Quantum many-body dynamics in two dimensions with artificial neural networks

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QM³ seminar Lisbon

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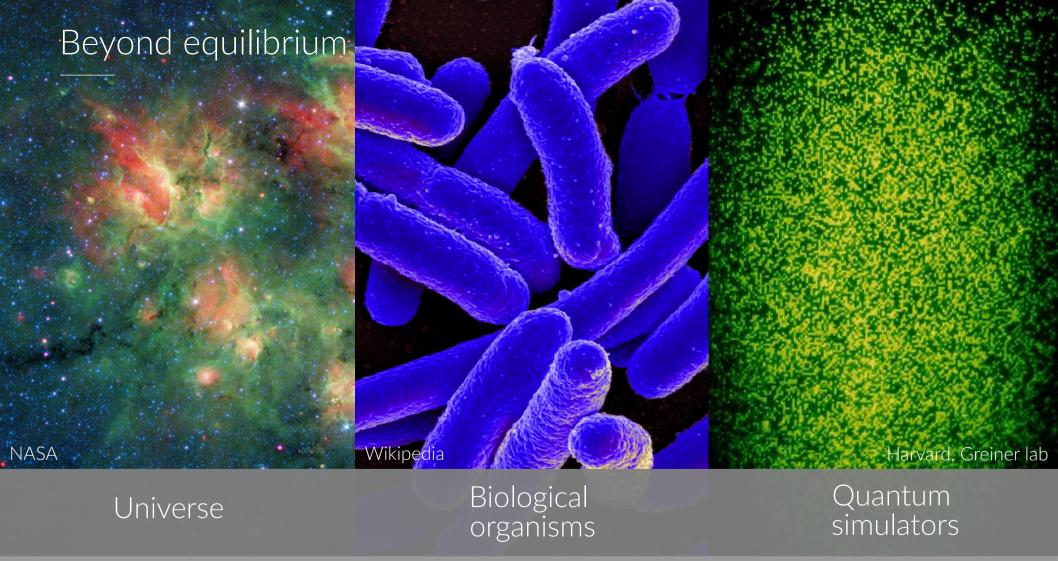








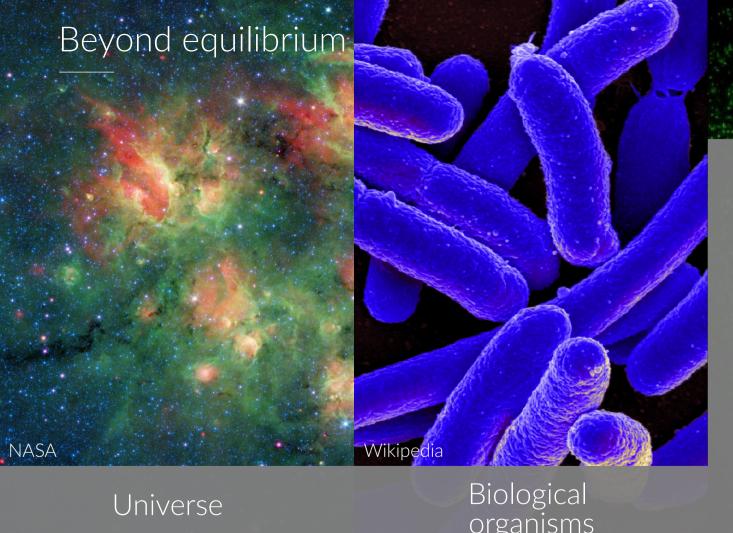




Markus Hey

Machine learning quantum dynamics

2



Coherent dynamics of quantum matter

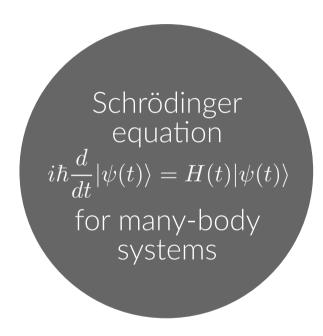
$$i\hbar \frac{d}{dt}|\psi(t)\rangle = H(t)|\psi(t)\rangle$$

