

Geometric quantization, fusion categories, and Rozansky–Witten theory

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Summary:

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- In many examples this projectivity can be understood as appearing via the homotopy theory of the higher automorphism group of the quantum theory itself.

$$\begin{array}{ccc} & & \pi_0 \text{Aut}(F) \\ & \nearrow & \uparrow \\ G & \dashrightarrow & \text{Aut}(F) \end{array}$$

Context and motivation

- Projective TQFTs in general are relevant to the classification of topological orders: a gapped quantum system is well-approximated at low energy by a *projective* field theory which is topological [Freed].

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 - Also see [Teleman] and [Braverman-Dhillon-Finkelberg-Raskin-Travkin-Johnson-Freyd].

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 - Also see [Teleman] and [Braverman-Dhillon-Finkelberg-Raskin-Travkin-Johnson-Freyd].
 - Relatedly, **RW** is the B-side of 3d mirror symmetry. See e.g. [Raskin-Hilburn, Gammage-Hilburn-Mazel-Gee].

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- 1 Quantum mechanics
- 2 Interlude: Families, symmetries, and anomalies
- 3 Fusion categories
- 4 Rozansky-Witten theory

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Deformation and geometric quantization

- Classical phase space: symplectic vector space (V, ω) , say over \mathbb{R} .

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where we have chosen a polarization $V \simeq \ell \oplus \ell^\vee$.

- The group $\mathrm{Sp}_{2n}(\mathbb{R})$ only acts *projectively* on \mathcal{H} . Equivalently, a central extension, classified by $w_2 \in H^2(B\mathrm{Sp}_{2n}(\mathbb{R}))$, acts linearly on \mathcal{H} :

$$g, h \mapsto T_g, T_h$$
$$T_h \circ T_g = c T_{gh}$$

$$\mathbb{Z}/2 \hookrightarrow \mathrm{Mp} \rightarrow \mathrm{Sp}_{2n}(\mathbb{R})$$

Table of analogies

$d = 1$ (QM)	(V, ω)	$\mathcal{O}(V), *_{\omega}$	$L^2(\ell)$	$w_2 \in H^2(B\mathrm{Sp}_{2n}(\mathbb{R}))$
$d = 3$ (TV)				
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Families

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$$\begin{array}{ccc} \mathbf{Bord}_d^X & & \\ \uparrow & \dashrightarrow^F & \\ \mathbf{Bord}_d & \xrightarrow{F_0} & \mathcal{T} \end{array}$$

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Two equivalent ways of encoding this notion:

$$\boxed{F: \mathbf{Bord}_d^X \rightarrow \mathcal{T} \quad F: 1 \rightarrow \sigma_X^{d+1}}$$

- Notation: The $(d+1)$ -dimensional gauge theory associated to X as in [\[Freed-Moore-Teleman\]](#) is $\sigma_X^{d+1}: \mathbf{Bord}_{d+1} \rightarrow \mathbf{Alg}(\mathcal{T})$.

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- The RHS is an (op)lax natural transformation as in [Johnson-Freyd-Scheimbauer].

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- On the other hand, the theory σ_{BG}^2 is a functor valued in the Morita category of algebras:

$$\mathbf{Bord}_2 \xrightarrow{\sigma_{BG}^2} \mathbf{Alg}$$

$$* \longmapsto (\mathbb{C}[G], *)$$

← 2-cobordism

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$$\begin{aligned} 1 &\rightarrow \sigma_{BG}^2 \\ * &: \mathbb{C} \rightarrow \mathbb{C}[G] \\ &\quad \mathfrak{M} \end{aligned}$$

- A theory defined relative to the theory σ_{BG}^2 is a morphism in the Morita category from \mathbb{C} to the group algebra, i.e. a module over the group algebra.

Anomalies

Let X be a (π -finite) groupoid (space), and now consider *projectivity data* on X :

$$\begin{array}{ccc} B^d \mathbb{C}^\times & \longrightarrow & \tilde{X} \\ & & \downarrow \\ & & X \xrightarrow{\alpha} B^{d+1} \mathbb{C}^\times \end{array}$$

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Projective/anomalous d -dimensional TQFT with background X -fields:

Theorem (VD23)

TFAE:

$F: 1_X \rightarrow \alpha$	$F: 1 \rightarrow \sigma_{X,c}$
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E.g. $X = BG$, $\tilde{X} = B\tilde{G}$ for \tilde{G} a central extension of G classified by α .

