Hawking's Singularity Theorem in General Relativity

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LisMath Seminar

Complexo Interdisciplinar da Universidade de Lisboa  
Lisbon, 13<sup>th</sup> March 2015

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# Outline

- 1 Fundamentals of Lorentzian geometry
- 2 Fundamentals of general relativity
  - The Einstein field equations
  - Spacetimes with timelike singularities
- 3 Causality
- 4 Hawking's singularity theorem

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# Lorentzian manifolds

- A **Lorentzian manifold** is a pair  $(M, g)$ , where  $M$  is a differentiable manifold and  $g$  is a nondegenerate symmetric 2-tensor with signature  $(- + \dots +)$  (called the **metric**).
- $M$  admits a Lorentzian metric iff it is noncompact or  $\chi(M) = 0$ .
- The simplest example is the **Minkowski spacetime**:

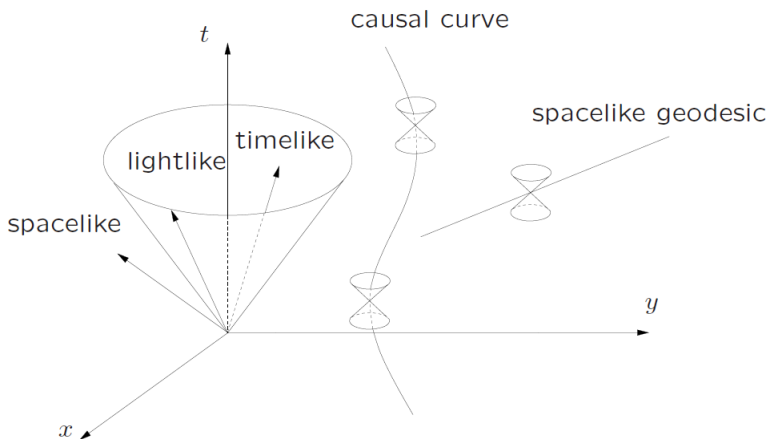
$$M = \mathbb{R}^{n+1}, \quad g = -dx^0 \otimes dx^0 + dx^1 \otimes dx^1 + \dots + dx^n \otimes dx^n$$

# Comparison with Riemannian geometry

- Most classic results from Riemannian geometry are valid.
  - But not all:
    - Triangle inequality (twin paradox);
    - Hopf-Rinow theorem (geodesic completeness);
    - Vectors  $v$  come in three types:
      - **Timelike:**  $g(v, v) < 0$
      - **Lightlike:**  $g(v, v) = 0$
      - **Spacelike:**  $g(v, v) > 0$
- ( $v$  is **causal** if it is not spacelike).

# Comparison with Riemannian geometry

(Minkowski spacetime)





# Basic concepts

- A curve  $c: I \subset \mathbb{R} \rightarrow M$  is called **timelike**, **lightlike**, **spacelike** or **causal** if so is  $\dot{c}$ .
- If  $c$  is a **geodesic**, then its type cannot change.
- $\tau(c)$ , the **length** of the timelike curve  $c: [a, b] \rightarrow M$ , is the **proper time** between **events**  $c(a)$  and  $c(b)$ .

## Definition

$(M, g)$  is **singular** if it is not geodesically complete.

# Einstein field equations

- Einstein postulated the equation system  $Ric - \frac{S}{2}g = T$  and Hilbert derived it from the variational principle.
- The spherically symmetric vacuum ( $T = 0$ ) case is the **Schwarzschild solution**. In 1 + 3 dimensions,

$$g = - \left(1 - \frac{2M}{r}\right) dt \otimes dt + \left(1 - \frac{2M}{r}\right)^{-1} dr \otimes dr + r^2 (d\theta \otimes d\theta + \sin^2 \theta d\varphi \otimes d\varphi)$$

The surface  $\{r = 2M\}$  is the **black hole** horizon. In the inner region, **every** timelike geodesic reaches the curvature singularity  $r = 0$  in **finite time**.

