

The symplectic geometry of the restricted three-body problem

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Practical aspects and applications

Astrodynamics and space missions

- The CR3BP is one of the basic models in **astrodynamics** and **space mission design**.

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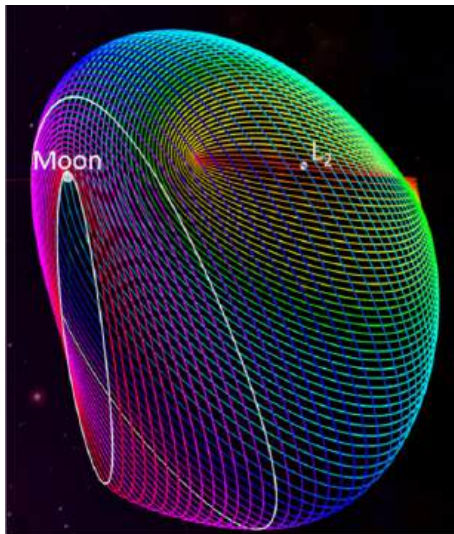
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- They come in families (varying the energy) connected through each other through *bifurcations*, codified in a graph of “options” for a putative mission.
- There is an infinitude of periodic orbits, many of which are used for real missions.
- The basic families in the CR3BP can be used to construct nominal orbits which are much more complex, using high-fidelity (*ephemeris*) models deforming the CR3BP.

Halo orbits



Halo orbits for Earth–Moon.

Current context: systems of interest

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- Jupiter-Europa; and
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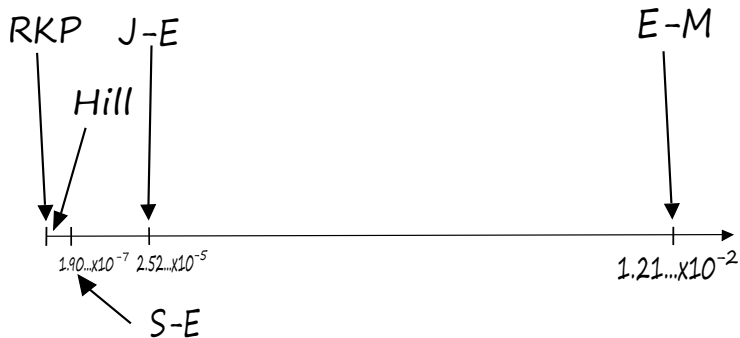
- Jupiter-Europa; and
- Saturn-Enceladus.

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In the context of deep space exploration, NASA's ARTEMIS program (e.g. Gateway, halo orbits) is interested in:

- Earth–Moon.

Systems of interest



Values of μ for different systems of interest.

Symplectic Data Analysis

Philosophy

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Symplectic Data Analysis (SDA).

What follows is based on work with:

- Urs Frauenfelder (Augsburg)
- Dan Scheeres (aerospace engineer, CO Boulder)
- Otto van Koert (Seoul)
- Cengiz Aydin (Heidelberg)
- Bhanu Kumar (Michigan)
- Dayung Koh (JPL-NASA: Gateway, Europa Clipper, Psyche)

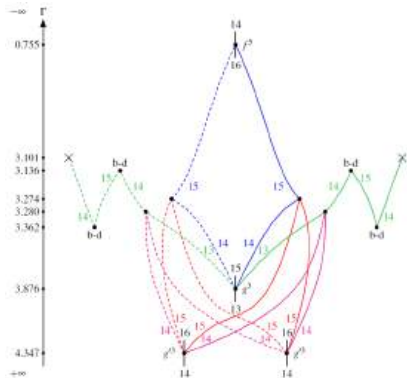
The symplectic toolkit

With Urs Frauenfelder and Cengiz Aydin, we developed a “**symplectic toolkit**” for engineers (that made it to *Quanta Magazine*):