projective rep	mod over the twisted gp alg
rep of the ext'n	mod over gp alg of the ext'n

- We will focus on the analogue of a module over the twisted group algebra:

$$\sigma_{X,c}^{d+1}: \mathbf{Bord}_{d+1} \rightarrow \mathbf{Alg}(\mathcal{T}) \quad F: \mathbf{1} \rightarrow \sigma_{X,c}$$

- Caveat: For some of the examples we will consider, the theory $\sigma_{X,c}^{d+1}$ has not been formally constructed.
- In these cases, one can consider the analogue of a projective representation instead: $\mathbf{1} \rightarrow \alpha$ where everything is rigorous. See [\[VD23, Hypothesis Q\]](#) for more details.

Recasting quantization

- The theory of states defines a 1-dimensional QFT **GQ**.

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- The classical symmetry group $G = \mathrm{Sp}_{2n}(\mathbb{R})$ acting projectively on \mathcal{H} is equivalent to the theory living relative to twisted G -gauge theory:

$$\begin{array}{ccccc}
 \mathrm{U}(1) & \longrightarrow & \mathrm{U}(\mathcal{H}) & \longrightarrow & \mathrm{U}(\mathcal{H})/\mathrm{U}(1) \\
 \uparrow & & \uparrow & \nearrow & \uparrow \\
 \mathbb{Z}/2 & \longrightarrow & \mathrm{Mp} & \longrightarrow & \mathrm{Sp}
 \end{array}$$

$$\begin{array}{ccc}
 B\mathrm{Mp} & & \\
 \downarrow & & \\
 B\mathrm{Sp} & \xrightarrow{w_2} & B^2\mathbb{Z}/2
 \end{array}$$

$$g \longmapsto T_g \mathcal{B}\mathcal{H}$$

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$$\begin{array}{l}
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 c \in \mathcal{U}(1)
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$$U(1) \rightarrow U(\mathcal{H}) \rightarrow U(\mathcal{H})/U(1)$$

$$\begin{array}{ccccc} \uparrow & & \uparrow & \swarrow & \uparrow \\ \mathbb{Z}/2 & \longrightarrow & \mathrm{Mp} & \longrightarrow & \mathrm{Sp} \end{array}$$

$$\begin{array}{ccc} B\mathrm{Mp} & & \\ \downarrow & & \\ B\mathrm{Sp} & \xrightarrow{w_2} & B^2\mathbb{Z}/2 \end{array}$$

A commutative diagram illustrating the relationship between various sigma-2 maps and gauge theories. The diagram consists of several nodes and arrows:

- Top node: $\sigma_{U(\mathcal{H})}^2$
- Middle node: σ_{BG, w_2}^2
- Bottom node: σ_{BG}^2
- Left node: 1
- Right node: 1

Arrows and their directions:

- $\sigma_{U(\mathcal{H})}^2 \rightarrow \sigma_{BG, w_2}^2$ (solid black arrow)
- $\sigma_{BG, w_2}^2 \rightarrow \sigma_{BG}^2$ (solid black arrow)
- $\sigma_{BG}^2 \rightarrow 1$ (solid black arrow)
- $1 \rightarrow \sigma_{BG, w_2}^2$ (solid blue arrow)
- $1 \rightarrow \sigma_{BG}^2$ (dashed red arrow)
- $1 \rightarrow 1$ (curved arrow at the bottom labeled **GQ**)
- $\sigma_{U(\mathcal{H})}^2 \rightarrow 1$ (solid black arrow)
- $\sigma_{BG, w_2}^2 \rightarrow 1$ (solid black arrow)

Handwritten annotations:

- Blue text: $\mathbb{C}[\mathrm{Sp}]^{w_2}$ next to the blue arrow.
- Red text: $\mathbb{C}[\mathrm{Sp}]$ next to the red arrow.
- Black text: $\tilde{\mathrm{GQ}}$ next to the top-left arrow.

