

Def Assume (u_L, u_R) satisfies RH-conditions: $s[u] = [f]$.

We say that (u_L, u_R) satisfies the "Lax (entropy) condition" if

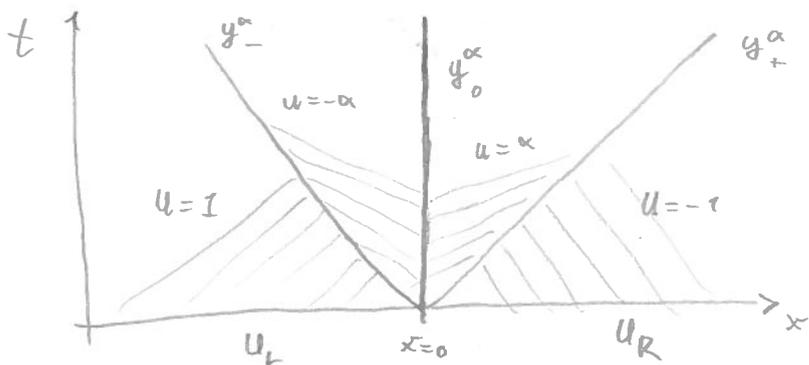
$f'(u_L) > s > f'(u_R)$

(LC)

Remark:

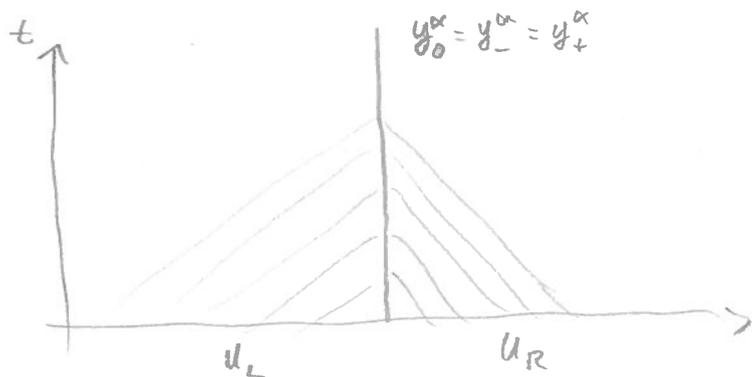
We get uniqueness of discontinuous solutions by only allowing for discontinuities that satisfy (LC).

In Example 1, only the solution u_α for $\alpha = 1$ meets (LC)!



Case
 $\alpha \geq 1$

Discontinuity across y_0^α violates (LC) $\forall \alpha > 1$, since $f'(u_L) = -\alpha, s = 0, f'(u_R) = \alpha$ (recall $f(u) = \frac{1}{2}u^2, f'(u) = u$).



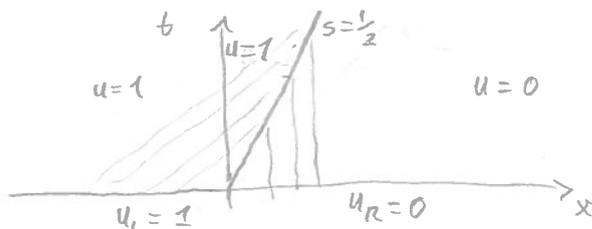
Case
 $\alpha = 1$

All good!

Example 2: Consider $u_t + uu_x = 0$ with $u_0 = \begin{cases} +1, & x < 0 \\ 0, & x > 0 \end{cases}$

To meet (LC), & RH-condition we solve this with shock surface of speed $s = \frac{1}{2}$, and constant states:

$$u(t, x) = \begin{cases} +1, & x < \frac{1}{2}t \\ 0, & x > 0 \end{cases}$$

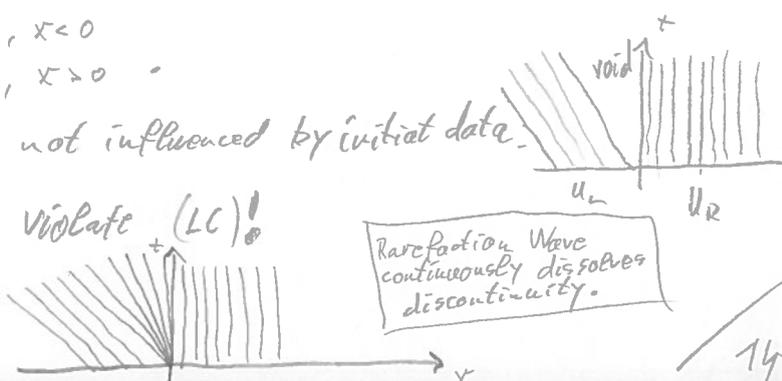


Consider $u_t + uu_x = 0$ with $u_0 = \begin{cases} 0, & x < 0 \\ +1, & x > 0 \end{cases}$.

