

# On the Fukaya-Morse $A_\infty$ category

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Joint work with O. Chekeres, A. Losev and D. Youmans,  
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# Plan

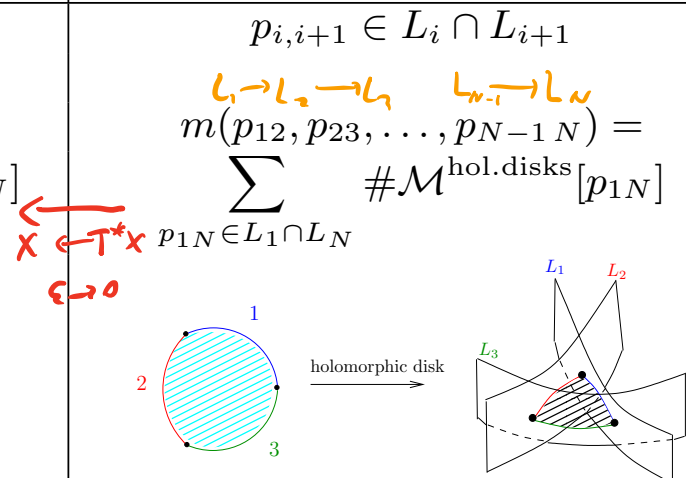
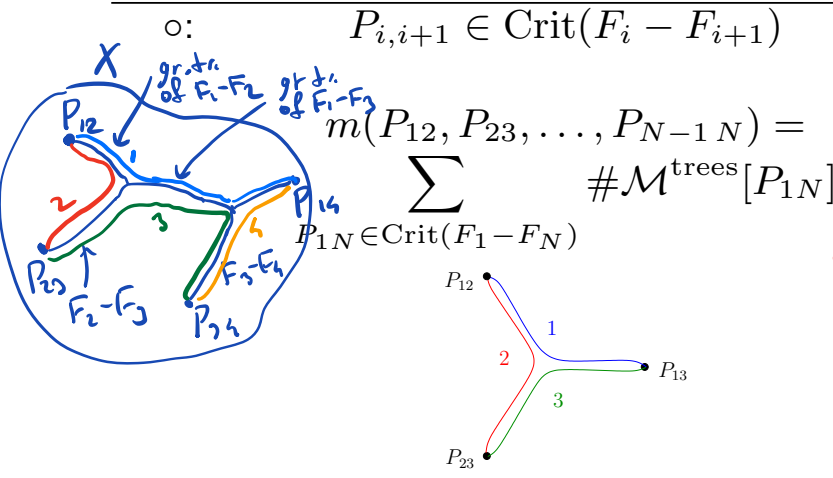
- 1 Fukaya-Morse  $A_\infty$  category (+ “enhancement”)
- 2 Picture 1: effective field theory (second quantization) **X**
- 3 Picture 2: HTQM (first quantization approach)

# Fukaya $A_\infty$ category and Morse degeneration

Fix  $X$  – a Riemannian manifold

Fukaya '93  
 Fukaya-Oh '97  
 Kontsevich-Soibelman '00

Fukaya-Morse $A_\infty$ category on $X$	Fukaya $A_\infty$ category on $T^*X$
Ob: $F_1, \dots, F_N$	Lagrangians $L_a = \text{graph}(\epsilon dF_a)$ $\xrightarrow{\epsilon \rightarrow 0}$
Mor: $\text{Mor}(F_a, F_b) = \text{Span}_{\mathbb{Z}}(\text{Crit}(F_a - F_b))$	$\text{Mor}(L_a, L_b) = \text{Span}_{\mathbb{Z}}(\text{intersection points of } L_a, L_b)$ $d(F_a - F_b) = 0$



# Enhancement

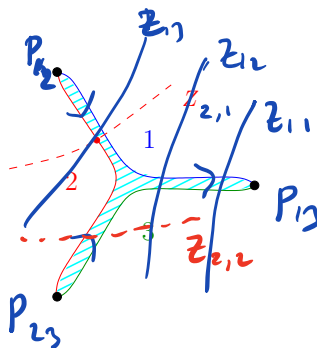
**Enhancement:**  $\text{Mor}(F_a, F_a) = C_{\bullet}^{\text{sing}}(X)$