Unprecedented control at the quantum level

Quantum simulators

organisms

Dynamics in correlated quantum matter



Outline

- Quantum dynamics in 2D: motivation and challenges
- Classical networks and artificial neural network wave functions
- The inversion problem: how to stabilize dynamics with neural networks
- Outlook

Quantum dynamics in 2D: motivation and challenges

Quantum Dynamics in 2D

Today: key turning point in theory and experiment

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Quantum simulators: Coherent dynamics of quantum matter in **2d**Harvard, MPQ, Princeton, Paris,...

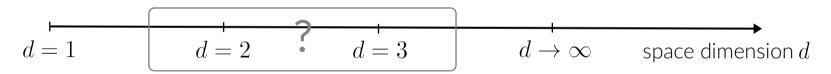
Theory: severe limitation for (numerically) exact methods

Quantum Dynamics in 2D

Today: key turning point in theory and experiment

Quantum simulators: Coherent dynamics of quantum matter in 2d Harvard, MPQ, Princeton, Paris,...

Theory: severe limitation for (numerically) exact methods



Tensor network methods

Limited by entanglement

Review: Schollwöck '11

Dynamical mean-field theory
Limited by spatial resolution
Review: Aoki '13

Dynamics in correlated quantum matter

Challenge: Hilbert space exponential

Quantum many-body state:
$$|\psi\rangle=\sum_s\psi(s)|s\rangle$$
 # amplitudes ~ 2^N (N spin-1/2's, fermions,...)

Exponential computational resources required (without efficient compression)

Dynamics in correlated quantum matter

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Our approach

"Don't store. Generate on the fly."

$$|s\rangle$$
 — Machine $\tilde{\psi}(s) \approx \psi(s)$

sample with Monte Carleo (MC)

Classical networks and artificial neural network wave functions

Classical networks

 $|s\rangle \longrightarrow \text{Classical network} \longrightarrow \tilde{\psi}(s,t) = e^{\mathcal{H}_{\mathrm{eff}}(s,t)}$

Effective classical Hamiltonian

- Structure obtained from cumulant expansion (around a classical limit)
- Perturbatively controlled
- Further variationally optimized

$$\mathcal{H}_{\mathrm{eff}}(s,t) = h_0(s,t) + \epsilon h_1(s,t) + \epsilon^2 h_2(s,t) + \dots$$
 ϵ : small parameter

(similar to Jastrow or Huse & Elser wave functions)

Classical networks

Classical network Schmitt & MH SciPost '18

Effective classical Hamiltonian

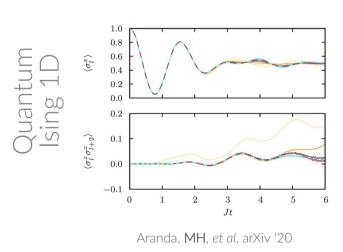
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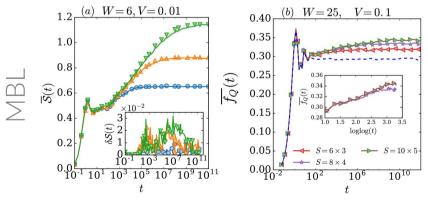
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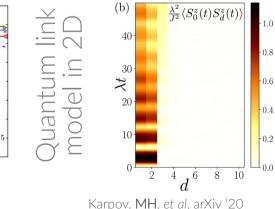
(similar to Jastrow or Huse & Elser wave functions)