This results in void region which is not influenced by initial data.

Any discontinuity in void region would violate (LC)!

↳ solve with "Rarefaction Wave":



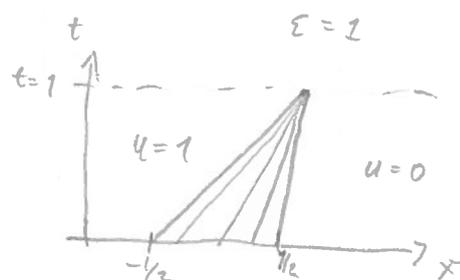
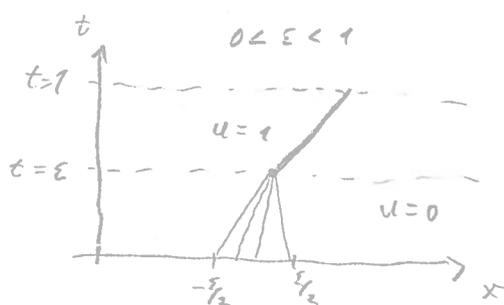
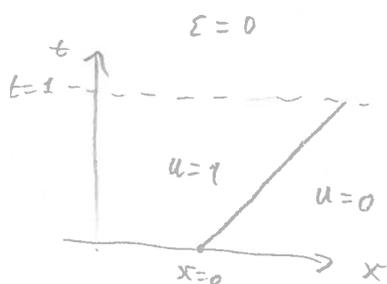
III.5) Irreversibility

Example:

For each $\varepsilon \in [0, 1]$, we define solution u_ε of $u_t + uu_x = 0$ by

$$\text{For } t \leq \varepsilon: \quad u_\varepsilon(t, x) := \begin{cases} 1, & x < t - \varepsilon/2 \\ \frac{x - \varepsilon/2}{t - \varepsilon}, & x \in (t - \varepsilon/2, \varepsilon/2) \\ 0, & x > \varepsilon/2 \end{cases}$$

$$\text{For } t > \varepsilon: \quad u_\varepsilon(t, x) := \begin{cases} 1, & x < \frac{t}{2} \\ 0, & x > \frac{t}{2} \end{cases}$$



Each solution u_ε satisfies Lax condition, but still, knowing solution at $t=1$ does not suffice to determine solution for $t < 1$.

↳ Irreversibility of discontinuous solutions!

Remark:

- Lax condition only gives unique solution "forward" in time, "backward" in time Lax condition is automatically violated which results in "backward in time non-uniqueness" (=irreversibility).
- Discontinuous solutions of conservation laws can be approximated by parabolic PDE's (Method of artificial viscosity),

$$u_t + f(u)_x = 0$$

$$\uparrow \varepsilon \gg 0$$

$$u_t + f(u)_x = \varepsilon B u_{xx}$$

and parabolic PDE's fail to be time-reversible.

IV) Riemann Problems

IV.1) Basic Notions

• Let $u(t,x), f(u) \in \mathbb{R}^N$.

Consider system of Conservation Laws

$$\boxed{u_t + f(u)_x = 0} \quad (CL)$$

• Def:

Let $u_L, u_R \in \mathbb{R}^N$ be constant, (so called "constant states").
Solving (CL) for u with initial data $u(0,x) = \begin{cases} u_L, & x < 0 \\ u_R, & x > 0 \end{cases}$

is called "Riemann Problem" for (CL) with data (u_L, u_R) .

• Observe, (CL) $\Leftrightarrow u_t + \underbrace{df|_u}_{=: A(u)} \cdot u_x = 0 \quad (CL')$
 $=: A(u), N \times N$ -matrix

Def: Let $M \subset \mathbb{R}^N$, open.

We say (CL) is "strictly hyperbolic" in M if there exists N mutually distinct eigenvalues $\lambda_1(u), \dots, \lambda_N(u)$ of $A(u) \forall u \in M$.
Then, we assume WLOG that $\lambda_1(u) < \dots < \lambda_N(u)$.

Remark:

If (CL) is strictly hyperbolic, then (CL) is symmetrizable.

• From now on we assume (CL) is strictly hyperbolic.

• Note, if $A(u) = \text{diag}(\lambda_1(u), \dots, \lambda_N(u))$, then (CL') were a system

of N scalar conservation laws $u_t^h + \lambda_h(u) u_x^h = 0, h \in \{1, \dots, N\}$.