# Einstein field equations

- Consider the Riemannian 3-manifold  $(\Sigma, h)$  with constant curvature  $k$ .  $(\mathbb{R} \times \Sigma, -dt \otimes dt + a^2(t)h)$  is the FLRW universe.
- Friedmann used the Einstein equations to obtain

$$\frac{\dot{a}^2}{2} - \frac{\alpha}{a} = -\frac{k}{2}$$

where the density is  $\rho = \frac{6\alpha}{a^3}$ .

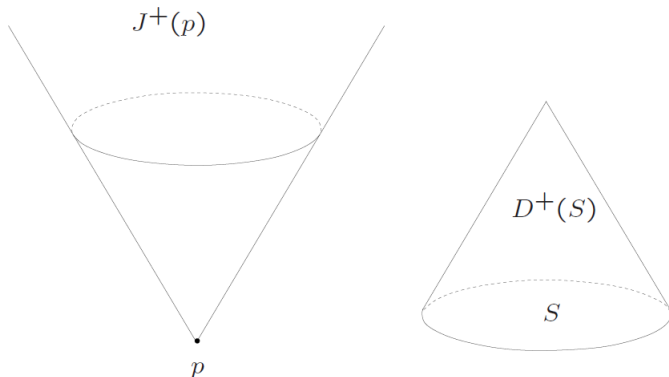
- If  $\alpha > 0$ , then  $a$  either blows up or goes to zero in finite time (**Big Bang**).

# Basic concepts

- Timelike  $v$  and  $w$  have the same **time orientation** if  $\langle v, w \rangle < 0$ .
- $(M, g)$  is **time-orientable** if  $\exists$  nonvanishing timelike  $T \in \mathfrak{X}(M)$ .  
A timelike vector  $v \in T_p M$  is **future-pointing** if it has the same time orientation as  $T_p$ . A timelike curve  $c$  is **future-directed** if  $\dot{c}$  is future-pointing.
- **Chronological future** of  $p \in M$  – set  $I^+(p)$  of all points to which  $p$  can be connected by a by a future-directed timelike curve.  
**Causal future** of  $p \in M$  – set  $J^+(p)$  of all points to which  $p$  can be connected by a by a future-directed causal curve.

# Basic concepts

(Minkowski spacetime)



# Main results

## Proposition

**Locally, the topological properties** of a time-oriented Lorentzian manifold are similar to those of the Minkowski spacetime. In particular:

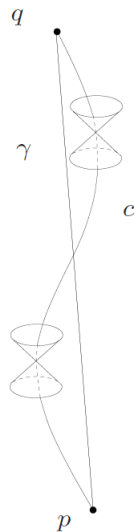
- $q \in I^+(p)$  and  $r \in I^+(q) \Rightarrow r \in I^+(p)$ ;
- $I^+(p)$  is open.

## Proposition (twin paradox)

Let  $(M, g)$  be time-oriented. Every  $p_0 \in M$  has a geodesically convex open neighbourhood  $V \subset M$  such that, in  $(V, g|_V)$ , timelike geodesics have maximal length.

# The twin paradox

$$\tau(\gamma) > \tau(c)$$



## Some further concepts

- $(M, g)$  satisfies the **chronology condition** if it does not contain closed timelike curves (cannot be compact).
- $(M, g)$  is **stably causal** if  $\exists$  a **global time function**, i.e., a smooth function  $t: M \rightarrow \mathbb{R}$  such that  $\text{grad } t$  is timelike.
- A future-directed causal curve  $c: ]a, b[ \rightarrow M$  is **future-inextendible** if  $\lim_{t \rightarrow b} c(t)$  does not exist.
- The **future domain of dependence** of  $S \subset M$  is

$$D^+(S) = \{p \in M \mid \text{every past-directed inextendible causal curve intersects } S\}$$

### Proposition

If  $p \in D^+(S)$ , then  $D^+(S) \cap J^-(p)$  is compact.