- Perturbatively controlled
- Further variationally optimized

Successful for tailored problems

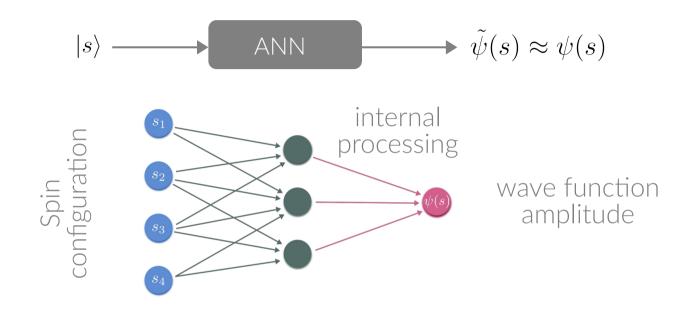






De Tomasi, Pollmann, & MH PRB (R) '19

Encoding quantum states in artificial neural networks (ANN)

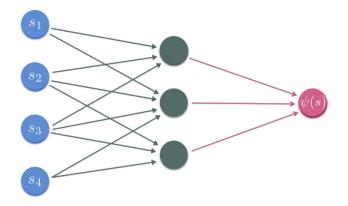


Enabling idea: ANNs universal function approximators

→ Any quantum state can be encoded in a (sufficiently large) ANN

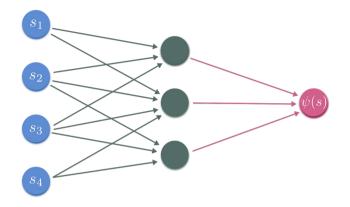
Carleo & Troyer Science '17

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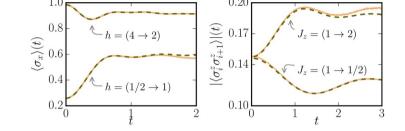
ANN is not just a black box \rightarrow Numerically exact approach

Convergence parameter: size of ANN

Complexity of algorithm: poly(size of ANN, system size)

Challenging for nonequilibrium dynamics

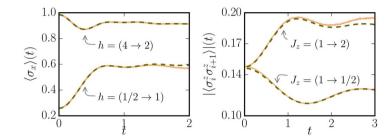
Results in the literature: limitations (either in system size, time,...)



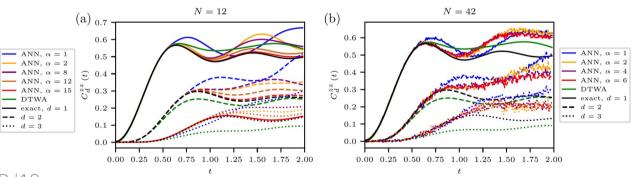
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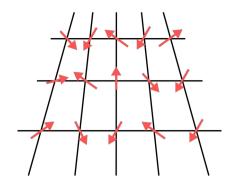
Systematic discrepancy even for increasing size of net

Czischek et al PRB '18

Overcoming the challenges

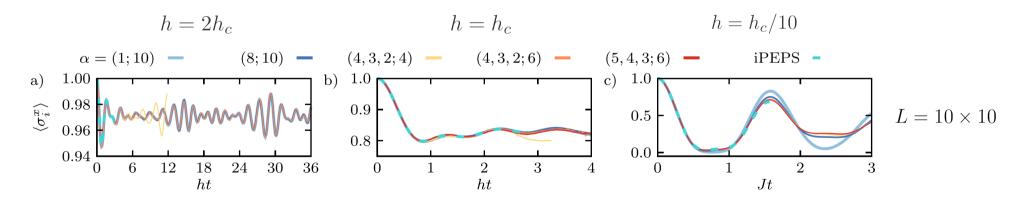
2D transverse-field Ising model

$$H = -J\sum_{\langle i,j\rangle} \sigma_i^z \sigma_j^z - h\sum_j \sigma_j^x$$



Nonequilibrium quantum quench: $|\psi_0\rangle = |\to\rangle \implies |\psi_0(t)\rangle = e^{-iHt}|\psi_0\rangle$

M. Schmitt & MH PRL '20



iPEPS data from Czarnik et al. PRB '19

The inversion problem: how to stabilize dynamics with neural networks 21

The setup

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Approximate time-evolved state by variational wave function

$$|\psi(t)\rangle \approx |\psi_{\eta(t)}\rangle = \sum_{s} \psi_{\eta(t)}(s)|s\rangle$$

Goal: choose variational parameters optimally. How?