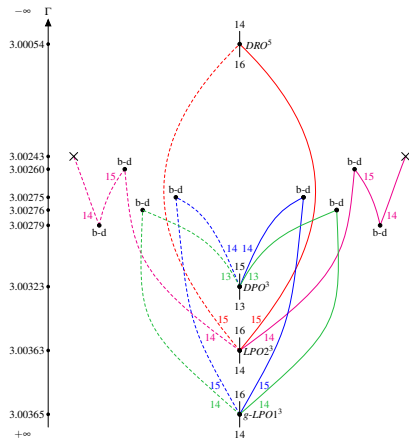
- (1) **The B-signs:** A \pm sign, helps predict/obstruct orbits and bifurcations.
- (2) **Global topological methods:** for bifurcations and stability.
- (3) **Conley-Zehnder index:** organizes families.
- (4) **Floer numbers:** A test to detect potentially overlooked families.

Bifurcation graphs

The symplectic toolkit helps organize and find orbits in a systematic way.



Hill lunar problem.

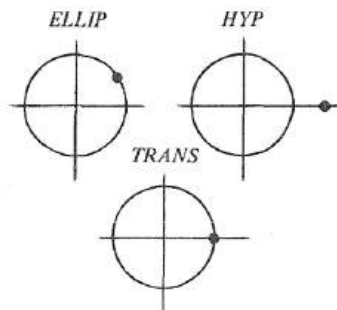
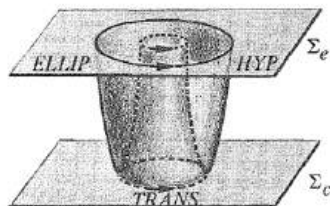


Jupiter-Europa.

Bifurcations

Monodromy matrix: symplectic matrix (linearization of the flow along an orbit).

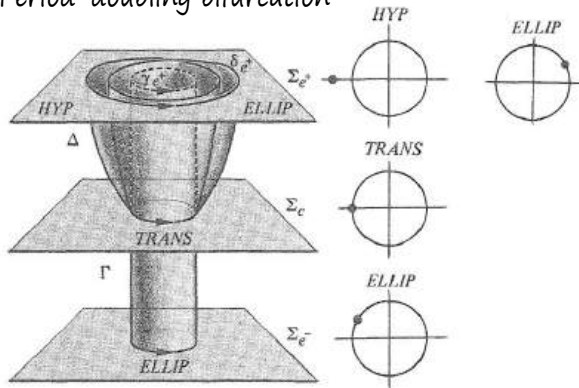
Bifurcation: e-value 1 is crossed in a family.



Creation or birth/death.

Bifurcations

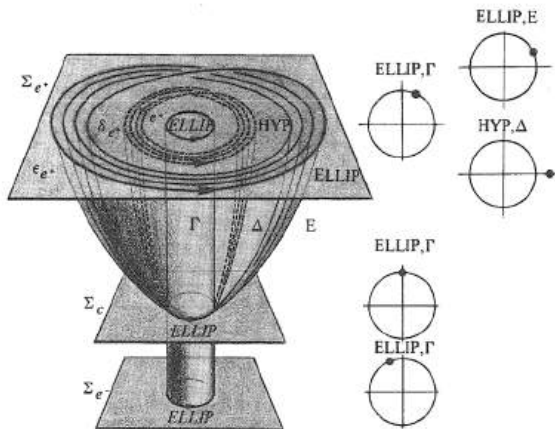
Period-doubling bifurcation



"Foundations of Mechanics", Abraham-Marsden.

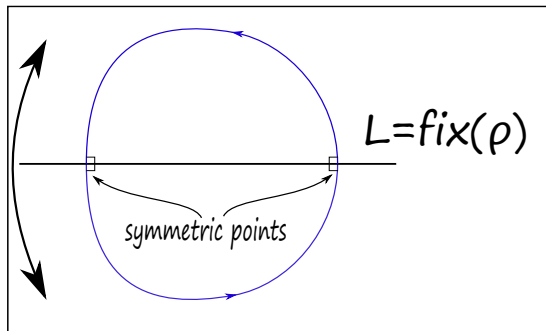
Period-doubling bifurcation or subtle division.