# Essential conditions

## Definition

$(M, g)$  is said to be **globally hyperbolic** if it is stably causal and all time slices  $S_a = t^{-1}(a)$  satisfy  $D(S_a) = M$ .

- $(M, g)$  will always be assumed to be globally hyperbolic and  $S$  will always denote a time slice.
- **Strong energy condition:**  $Ric(V, V) \geq 0$  for every timelike  $V$ .
- Let  $n$  be the future-pointing unit normal of  $S$  and  $c_p$  the geodesic orthogonal to  $S$  and tangent to  $n$  at  $p \in S$ .  
The critical values of the **exponential map** ( $\exp(t, p) = c_p(t)$ ) are said to be **conjugate points** to  $S$ .

# Main results

- If  $c_p$  has no conjugate points between  $c_p(0) = p$  and  $c_p(t_0)$ ,  $\exists$  an open neighbourhood  $V$  of  $c_p([0, t_0])$  where  $\exists$  a family of timelike geodesics tangent to a unit vector field  $X = -\text{grad } t$ .
- $\theta = \text{div } X$  is called the **expansion**.

## Proposition

Let  $(M, g)$  satisfy the SEC,  $p \in S$  and  $\theta = \theta_0 < 0$ .

Then, if  $c_p$  can be extended to a length  $\tau_0 = -n/\theta_0$  to the future of  $S$ , it has a conjugate point.

# Main results

## Proposition

Let  $c$  be a geodesic through  $p \in M$  orthogonal to  $S$ .

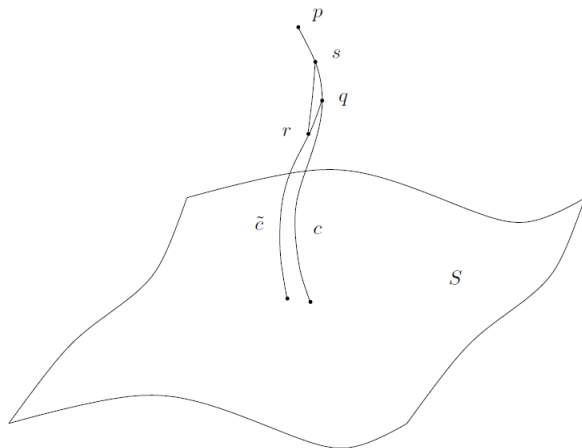
Then, if  $\exists$  a conjugate point between  $S$  and  $p$ ,  $c$  does not maximize length.

## Proposition

$\exists$  a timelike curve connecting  $S$  to  $p$  which has maximal length.

This curve is a geodesic.

# Conjugate points and (non)-maximizing geodesics



# Hawking's singularity theorem

## Theorem

If  $(M, g)$  satisfies the SEC and  $\theta \leq \theta_0 < 0$ , then it is singular.

## Proof

- Suppose  $\exists$  future-directed timelike geodesic  $c$  orthogonal to  $S$ , extendible to proper time  $\tau(c) = \tau_0 + \epsilon > \tau_0 = -\frac{n}{\theta_0}$ ;
- Then,  $\exists$  maximal timelike geodesic  $\gamma$  connecting  $S$  to  $c(\tau_0 + \epsilon)$ ;
- As such,  $\tau(\gamma) \geq \tau_0 + \epsilon$ ;
- However,  $\gamma$  has a conjugate point at a distance of at most  $\tau_0$  and ceases to be maximizing beyond this point;
- Hence, this yields a contradiction.



# References

- Papers:

- ① S. Hawking, *The Occurrence of Singularities in Cosmology. III. Causality and singularities*, Proc. Roy. Soc. Lon. A 300 (1967), 187201.

- Books:

- ① G. Naber, *Spacetime and Singularities: An Introduction*, Cambridge University Press, 1988.
- ② L. Godinho and J. Natário, *An Introduction to Riemannian Geometry: With Applications to Mechanics and Relativity* (Universitext), Springer, 2014.