Then each component u^h is constant along $t \mapsto \gamma_a^h(t)$, characteristic

line, defined by $\begin{cases} \dot{\gamma}_a^h(t) = \lambda_h(u(t, \gamma_a^h(t))) \\ \gamma_a^h(0) = a \end{cases}$

• This motivates for general case:

Def: We define the h -th "characteristic curve" of (CL), at $a \in \mathbb{R}$,
as the solution $\gamma_a^h(t)$ of ODE $\begin{cases} \dot{\gamma}_a^h(t) = \lambda_h(u(t, \gamma_a^h(t))) \quad [\in \mathbb{R}] \\ \gamma_a^h(0) = a \end{cases}$

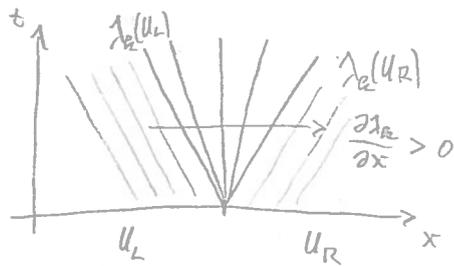
• Def

Let u be a C^1 -solution of (CL).

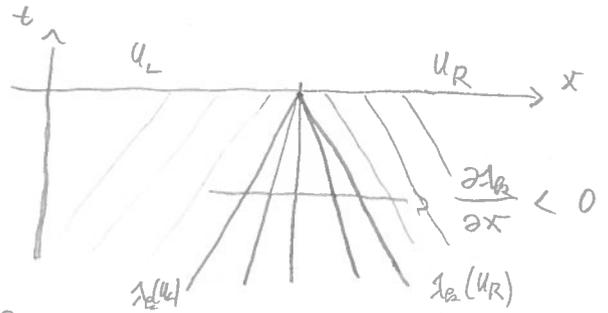
We call u a " k -simple wave" if u is constant along the k -th characteristic curve.

We call a k -simple wave u a " k -rarefaction wave", (" k -compression wave")

if $\frac{\partial}{\partial x} \lambda_k(u(t,x)) > 0$, (if $\frac{\partial}{\partial x} \lambda_k(u(t,x)) < 0$).



k -Rarefaction Wave



k -Compression Wave

Recall that $\lambda_k(u)$ gives the slope of the k -th characteristic curve. Thus, $\frac{\partial \lambda_k}{\partial x} > 0$ means this slope is increasing as one moves from left to right in above picture.

• For (CL) to admit k -rarefaction waves, (CL) must be genuinely nonlinear:

Def: Let $v_k(u)$ be eigenvector of $A(u)$ with $A v_k = \lambda_k v_k$, and $|v_k| = 1$.

We call (λ_k, v_k) the " k -th characteristic field" of (CL).

We call (λ_k, v_k) "genuinely nonlinear" in $M \subset \mathbb{R}^n$ if:

$$\langle v_k(u), \nabla \lambda_k(u) \rangle_{\mathbb{R}^n} > 0 \quad \forall u \in M.$$

We call (λ_k, v_k) "linearly degenerate" in M if

$$\langle v_k(u), \nabla \lambda_k(u) \rangle_{\mathbb{R}^n} = 0 \quad \forall u \in M.$$

Generalizes condition $f''(u) > 0$ for scalar systems.

- Recall that a weak solution of (CL), which is discontinuous across the line $\{x=st\}$, satisfies the Rankine-Hugoniot condition

$$\boxed{S[u] = [f(u)]} \quad (\text{RH})$$

Def: Assume the k -th characteristic field is linearly degenerate.

We call a (weak) solution of the Riemann problem (u_L, u_R) for (CL) a "contact discontinuity", if it is discontinuous across $\{x=st\}$ and C^1 otherwise, such that

$$\lambda_k(u_L) = S = \lambda_k(u_R).$$

Def: Assume the k -th char. field is genuinely nonlinear.