Composition maps:

$$F_1 \xrightarrow{z_{1,1}} F_1 \xrightarrow{z_{1,2}} \dots \xrightarrow{z_{1,n}} F_1 \xrightarrow{P_{12}} F_2 \xrightarrow{z_{2,1}} F_2 \xrightarrow{z_{2,2}} \dots \xrightarrow{z_{2,n}} F_2 \xrightarrow{P_{2,3}} F_3 \dots \xrightarrow{P_{N-1,N}} F_{N-1} \xrightarrow{P_{N-1,N}} F_N \xrightarrow{\dots} F_N$$

$$m: \text{Mor}(F_1, F_1)^{\otimes k_1} \otimes \text{Mor}(F_1, F_2) \otimes \dots \otimes \text{Mor}(F_{N-1}, F_N) \otimes \text{Mor}(F_N, F_N)^{\otimes k_N} \rightarrow \text{Mor}(F_1, F_N)$$

$$m(\{Z_{1,\alpha}\}, P_{12}, \{Z_{2,\alpha}\}, \dots, P_{N-1,N}, \{Z_N,\alpha\}) = \sum_{P_{1N}} \#\mathcal{M} \cdot [P_{1N}]$$



- differentials:  $m(Z) = \partial Z$ ,  $m(P_{12}) = d_{\text{Morse}} P_{12}$
- $m(Z_1, Z_2) = Z_1 \cap Z_2$
- $m(Z_1, \dots, Z_n) = 0$ ,  $n \geq 3$

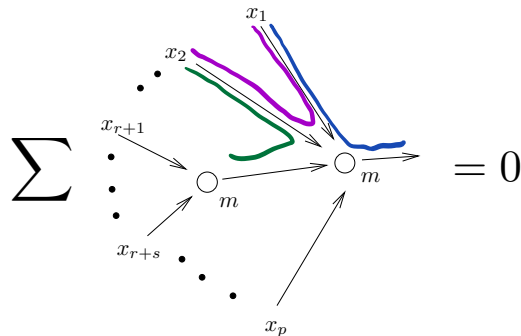
# $A_\infty$ relations

$$F_{a_1} \xrightarrow{x_1} F_{a_2} \xrightarrow{x_2} F_{a_3} \xrightarrow{x_3} \dots \xrightarrow{x_p} F_{a_{p+1}}$$

Let  $x_i \in \text{Mor}(F_{a_i}, F_{b_i})$ ,  $i = 1 \dots p$ , with  $b_i = a_{i+1}$  - composable sequence of morphisms. Then:

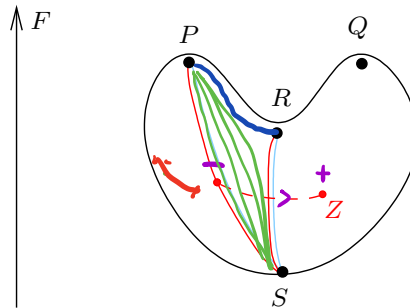
$$\sum_{r,s} m(x_1, \dots, x_r, m(x_{r+1}, \dots, x_{r+s}), x_{r+s+1}, \dots, x_p) = 0$$

Or:



# Example of an $A_\infty$ relation: heart-shaped sphere

**Example:**  $X = S^2$ ,  $N = 2$  functions,  $F = F_1 - F_2$ .



$$\begin{aligned}
 & F_1 \xrightarrow{P} F_2 \xrightarrow{Z} F_2 \\
 & \underbrace{m(m(P), Z)}_R \\
 & + \underbrace{m(P, m(Z))}_{-S} \\
 & + \underbrace{m(m(P, Z))}_0 \\
 & = 0
 \end{aligned}$$