Bifurcations



Emission, or k -fold bifurcation ($k = 4$).

Symmetries



Symmetric periodic orbits.

Wonenburger matrices

The monodromy matrix of a symmetric orbit at a **symmetric point** is a **Wonenburger** matrix:

$$M = M_{A,B,C} = \begin{pmatrix} A & B \\ C & A^T \end{pmatrix} \in Sp(2n), \quad (1)$$

where

$$B = B^T, \quad C = C^T, \quad AB = BA^T, \quad A^T C = CA, \quad A^2 - BC = \mathbb{1}. \quad (2)$$

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The eigenvalues of M are determined by those of the first block A .

Lemma

- λ *e-val* of $M \rightsquigarrow$ its **stability index** $a(\lambda) = \frac{1}{2}(\lambda + 1/\lambda)$ *e-val* of A .
- a *e-val* of $A \rightsquigarrow \lambda(a) = a + \sqrt{a^2 - 1}$ *e-val* of M .

Toolkit

Global topological methods

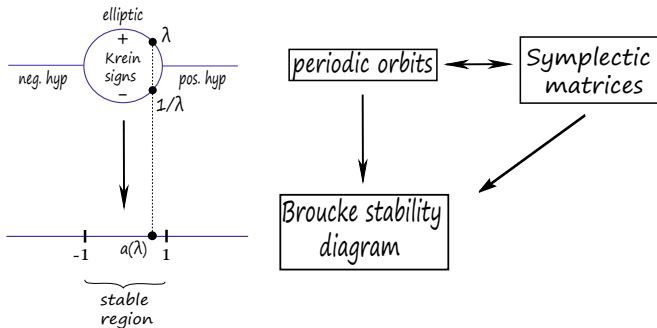
These methods encode:

- Bifurcations;
- stability;
- eigenvalue configurations;
- obstructions to existence of regular families;
- *B-signs*,

in a visual and resource-efficient way.

Broucke's stability diagram: 2D

Let $n = 2$, λ eigenvalue of monodromy, with **stability index** $a(\lambda) = \frac{1}{2}(\lambda + 1/\lambda)$.



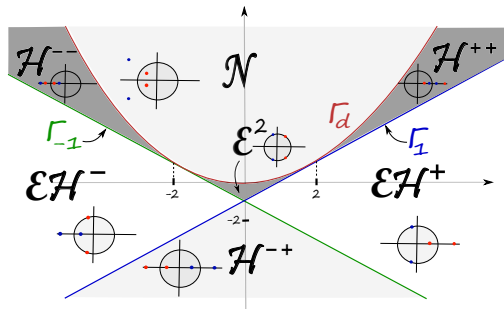
If we know that two points lie in different components, then one should expect bifurcations in any path between them.

Broucke's stability diagrams: 3D

Let $n = 3$. Given $M_{A,B,C}$, its **stability point** is $p = (\text{tr}(A), \det(A)) \in \mathbb{R}^2$.

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Let $n = 3$. Given $M_{A,B,C}$, its **stability point** is $p = (\text{tr}(A), \det(A)) \in \mathbb{R}^2$. The plane splits into regions corresponding to the eigenvalue configuration:



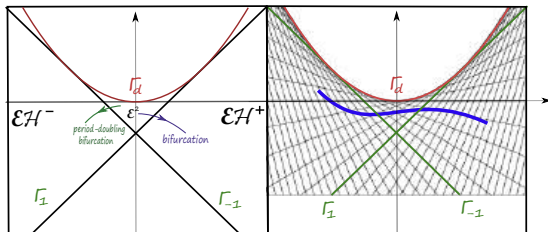
- $\Gamma_{\pm 1} = \text{e-val } \pm 1$.
- $\Gamma_d = \text{double e-val.}$
- $\mathcal{E}^2 = \text{doubly elliptic (stable region).}$
- etc.

Bifurcations in the Broucke diagram

An orbit family $t \mapsto x_t$ induces a path $t \mapsto p_t \in \mathbb{R}^2$ of stability points.