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$$\begin{array}{ccccc}
 U(1) & \longrightarrow & U(\mathcal{H}) & \longrightarrow & U(\mathcal{H})/U(1) \\
 \uparrow & & \uparrow & \swarrow \text{red dashed} & \uparrow \\
 \mathbb{Z}/2 & \longrightarrow & Mp & \longrightarrow & Sp
 \end{array}$$

$$\begin{array}{ccc}
 B Mp & & \\
 \downarrow & & \\
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 \end{array}$$

" $F \hookrightarrow \text{Amb}(F)$ "

$$\begin{array}{ccc}
 \underbrace{F}_{\mathbb{F}} & \xrightarrow{\sigma_{U(\mathcal{H})}^2} & \\
 \uparrow \text{blue} & \downarrow & \\
 \tilde{GQ} & \xrightarrow{\sigma_{BG, w_2}^2} & 1 \\
 \uparrow \text{red dashed} & \uparrow & \\
 1 & \xrightarrow{\sigma_{BG}^2} & 1 \\
 \text{arc } GQ & &
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Fusion categorical analogue of deformation quantization

- Analogue of classical phase space: finite abelian group Λ equipped with a quadratic form $q: \Lambda \rightarrow \mathbb{C}^\times$.

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- Analogue of deformation quantization: braided fusion category

$$(\mathbf{Vect}[\Lambda], *, \beta_q) \quad \beta_{\mathbb{C}_\ell, \mathbb{C}_k}: \mathbb{C}_\ell * \mathbb{C}_k \xrightarrow{\langle \ell, k \rangle_q \text{id}} \mathbb{C}_k * \mathbb{C}_\ell .$$

$$\langle -, - \rangle_q = \frac{q(-+ -)}{q(-)q(-)}$$

$$\mathbb{C}_{\ell+k}$$

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- Retains an action of $O(\Lambda)$.

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$$\mathbb{Z}(\text{Vect}[L]) \simeq \mathcal{A} \quad \textcircled{\mathcal{G}} \quad (\mathbf{Vect}[L], *) = \mathcal{E} \quad \mathcal{TV}: \begin{matrix} \psi \mapsto \mathcal{E} \\ \psi' \mapsto \mathcal{A} \end{matrix}$$

where we have chosen $\Lambda \simeq L \oplus L^\vee$ such that $q = \text{ev}$.

3 type $B^3C^2 \hookrightarrow \text{BAut}_{\text{int}}(\text{Vect}[L-1])$

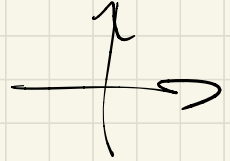
2 type $\text{BAut}_{\text{EgPer}}(\mathbb{Z}(e)) \simeq \Pi_{\leq 2} \quad \downarrow \Sigma$

Per Eg cl of $\mathbb{Z}(e)$

deformations

Moduli class of $e = \text{Vect}[L]$

has only up to 2 deformations



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Vect [L]

where we have chosen $\Lambda \simeq L \oplus L^\vee$ such that $q = \text{ev}$.

- The group $O(\Lambda)$ only acts projectively on **Vect** [L].

The fusion categorical anomaly

Recasting the previous slide in terms of TQFTs:

- There is a (framed, fully-extended) Turaev-Viro theory, which sends the point to **Vect** $[L]$, in the Morita 3-category of fusion categories [Douglas-Schommer-Pries-Snyder].

The fusion categorical anomaly

Recasting the previous slide in terms of TQFTs:

- There is a (framed, fully-extended) Turaev-Viro theory, which sends the point to $\mathbf{Vect}[L]$, in the Morita 3-category of fusion categories [Douglas-Schommer-Pries-Snyder].
- Using the obstruction theory of [Etingof-Nikschyck-Ostrik]:

Theorem (VD23)

The Turaev-Viro theory for $\mathbf{Vect}[L]$ can be upgraded to a theory defined relative to a twisted gauge theory for $O(\Lambda)$:

$$\mathbf{TV}: 1 \rightarrow \sigma_{BO(\Lambda),c}^4 = \mathcal{O}_4$$

Table of analogies (reprise)

$d = 1$ (QM)	(V, ω)	$\mathcal{O}(V), *_{\omega}$	$L^2(\ell)$	$w_2 \in H^2(B\mathrm{Sp}_{2n}(\mathbb{R}))$
$d = 3$ (TV)	(Λ, q)	$\mathbf{Vect}[\Lambda], \beta_q$	$\mathbf{Vect}[L]$	$O_4 \in H^4(B\mathrm{O}(\Lambda))$
$d = 3$ (RW)				