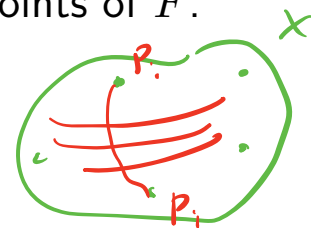
$A_\infty$  relation:  $x_1 = P \in \text{Mor}(F_1, F_2)$ ,  $x_2 = Z \in \text{Mor}(F_2, F_2)$ .

$$\boxed{m(d_{\text{Morse}}(P), Z) + m(P, \partial Z) + d_{\text{Morse}}m(P, Z) = 0}$$

# Example: deformation of Morse differential by cycles

$N = 2, F = F_1 - F_2$ . Fix  $\{C_\alpha\}$  -cycles on  $X$ .  $\{P_i\}$  - crit. points of  $F$ .  
 Generating function for compositions

$$m: \text{Mor}(F_1, F_2) \otimes \text{Mor}(F_2, F_2)^{\otimes k} \rightarrow \text{Mor}(F_1, F_2),$$



$$m_i^j(T) = \sum_{k \geq 1} \sum_{\alpha_1, \dots, \alpha_k} \underbrace{\# \mathcal{M}(P_i, C_{\alpha_1}, \dots, C_{\alpha_k}, P_j)}_{\# \text{grad traj } P_i \rightarrow P_j \text{ passing through cycles}} T_{\alpha_1} \cdots T_{\alpha_k}$$

$T_\alpha$  - generating parameters,  $|T_\alpha| = 1 - \text{codim} C_\alpha$ .

$$A_\infty \text{ relations} \Rightarrow \boxed{(d_{\text{Morse}} + m(T))^2 = 0}.$$

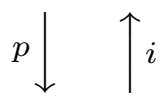
$$(d_{\text{Morse}} + m(T)) \circ MC(F) = 0$$

Explanation from HPT:

$$G = L_v + \varepsilon d^*$$

$$H = \underline{\underline{L_v + \varepsilon \Delta}}$$

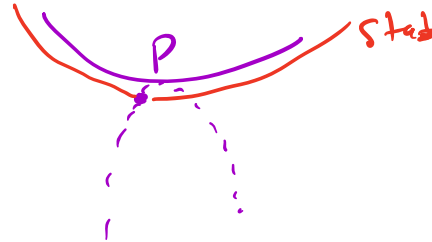
$$K \hookrightarrow \Omega^\bullet(X), d + \sum_\alpha T_\alpha \delta_{C_\alpha} \wedge$$



$$MC(X, F), d_{\text{Morse}} + m(T)$$

$$p d_{2i} - p d_2 K d_2 i - p d_1 K d_2 K d_2 i - \dots$$

# Morse contraction



- $i: \underset{\text{crit. point}}{P} \mapsto \delta_{\text{Unstab}_P}$

- $p: \omega \mapsto \sum_P \left( \int_X \omega \wedge \delta_{\text{Stab}_P} \right) \cdot [P]$

- $K = \int_0^\infty dt \iota_v e^{-t\mathcal{L}_v} : \Omega^\bullet(X) \rightarrow \Omega^{\bullet-1}(X).$

$v$  – gradient vector field.

Integral kernel:  $\delta_Y \in \Omega_{\text{distr}}(X \times X);$

$Y = \{(x, y) \mid x = \text{Flow}_t(v) \circ y \text{ for some } t > 0\}$





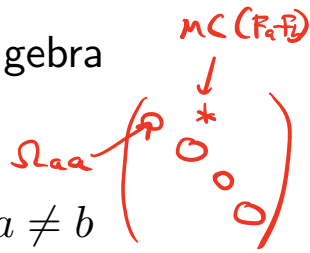
# Picture 1a: homotopy transfer

$$\underline{K} \hookrightarrow V = \Omega^\bullet(X) \otimes \text{Mat}_{N \times N} = \bigoplus_{a,b=1}^N \Omega_{ab}^\bullet(X) \quad \text{-- dg algebra}$$

$$\begin{array}{ccc} & & \\ & \downarrow \underline{p} & \uparrow \underline{i} \\ & & \end{array}$$