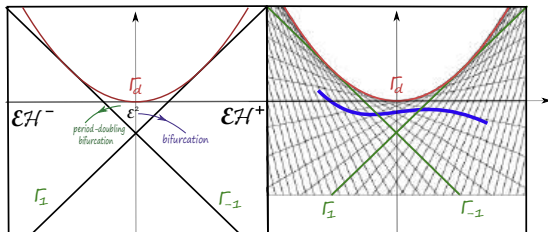
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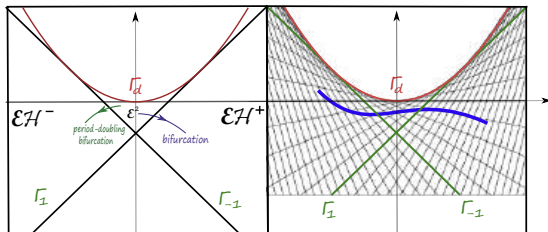
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If we know that two points lie in different components, then one should expect bifurcations in any path between them.

B-signs

Let x be a symmetric orbit with monodromy

$$M = \begin{pmatrix} A & B \\ C & A^T \end{pmatrix}.$$

Assume λ simple, elliptic or hyperbolic e-value. Let v satisfy $A^T v = a(\lambda) \cdot v$. The **B-sign** of λ is

$$\epsilon(\lambda) = \text{sign}(v^T B v) = \pm 1.$$

- $n = 2$ two B -signs, one for each symmetric point.
- $n = 3$ two **pairs** of B -signs, one for each symmetric point and each eigenvalue.

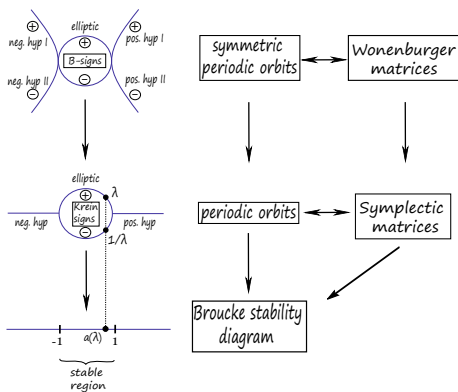
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Theorem (Frauenfelder–M. '23)

A planar symmetric orbit is negative hyperbolic iff the B -signs of its two symmetric points differ.

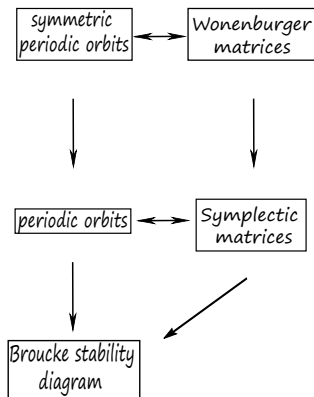
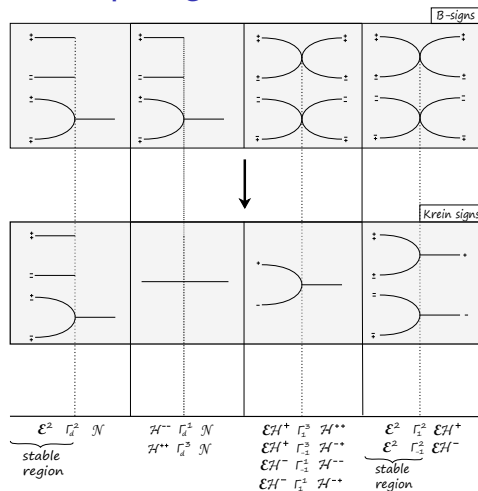
Global topological methods: GIT sequence, 2D

GIT sequence = refinement of Broucke diagram for **symmetric** orbits.



- B -signs “separate” hyperbolic branches, for symmetric orbits.
- If two points lie in the same component, but B -signs differ, one should *also* expect bifurcation in any path joining them.

Global topological methods: GIT sequence, 3D



The branches are two-dimensional, and come together where we cross from one region to another. If the B -signs are not mixed, then a stable orbit cannot cross from \mathcal{E}^2 to \mathcal{N} (Krein–Moser theorem).