Do (infinitesimal) step

$$|\psi_{\eta(t)}\rangle \mapsto e^{-iH\Delta t}|\psi_{\eta(t)}\rangle$$

Is now out of the variational subspace

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Project back and identify a new optimal set of variational parameters

$$F_{\eta'} = |\langle \psi_{\eta'} | e^{-iH\Delta t} | \psi_{\eta(t)} \rangle| \implies \sup_{\eta'} F_{\eta'}$$

For an infinitesimal time step: Taylor expansion

$$\eta' = \eta + \Delta \eta, \quad e^{-iH\Delta t} = 1 - iH\Delta t - \frac{1}{2}H^2\Delta t + \dots$$

Quantum dynamics o Nonlinear classical ODE for network weights η_k

$$S_{k,k'} = \langle \langle O_k^* O_{k'} \rangle \rangle_c \longrightarrow S_{k,k'} \dot{\eta}_{k'} = F_k \longrightarrow F_k = -i \langle \langle O_k^* E_{\text{loc}} \rangle \rangle_c$$

$$O_k(s) = \frac{\partial \ln \psi_{\eta}(s)}{\partial \eta_k} \qquad E_{\text{loc}}(s) = \sum_{s'} \langle s | H | s' \rangle \frac{\psi_{\eta}(s')}{\psi_{\eta}(s)}$$

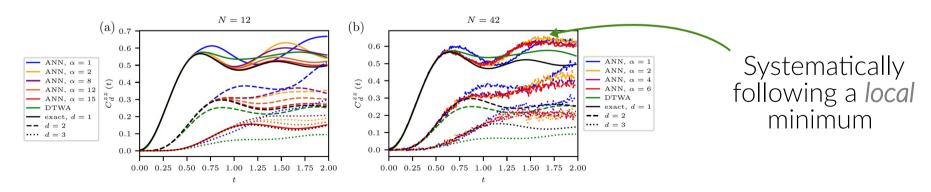
When solved exactly: prescription to follow the minimum in variational manifold step by step

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Key solution: properly account for noisy estimates

$$S\dot{\eta} = F \quad \Rightarrow \quad \dot{\eta} = S^{-1}F$$

M. Schmitt & MH PRL '20

Problem: S not invertible...

→ Pseudo-inverse: threshold on singular values

But: We do MC and there is noise everywhere (also in the singular values)

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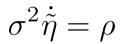
But: We do MC and there is noise everywhere (also in the singular values)

Solution: Represent TDVP in the diagonal basis of the S matrix

$$S_{k,k'}\dot{\eta}_{k'} = F_k \qquad \qquad \sigma_k^2\dot{\tilde{\eta}}_k = \langle\langle Q_k^*E_{loc}\rangle\rangle_c \equiv \rho_k$$

$$S_{k,k'} = V_{k,l}\sigma_l^2(V^\dagger)_{l,k'}, \ Q_k = (V^\dagger)_{k,k'}O_k,$$

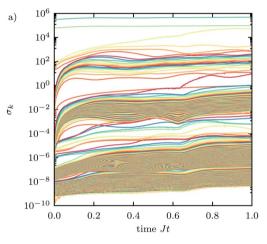
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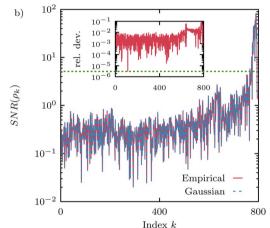