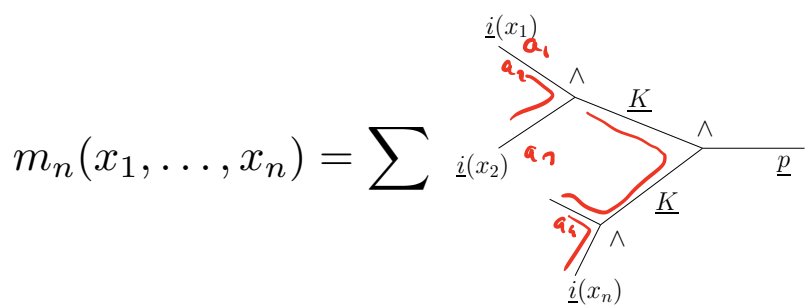
$$\mathbb{M} = \bigoplus_{a,b} \mathbb{M}_{ab}$$

$$\mathbb{M}_{ab} = \begin{cases} MC(F_a - F_b), & a \neq b \\ \Omega_{aa}^\bullet, & a = b \end{cases}$$



$$\underline{i}, \underline{p}, \underline{K} = \begin{cases} \text{Morse contraction for } F_a - F_b, & a \neq b \\ \text{trivial } (\underline{i} = \underline{p} = \text{id}, \underline{K} = 0), & a = b \end{cases}$$

**Induced  $A_\infty$  algebra structure on  $\mathbb{M}$ :**



Kadeishvili  
 - Kontsevich-Solterman

$$\text{inputs: } x_i \in \begin{array}{l} \mathbb{M}_{a_i b_i} \\ \text{with } b_i = a_{i+1} \end{array} = \begin{cases} \text{Morse chain,} & a_i \neq b_i \\ \text{form/sing. chain } \delta_Z, & a_i = b_i \end{cases}$$

# Picture 1b: effective action

$$\text{BF theory: } S = \int_X \langle B \wedge dA + \frac{1}{2}[A, A] \rangle$$

**Fields:**

$$\bullet A \in \Omega^\bullet(X) \otimes \underbrace{(\text{Mat}_{N \times N} \otimes \mathbb{A})}_{\mathfrak{g}}[1] = \bigoplus_{a,b} \Omega_{ab} \otimes \mathbb{A}[1]$$

where  $\mathbb{A} =$  upper-triangular  $\tilde{N} \times \tilde{N}$  matrices.

$$\bullet B \in \Omega^\bullet(X) \otimes \mathfrak{g}^*[d-2]$$

**Next:** integrate out off-diagonal components  $A_{ab}$ ,  $a \neq b$  subject to gauge-fixing

$$\iota_{v_{ab}} A_{ab} = 0$$

– axial gauge but in different directions for different components!  
We induce the effective action on diagonal fields + remnants of off-diagonal fields.

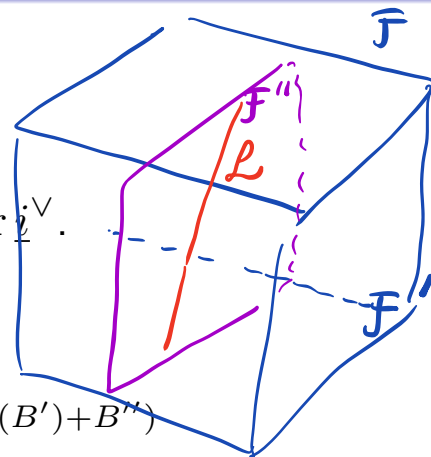
# Picture 1b: effective action cont'd

**More explicitly:**

BV pushforward:  $\mathcal{F}_{\text{fields}} \rightarrow \mathcal{F}'_{\text{IR fields}} = T^*[-1](\mathbb{M} \otimes \mathbb{A}[1])$

Splitting of fields:  $\mathcal{F} = \mathcal{F}' \oplus \mathcal{F}''$  with  $\mathcal{F}'' = \ker \underline{p} \oplus \ker \underline{i}^\vee$ .