Conley–Zehnder index

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- CZ-index stays constant if no bifurcation occurs;
- Then helps understand which families of orbits connect to which.
- It can be thought of as an intersection number with the *Maslov cycle*, consisting of symplectic matrices with 1 as an eigenvalue.

Conley–Zehnder index

$n=2$ x planar orbit with (reduced) monodromy M_x , x^k k -fold cover.

- **Elliptic case:**

$$M_x \sim \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix},$$

Then

$$\mu_{CZ}(x^k) = 1 + 2 \cdot \lfloor k \cdot \theta / 2\pi \rfloor$$

In particular, it is odd, and jumps by ± 2 if the e-val 1 is crossed.

- **Hyperbolic case:**

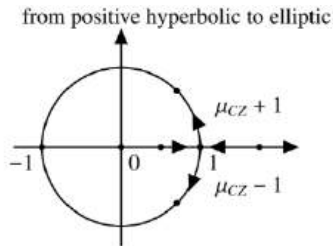
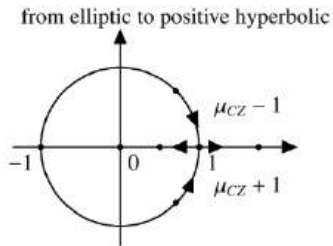
$$M_x \sim \begin{pmatrix} \lambda & 0 \\ 0 & 1/\lambda \end{pmatrix},$$

Then

$$\mu_{CZ}(x^k) = k \cdot n,$$

where linearized flow rotates eigenspaces by angle $\frac{\pi n t}{T}$, with n even/odd if x pos./neg. hyp.

CZ-jumps



μ_{CZ} jumps by ± 1 when crossing 1, according to direction of bifurcation. If it stays elliptic, the jump is by ± 2 .

Conley–Zehnder index

$n = 3$, $x \in \mathbb{R}^2$ planar orbit in 3BP, then

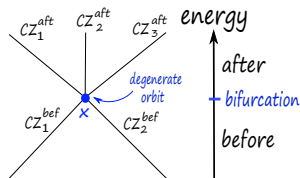
$$\mu_{CZ}(x) = \mu_{CZ}^p(x) + \mu_{CZ}^s(x),$$

a spatial and a planar index.

- Planar to planar bifurcations correspond to jumps in μ_{CZ}^p .
- Planar to spatial bifurcations correspond to jumps of μ_{CZ}^s .

The jump is determined by the B -sign. The CZ-index may be computed directly (numerically or via above formulas), or via deformation of known cases.

Floer numerical invariants



Given a bifurcation at x , the *Floer number* of x is

$$\chi(x) = \sum_i (-1)^{CZ_i^{bef}} = \sum_j (-1)^{CZ_j^{aft}}.$$

The sum on the LHS is over (**good**) orbits *before* bifurcation, and RHS is over (**good**) orbits *after* bifurcation.

Invariance

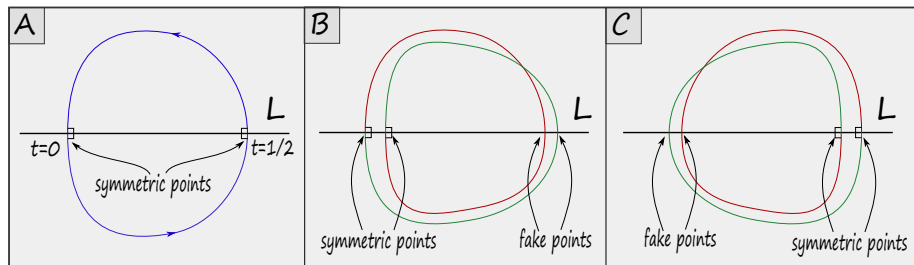
The fact that the sums agree before and after –*invariance*– follows from very deep results in **Floer theory** from symplectic geometry.



In Memoriam Andreas Floer, 1956-1991.