We call a (weak) solution of the Riemann problem (u_L, u_R) for (CL)

which is discontinuous across $\{(t, x=st) \mid t \geq 0\}$ and C^1 otherwise

an "(admissible) k -shock wave" if s meets "Lax condition":

$$\boxed{\lambda_k(u_L) > S > \lambda_k(u_R)} \quad (\text{LC})$$

Example:

The Euler equations (MA), (MO), (E) can be written in Lagrangian coordinates, replacing the internal energy e by the entropy S by 1st Law of Thermody.

$$ds \begin{cases} v_t - u_x = 0 \\ u_t + p_x = 0 \\ S_t = 0 \end{cases}, \quad \Leftrightarrow \quad \partial_t \begin{pmatrix} v \\ u \\ S \end{pmatrix} + \overbrace{\begin{bmatrix} 0 & -1 & 0 \\ p_v & 0 & p_s \\ 0 & 0 & 0 \end{bmatrix}}{=: A} \partial_x \begin{pmatrix} v \\ u \\ S \end{pmatrix} = 0,$$

where $v := \frac{1}{\rho}$ & $p = p(v, S)$, such that $p_v < 0$.

The eigenvalues of A are

$$\underline{\lambda_1 = -\sqrt{-p_v}}, \quad \underline{\lambda_2 = 0}, \quad \underline{\lambda_3 = \sqrt{-p_v}}.$$

\Rightarrow System is strictly hyperbolic.

& If $p_{vv} > 0$, then λ_1 & λ_3 are genuinely nonlinear,

while λ_2 is linearly degenerate.

• Thm ("Lax's Existence Theorem", 1957)

Assume (CL) is strictly hyperbolic and that each characteristic field is either linearly degenerate or genuinely nonlinear. Let $u_L \in \mathbb{R}^N$ be a constant state.

Then there exists a neighborhood $\mathcal{U} \subset \mathbb{R}^N$ of u_L such that for each $u_R \in \mathcal{U}$ the Riemann problem (u_L, u_R) for (CL) has a solution.

This solution consists of at most $(N+1)$ -constant states separated by admissible shock waves, rarefaction waves and contact discontinuities.

There is precisely one solution of this form.

• Idea of proof:

Omit contact discontinuities! (\rightarrow All char. fields are gen. nonlinear)

Step 1: Parameterize state space \mathcal{U} by shock & rarefaction curves!

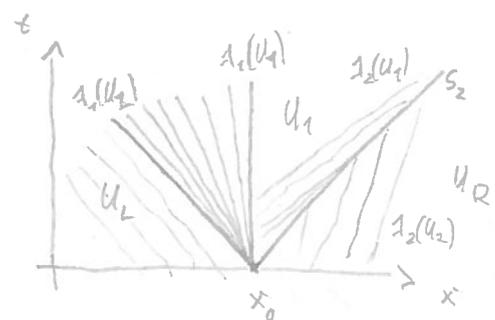
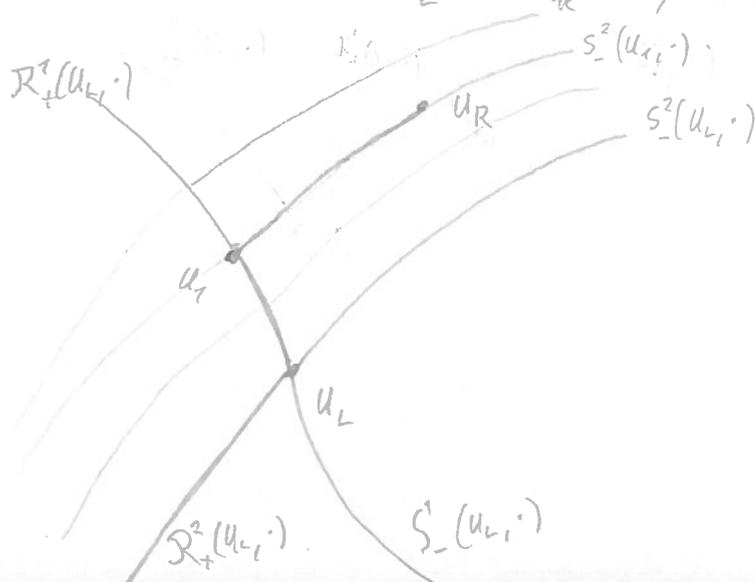
• Rarefaction curves $\mathcal{R}^1, \dots, \mathcal{R}^N: \frac{d\mathcal{R}^k}{ds} = r_k(\mathcal{R}^k)$, (\mathcal{R}^k k-th rarefaction curve)

$\hookrightarrow \mathcal{R}^k$ is flow of r_k , k-th eigenvector

• Shock curves S^1, \dots, S^N are implicit solutions of "Hugoniot locus"

$$u - u_L - s(f(u) - f(u_L)) = 0$$

Step 2: Find the unique intermediate state $u_1, \dots, u_{N-1} \in \mathcal{U}$ which connect u_L to u_R by the parameterization!



• Order curves so that wave speeds increase!

• Use shock curve for negative parameter and rarefaction curve for positive parameter