Gauge-fixing Lagrangian:  $\mathcal{L} = \text{im}(K) \oplus \text{im}(K^\vee) \subset \mathcal{F}''$



$$e^{\frac{i}{\hbar} S_{\text{eff}}(A', B')} = \int_{\mathcal{L}} \mathcal{D}A'' \mathcal{D}B'' e^{\frac{i}{\hbar} S(\underline{i}(A') + A'', \underline{p}^\vee(B') + B'')}$$

$$S_{\text{eff}} = \sum_{\text{trees}} \begin{array}{c} \underline{i}(A') \\ \diagdown \\ \text{[ ]} \\ \diagup \\ \underline{i}(A') \end{array} \begin{array}{c} \underline{K} \\ \diagdown \\ \text{[ ]} \\ \diagup \\ \underline{K} \\ \diagdown \\ \text{[ ]} \\ \diagup \\ \underline{i}(A') \end{array} \langle B', \underline{p}(\dots) \rangle = \sum_{n \geq 1} \frac{1}{n!} \langle B', l_n(A', \dots, A') \rangle$$

$\{l_n\}$  –  $L_\infty$  algebra operations on  $\mathbb{M} \otimes \mathbb{A}$ .

$L_\infty$  relations  $\Leftrightarrow$  BV master equation  $\{S_{\text{eff}}, S_{\text{eff}}\} = 0$

# From $L_\infty$ back to $A_\infty$

Recovering  $A_\infty$  products on  $\mathbb{M}$ :

$$l_n(x_1 \otimes t_{12}, \dots, x_n \otimes t_{n, n+1}) = m_n(x_1, \dots, x_n) \otimes t_{1, n+1}$$

$t_{ij} \in \mathbb{A}$  matrix with  $(i, j)$ -entry 1 and all other entries 0.

# Picture 2: HTQM

## Topological quantum mechanics:

- Space of states:  $\mathcal{H}_{ab} = \Omega^\bullet(X)$   
 (for a particle of  $(a, b)$ -type,  $a \neq b$ ).
- BRST operator  $Q = d$ .
- Hamiltonian  $H = \mathcal{L}_{v_{ab}} = [Q, G]_+$ .
- $G = \iota_{v_{ab}}$ .
- Evolution operator (superpropagator):  
 $U(t, dt) = e^{-tH - dtG} \in \Omega^\bullet(\mathbb{R}_+) \otimes \text{End}(\mathcal{H}_{ab})$

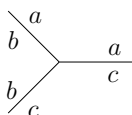
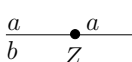
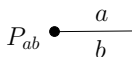


}  $N=2$  SUSY QM

$$(d_t + \text{ad}_Q)U = 0$$

# HTQM on metric trees

## HTQM on metric trees:

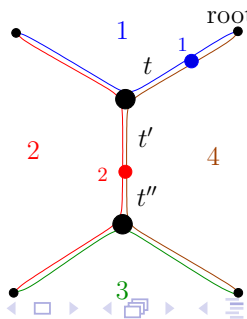
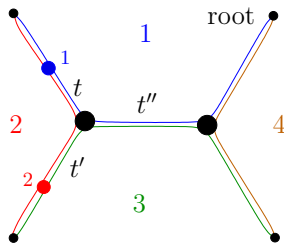
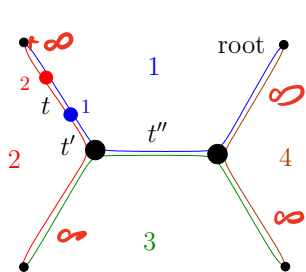
- 3-valent vertex   $\sim \mathcal{H}_{ab} \otimes \mathcal{H}_{bc} \xrightarrow{\wedge} \mathcal{H}_{ac}$
- 2-valent vertex   $\sim \text{operator } \mathcal{H}_{ab} \xrightarrow{\wedge \delta_Z} \mathcal{H}_{ab}$
- 1-valent vertex   $\sim \text{state } \delta_{\text{Unstab}_{P_{ab}}}$
- $(a, b)$ -edge of length  $t \sim U_{ab}(t, dt)$