The Floer number can be used as a **test**: if the sums do *not* agree, we know the algorithm missed an orbit.

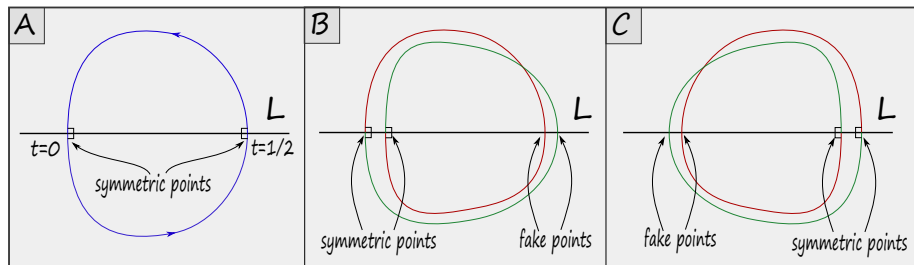
Example: symmetric period doubling bifurcation



The simple symmetric orbit x goes from elliptic to negative hyperbolic.

- A priori there could be two bifurcations for each symmetric point (B or C).

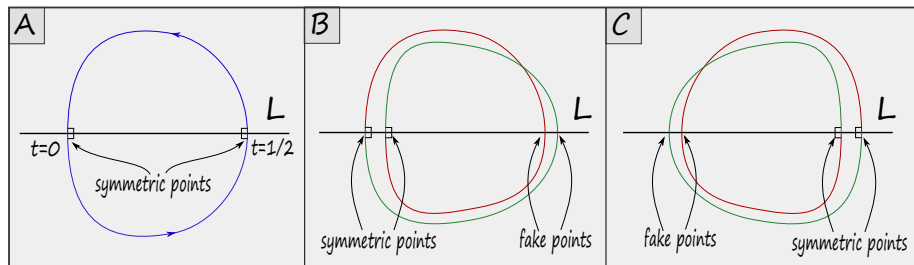
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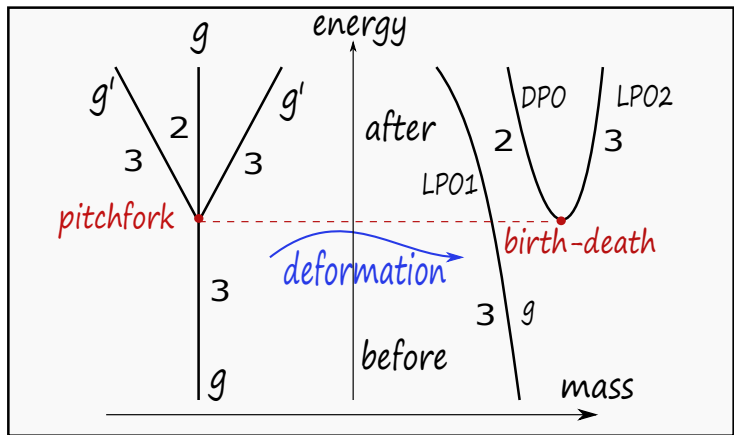
- A priori there could be two bifurcations for each symmetric point (B or C).
- Invariance of $\chi(x^2)$ implies only one can happen (note x^2 is *bad*).
- Bifurcation happens at the symmetric point in which the B -sign does *not* jump.

Summary of toolkit

- (1) **The B-signs:** a \pm sign associated to elliptic/hyperbolic orbits, which helps predict bifurcations.
- (2) **Global topological methods:** the *GIT-sequence*, a topological refinement of *Broucke's stability diagram*, which encodes bifurcations and stability of orbits.
- (3) **Conley-Zehnder indices:** an integer associated to a (non-degenerate) orbit which only jumps at bifurcation, and so predicts which families connect to which.
- (4) **Floer numerical invariants:** numerical counts of orbits that stay the same before and after a bifurcation, and so help predict existence of orbits.

