



Evidence for the Holographic dual of $N = 3$
AdS₄ × M₆ Solution in Massive II A

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Iberian Strings, 1.16-19, 2017

1511.08223 1508.05376 with Junchen Rong

- This work is aimed at looking for the supergravity dual of the simplest Chern-Simons matter CFT, inspired by [[Jafferis Talk @ Texas AM '15](#), [Guarino, Jafferis, Varela' 15](#)].
- In $D = 3$, we know that the $\mathcal{N} = 6$ ABJM model possesses a IIA dual description. It consists of two gauge groups and 4 bifundamental adjoint chirals. There are also deformations or quiver generalizations ABJM whose dual supergravity backgrounds are known [[Ceresole, Dall'Agata, D'Auria, Ferrara, Fré, Gualtieri, Lozano, Martelli, Petrini, Sparks, Tomasiello, Zaffaroni...](#)]
- We also know that there are many Chern-Simons matter CFTs with a single gauge group [[Schwarz '04](#), [Gaiotto and Yin'07](#)]. Do they have supergravity duals?

- Chern-Simons matter theory with a single gauge group can have maximal $\mathcal{N} = 3$ SUSY.
- Most of the Chern-Simons matter CFTs with a single gauge group cannot have supergravity duals, because in the spectrum of BPS single trace operators, there are higher spin operators and their number grows exponential with energy. [Minwalla, P. Narayan, T. Sharma, V. Umesh and X. Yin '11]
- However, a particular $\mathcal{N} = 2$ Chern-Simons matter theory with a single gauge group admits supergravity description in the large N limit [Guarino, Jafferis, Varela' 15].

- This $\mathcal{N} = 2$ Chern-Simons matter theory consists of a single $SU(N)$ Chern-Simons vector, 3 $SU(N)$ adjoint chirals with cubic superpotential $W = \text{Tr}(X, [Y, Z])$ invariant under $SU(3)_f$.
- This theory does not have a Lagrangian description. It can be defined via RG flow from maximal SYM deformed by a CS term which breaks the supersymmetry from $\mathcal{N} = 8$ to $\mathcal{N} = 2$.
- In the large N limit, it is dual to the $\mathcal{N} = 2$ $SU(3) \times U(1)$ invariant warped $AdS_4 \times M_6$ solution in massive IIA, where M_6 is topologically S^6 . This solution is found by lifting the $\mathcal{N} = 2$ $SU(3) \times U(1)$ invariant critical point in dyonic ISO(7) gauged maximal supergravity [Guarino, Jafferis, Varela' 15].

Can we find a CFT even simpler than the $\mathcal{N} = 2$ CFT which admits a supergravity dual?

- There is a potential candidate. Deforming the $\mathcal{N} = 2$ CFT by adding $\mathcal{N} = 2$ mass term to the chiral multiplet Z flows to an $\mathcal{N} = 3$ $SU(N)$ Chern-Simons coupled to two chiral multiplets with quartic superpotential $W = \text{Tr}([X, Y][X, Y])$ invariant under $SO(3)$ flavor symmetry [Gaiotto and Yin'07].
- It admits a Lagrangian description.
- There are no higher spin operators present in the spectrum of BPS single trace operators. The number of BPS single trace operators grows with energy in a roughly Kaluza-Klein fashion [Minwalla, P. Narayan, T. Sharma, V. Umesh and X. Yin '11].

- It was known that dyonic ISO(7) gauged maximal SUGRA admits an $\mathcal{N} = 3$ SO(3)×SO(3) invariant critical point with negative cosmological constant [Gallerati, Samtleben and Trigiante'14].
- It can be lift to a warped $AdS_4 \times M_6$ solution in massive IIA using the uplift formulas found by [Guarino, Varela' 15]. Its free energy is lower than that of the $\mathcal{N} = 2$ solution. There may exist an $\mathcal{N} = 2$ domain wall interpolating between the $\mathcal{N} = 2$ and $\mathcal{N} = 3$ solutions (yet to be constructed).
- Could this solution be the dual of the $\mathcal{N} = 3$ SCFT Chern-Simons matter theory?

- To apply the uplift formulae, we need to construct the metric \mathcal{M}_{MN} (square of the 56-bein \mathcal{V}_M^N) on the scalar coset at the critical point.
- The $\mathcal{N} = 3$ critical point is captured by the $SO(3)_D \times SO(3)_R$ invariant sector of the dyonic ISO(7) gauged maximal SUGRA [Pang, Pope, Rong'15].
- See [Dall'Agata, Inverso and Marrani'14, Guarino and Varela'15] for details about dyonic ISO(7) gauged maximal SUGRA. SO(7) is gauged electrically, while \mathbb{R}^7 is gauged dyonically

$$D = d - gA^{IJ}t_{[I}^K \delta_{J]K} + (g\delta_{IJ}A^I - m\tilde{A}_J)t_8^J.$$

- $ISO(7) \equiv SO(7) \ltimes \mathbb{R}^7$

$$\begin{aligned} SO(7) &\supset SO(4) \times SO(3)_V \supset SO(3)_R \times SO(3)_L \times SO(3)_V \\ &\supset SO(3)_R \times [SO(3)_L \times SO(3)_V]_D \end{aligned}$$

- $\text{SO}(3)_R \times \text{SO}(3)_D$ invariant sector preserves $\mathcal{N} = 1$ SUSY:
 $(g_{\mu\nu}, \psi_\mu^8), (\phi_1, \sigma_1, \chi_1), (\phi_2, \sigma_2, \chi_2)$.
- The superpotential of the $\mathcal{N} = 1$ theory is given by A_{88} , from which the scalar potential can be constructed.