Out of these building blocks, we build a form on the space of metric trees:

$$I \in \Omega^\bullet(MT_{N;k_1, \dots, k_N}) \otimes \text{Hom}(\text{Mor}_{1,1}^{\otimes k_1} \otimes \text{Mor}_{1,2} \otimes \dots \otimes \text{Mor}_{N-1,N} \otimes \text{Mor}_{N,N}^{\otimes k_N}, \text{Mor}_{1,N})$$

where  $\text{Mor}_{a,b} := \text{Mor}(F_a, F_b)$ .

**Example:** three top-cells in  $MT_{4;1,1,0,0}$



# Example

$$I \left( \begin{array}{c} P_{12} \\ \text{2} \\ Z \\ t \\ \text{3} \\ Z' \\ t' \\ P_{23} \end{array} \begin{array}{c} \text{1} \\ t'' \\ \text{4} \\ \text{root} \\ P_{34} \end{array} \right) = \sum_{P_{14} \in \text{Crit}(F_1 - F_4)} \bar{I} \cdot [P_{14}]$$

$$\bar{I} = \int_X \delta_{\text{Stab}_{P_{14}}} \wedge U_{13}(t'', dt'') \left( U_{12}(t, dt) (\delta_Z \wedge \delta_{\text{Unstab}_{P_{12}}}) \wedge \right. \\ \left. \wedge U_{23}(t', dt') (\delta_{Z'} \wedge \delta_{\text{Unstab}_{P_{23}}}) \right) \wedge \delta_{\text{Unstab}_{P_{34}}}$$

# Properties of $I$

- 1  $(d_{MT} + Q)I = 0$
- 2 Factorization on IR boundary of  $MT$ :

$$I \left( \text{Diagram} \right) = \langle I(T_2) \wedge I(T_1) \rangle$$

The diagram shows two circles,  $T_1$  on the left and  $T_2$  on the right. Each circle has four arrows pointing inward from its perimeter. A line with an arrow points from the right side of  $T_1$  to the left side of  $T_2$ . The label  $+\infty$  is placed above this connecting line. The entire diagram is enclosed in large parentheses.

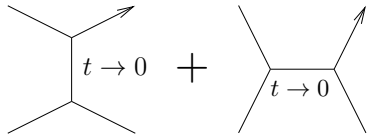
- 3 Period  $\int_{MT} I = m$  – the composition map in Fukaya-Morse category.



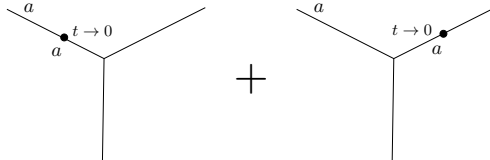
# $A_\infty$ relations from IR factorization of HTQM

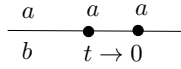
$$(d_{MT} + Q)I = 0 \quad \Rightarrow \quad \int_{\partial MT} I = -Q \underbrace{\int_{MT} I}_m \quad m(\dots m(\underline{\circ}) \dots)$$

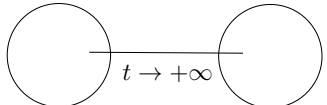
**L.h.s. (contributions of boundary strata of  $MT$ ):**

**a**  = 0

**b**  = 0

**c**  = 0

**d**   $\rightarrow$  terms  $m(\dots m(Z, Z') \dots)$  in  $A_\infty$  relation.

**e**   $\rightarrow$  terms  $m(\dots m(\underbrace{\dots}_{\geq 2 \text{ colors}}) \dots)$  in  $A_\infty$  relation.

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- O. Chekeres, A. Losev, P. Mnev, D. Youmans, “Two field-theoretic viewpoints on the Fukaya-Morse  $A_\infty$  category,” arXiv:2112.12756.