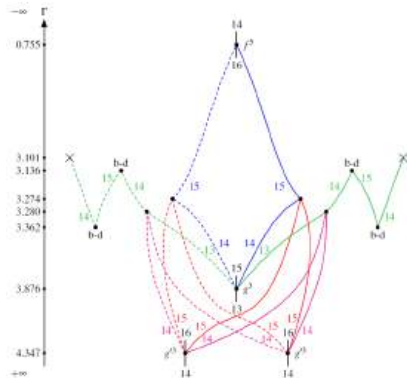
Numerical work

Example: pitchfork bifurcation

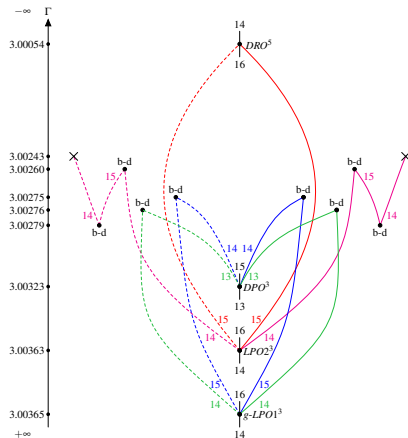


Hill's problem has plenty of symmetry: a (nno-generic) pitchfork bifurcation (Hénon) deforms to a generic bifurcation in Jupiter–Europa. Birth-death branch is hard to detect. The Floer number is easy to compute.

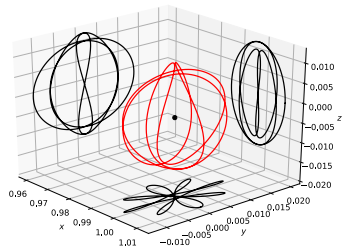
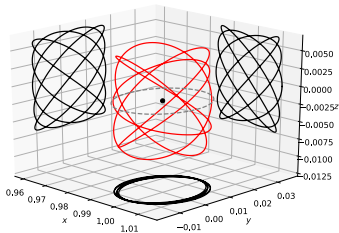
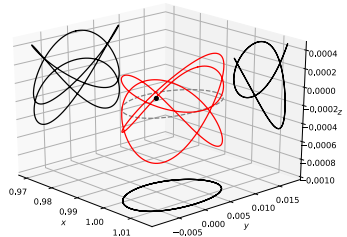
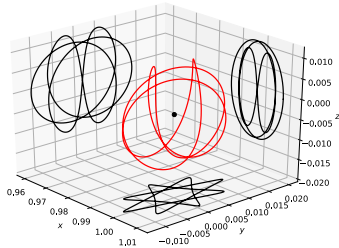
Bifurcation graphs



Hill lunar problem.

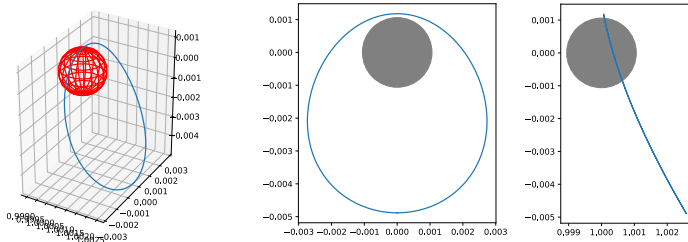


Jupiter-Europa.

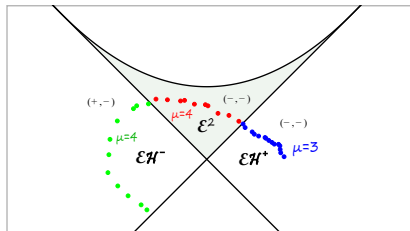


Jupiter-Europa: Spatial connection between retrograde planar orbit to direct planar orbit, CZ-index 15.

Halo orbits in Saturn-Enceladus



Halo orbit, altitude = 29km, CZ = 4 (Enceladus Orbilander, NASA).