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Noise independent of signal strength

$$SNR(\sigma_k) = \sqrt{N_{MC}/2}$$





Signal to noise ratio can vary over orders of magnitude

$$SNR(\rho_k) = \sqrt{\frac{N_{MC}}{1 + \frac{\sigma_k^2}{\rho_k^2} Var(H)}}$$

New regularization scheme for inversion: SNR threshold

Quantum quenches in the 2D transverse-field Ising model

$$|\psi_0\rangle = |\to\rangle \implies |\psi_0(t)\rangle = e^{-iHt}|\psi_0\rangle, \qquad H = -J\sum_{\langle i,j\rangle} \sigma_i^z \sigma_j^z - h\sum_j \sigma_j^x$$

$$h = 2h_c \qquad h = h_c \qquad h = h_c/10$$

$$\alpha = (1;10) \qquad (8;10) \qquad (4,3,2;4) \qquad (4,3,2;6) \qquad (5,4,3;6) \qquad \text{iPEPS} \qquad (0.98)$$

$$0.98 \qquad 0.98 \qquad 0.99 \qquad 0.94 \qquad 0.94$$

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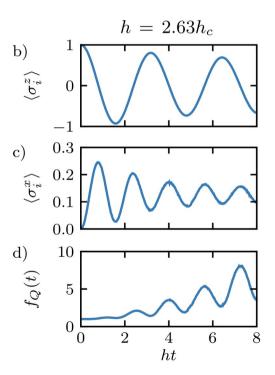
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iPEPS data from Czarnik et al. PRB '19

Collapse and revival oscillations in the 2D Ising model

$$|\psi_0\rangle = |\uparrow\rangle \implies |\psi_0(t)\rangle = e^{-iHt}|\psi_0\rangle$$



Order parameter

→ decay and revival of long-range order

Transverse magnetization

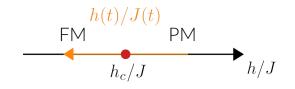
→ Buildup of thermal magnetization profile

Quantum Fisher information density

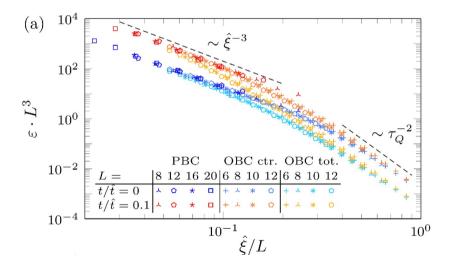
→ significant multipartite entanglement

$$f_Q(t) = \frac{1}{N} \sum_{i,j} \langle \sigma_i^z \sigma_j^z \rangle_c$$

$$H = -J \sum_{\langle i,j \rangle} \sigma_i^z \sigma_j^z - h \sum_j \sigma_j^x$$



Kibble-Zurek mechanism in 2D transverse-field Ising model (together with M. Schmitt, M. Rams, J. Dziarmaga, W. Zurek)

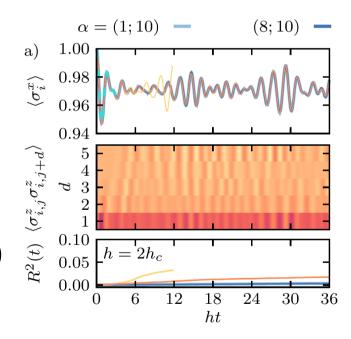


→ room for exploring yet inaccessible regimes with ANNs

Summary

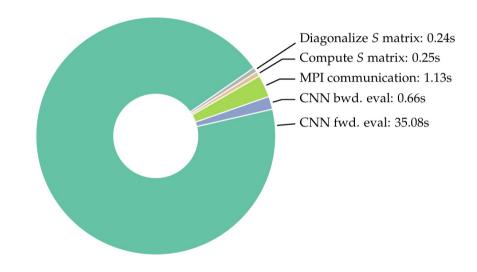
Powerful tool to describe nonequilibrium dynamics in 2D

- Competitive with (or even superior to) tensor networks
- Increasing system size or time at moderate polynomial costs
- Current limitations: still instabilities
- Not (so much) limited by network size (expressivity)



Computational resources

Distribution of compute time for one time step on 40 NVIDIA Tesla V100 GPUs



$$N_{MC} = 5 \cdot 10^5$$
 $N = 100$

Overall complexity: $\mathcal{O}(N_{MC} \times \max(N^2, P) \times P)$