3-yp with
inv. brn/walks

$$\mathrm{Aut}_{\mathrm{Fns}}(\mathrm{Vect}[L]) = \mathrm{BnPic}(\mathrm{Vect}[L])$$

$$\downarrow \quad \quad \quad \mathrm{Pic}(\mathbb{Z}(\tau))$$

$$G = \mathcal{O}(L \otimes L^{\vee}, \omega)$$

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 B^2\mathbb{C}^{\times} & \longrightarrow & \mathrm{Aut}(\mathbf{TV}) & \longrightarrow & \pi_{\leq 1} \mathrm{Aut}(\mathbf{TV}) \\
 \parallel & & \uparrow & \swarrow \text{---} & \uparrow \\
 B^2\mathbb{C}^{\times} & \longrightarrow & \tilde{O} & \longrightarrow & O(\Lambda)
 \end{array}$$

$U(n)/U(1)$
 \uparrow
 Sp

$$\begin{array}{ccc}
 B\tilde{O} & & \\
 \downarrow & & \\
 B\mathrm{O}(V) & \xrightarrow{O_4} & B^4\mathbb{C}^{\times}
 \end{array}$$

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 \parallel & & \uparrow & \swarrow \text{dashed red} & \uparrow \\
 B^2\mathbb{C}^{\times} & \longrightarrow & \tilde{O} & \longrightarrow & O(\Lambda) \\
 & & \downarrow B\tilde{O} & & \\
 & & BO(V) & \xrightarrow{O_4} & B^4\mathbb{C}^{\times}
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$$\mathcal{L} \rightsquigarrow \mathcal{M}$$

- The classical phase space: 2-shifted symplectic stack (M, ω) as in [Calaque, Pantev, Toën, Vaquié, Vezzosi].

Shifted deformation quantization

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- The deformation quantization can be thought of as a braided deformation of the derived category $\mathbf{QC}(M)$ of quasi-coherent sheaves.

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- The deformation quantization can be thought of as a braided deformation of the derived category $\mathbf{QC}(M)$ of quasi-coherent sheaves.
 - Think: the deformation quantization of k -shifted is an \mathbb{E}_{k+1} -algebra, modules over this form an \mathbb{E}_k -category.

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- The theory should fit into the framework of the AKSZ construction. [Alexandrov-Kontsevich-Schwarz-Zaboronsky, Qiu-Zabzine, Scheimbauer-Calaque-Haugseeng, Stefanich, Riva]

Table of analogies (reprise²)

$d = 1$ (QM)	(V, ω)	$\mathcal{O}(V), *_{\omega}$	$L^2(\ell)$	$w_2 \in H^2(B\mathrm{Sp}_{2n}(\mathbb{R}))$
$d = 3$ (TV)	(Λ, q)	$\mathbf{Vect}[\Lambda], \beta_q$	$\mathbf{Vect}[L]$	$O_4 \in H^4(B\mathrm{O}(\Lambda))$
$d = 3$ (RW)	(M, ω)	$\mathbf{QC}(M), \beta_{\omega}$	$\mathbf{RW}(*)$	$? \in H^4(B\mathrm{Sp}(M))$

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 \parallel & & \uparrow & \swarrow & \uparrow \\
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F	Class	obs	$F(\cdot)$	egw.
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"Shw(\sqrt{m})"

$L \subset \mathbb{T}^+ \sqrt{m}$

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 & \sigma_B^4 \mathrm{Aut}(\mathbf{RW}) & \\
 \widetilde{\mathrm{RW}} \nearrow & \downarrow & \searrow \\
 1 & \sigma_{\widetilde{\mathrm{Sp}}}^4 & 1 \\
 \uparrow & \uparrow & \\
 1 & \sigma_{\mathrm{Sp}}^4 & 1 \\
 \text{---} \swarrow & & \searrow \text{---} \\
 & \mathrm{RW} &
 \end{array}$$

(V, γ)	$cl(N, \gamma)$	$M \simeq \Lambda \times$	$H^2(BSO)$
(Λ, γ)	$Urb[\Lambda]$	$Urb[\mathbb{C}]$	$H^4(BSO)$

$$\Lambda = \text{U.S.} / F_p \quad \gamma \in H^4(B^2\Lambda, \mathbb{C}^*)$$

$$\simeq H^2(B^2\Lambda, B^2\mathbb{C}^*)$$

$$H_{cb}^3(\Lambda, \mathbb{C}^*)$$

My current work in this direction

- The analogue of Crane-Yetter in the **RW** context.
 - One strategy for accessing this projective Sp-action.
- General relationship between the prequantum k -gerbe and the anomaly $(k + 1)$ -gerbe.
 - One strategy for understanding the impact of changing the polarization on the shifted geometric quantization.

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