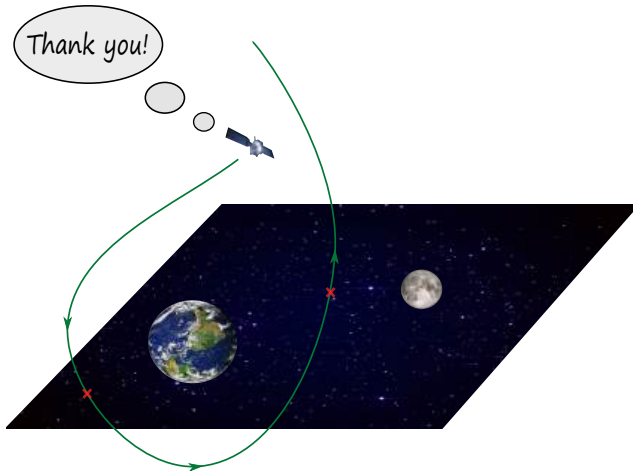
GIT plot of Halo orbits with CZ-index.

Remarks

- Otto van Koert produced a Python **CZ-index calculator** for CR3BP (available in GitHub); Bhanu Kumar has a MATLAB version.

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- Otto van Koert produced a Python **CZ-index calculator** for CR3BP (available in GitHub); Bhanu Kumar has a MATLAB version.
- With Dan Scheeres and Gavin Brown, we applied these tools to the HR4BP, a higher-fidelity, time-dependent deformation of the CR3BP (Astrodynamics Specialists Conference 2024).



References I



P. Albers, J. Fish, U. Frauenfelder, H. Hofer, O. van Koert.
Global surfaces of section in the planar restricted 3-body problem.
Arch. Ration. Mech. Anal. 204(1) (2012), 273–284.



Albers, Peter; Frauenfelder, Urs; van Koert, Otto; Paternain, Gabriel P.
Contact geometry of the restricted three-body problem.
Comm. Pure Appl. Math. 65 (2012), no. 2, 229–263.



Wanki Cho, Hyojin Jung, Geonwoo Kim.
The contact geometry of the spatial circular restricted 3-body problem.
[arXiv:1810.05796](https://arxiv.org/abs/1810.05796)



C. Conley.
On Some New Long Periodic Solutions of the Plane Restricted Three Body Problem.
Comm. Pure Appl. Math. 16 (1963), 449–467.

References II



Frauenfelder, Urs; Koh, Dayung; Moreno, Agustin.

On Floer-type numerical invariants, GIT quotients, and orbit bifurcations of real-life planetary systems.

To Appear.



U. Hryniewicz and Pedro A. S. Salomão.

Elliptic bindings for dynamically convex Reeb flows on the real projective three-space.

Calc. Var. Partial Differential Equations 55 (2016), no. 2, Art. 43, 57 pp.



U. Hryniewicz, Pedro A. S. Salomão and K. Wysocki.

Genus zero global surfaces of section for Reeb flows and a result of Birkhoff.

arXiv:1912.01078.

References III



Frauenfelder, Urs; Moreno, Agustin.

On GIT quotients of the symplectic group, stability and bifurcations of symmetric orbits

Preprint arXiv:2109.09147.



Koh, Dayung; Anderson, Rodney L.; Bermejo-Moreno, Ivan.

Cell-mapping orbit search for mission design at ocean worlds using parallel computing.

The Journal of the Astronautical Sciences, Volume 68, Issue 1, p.172-196.



McGehee, Richard Paul.

Some homoclinic orbits for the restricted three-body problem.

Thesis (Ph.D.), The University of Wisconsin - Madison. ProQuest LLC, Ann Arbor, MI, 1969. 63 pp.0.

References IV



Moreno, Agustin.

Pseudo-holomorphic dynamics in the restricted three-body problem.

Preprint [arXiv:2011.06568](#)



Moreno, Agustin; van Koert, Otto.

Global hypersurfaces of section in the spatial restricted three-body problem.

To appear in Nonlinearity.



Moreno, Agustin; van Koert, Otto.

A generalized Poincaré-Birkhoff theorem.

J. Fixed Point Theory Appl. 24 (2022), no. 2, Paper No. 32.



Poincaré, H.

Sur un théoreme de géométrie.

Rend. Circ. Matem. Palermo 33, 375–407 (1912).