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$$\xi_1 = \tanh \phi_1 e^{i\sigma_1}, \quad \xi_2 = \tanh \phi_2 e^{i\sigma_2},$$

the $\mathcal{N} = 3$ point is given by

$$\xi_1 = \frac{3}{5} - \frac{2i}{5}, \quad \xi_2 = \frac{i}{2},$$

which is a stationary point of the potential not the superpotential!
 The preserved three SUSYs lie outside the $\mathcal{N} = 1$ subsector.

- Using the uplift formulas [Guarino, Varela'15].
- In terms of the auxiliary coordinates on S^6

$$\begin{aligned} \mu^1 &= \sin \xi \cos \theta_1 \cos \chi_1, & \mu^2 &= \sin \xi \cos \theta_1 \sin \chi_1, \\ \mu^3 &= \sin \xi \sin \theta_1 \cos \psi, & \mu^4 &= \sin \xi \sin \theta_1 \sin \psi, \\ \nu^1 &= \cos \xi \cos \theta_2, & \nu^2 &= \cos \xi \sin \theta_2 \cos \chi_2, \\ \nu^3 &= \cos \xi \sin \theta_2 \sin \chi_2, \end{aligned}$$

The 10D metric is given by

$$L^{-2}d\hat{s}_{10}^2 = \Delta^{-1}\left(\frac{3\sqrt{3}}{16}ds_{AdS_4}^2\right) + g_{mn}dy^m dy^n,$$

$$\Delta = 3^{\frac{9}{8}}2^{-\frac{3}{4}}(\cos 2\xi + 3)^{-\frac{1}{8}}\Xi^{-\frac{1}{4}}, \quad \Xi = (24 \cos 2\xi + 3 \cos 4\xi + 37),$$

and the internal metric on the deformed S^6 is given as

$$g_{mn}dy^m dy^n = \frac{3\sqrt{3}}{4}(\Delta\Xi)^{-1} \left[-\sin^2 2\xi d\xi^2 + 8(\cos 2\xi + 3)d\mu \cdot d\mu \right. \\ \left. + 4(\cos 2\xi + 3)d\nu \cdot d\nu + 16\mu^A \eta_{AB}^i d\mu^B \epsilon^{ijk} \nu^j d\nu^k \right. \\ \left. - \frac{16}{\cos 2\xi + 3} (d\mu^A \eta_{AB}^i \mu^B \nu^i)^2 \right],$$

where η^i 's are the generators of $SO(3)_L$,

$$\eta^1 = \begin{pmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}, \quad \eta^2 = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix}, \quad \eta^3 = \begin{pmatrix} 0 & -1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

Denote $\mathcal{K}^i \equiv \mu^A \eta_{AB}^i d\mu^B$, various p-form fields can be expressed as

$$\begin{aligned}
L^{-1}e^{\frac{3}{4}\phi_0}\hat{A}_{(1)} &= \frac{2}{\cos 2\xi + 3}\nu^i\mathcal{K}^i, \\
L^{-2}e^{-\frac{1}{2}\phi_0}\hat{A}_{(2)} &= -\Xi^{-1}\left[-8d\nu^i \wedge \mathcal{K}^i + 6\sin 2\xi d\xi \wedge \nu^i\mathcal{K}^i \right. \\
&\quad + 2(3\cos 2\xi + 5)\nu^i\eta_{AB}^i d\mu^A \wedge d\mu^B \\
&\quad \left. - (3\cos 2\xi + 5)\epsilon^{ijk}\nu^i d\nu^j \wedge d\nu^k\right], \\
L^{-3}e^{\frac{1}{4}\phi_0}\hat{A}_{(3)} &= \Xi^{-1}\left[6\sin 2\xi\epsilon^{ijk}d\xi \wedge \mathcal{K}^i\nu^j \wedge d\nu^k \right. \\
&\quad - 2(3\cos 2\xi + 5)\epsilon^{ijk}\nu^i d\nu^j \wedge d\mu^A \wedge \eta_{AB}^k d\mu^B \\
&\quad \left. - 4\epsilon^{ijk}\mathcal{K}^i \wedge d\nu^j \wedge d\nu^k - \frac{8}{3}\csc^2\xi\epsilon^{ijk}\mathcal{K}^i \wedge \mathcal{K}^j \wedge \mathcal{K}^k\right] \\
&\quad + \frac{3\sqrt{3}}{8}\Omega_{(3)}.
\end{aligned}$$

where $d\Omega_{(3)} = \text{vol}(AdS_4)$, $L^2 = 2^{-\frac{1}{12}}g^{-25/12}m^{1/12}$ and $e^{\phi_0} = 2^{\frac{5}{6}}g^{\frac{5}{6}}m^{-\frac{5}{6}}$.

Holographic checks of the duality between $\mathcal{N} = 3$ solution in massive IIA and $\mathcal{N} = 3$ Chern-Simons matter CFT

- Matching the symmetry
bulk residual SUSY transform under $SO(3)_D \rightarrow$ CFT $SO(3)_r$ symmetry. Bulk $SO(3)_R \rightarrow$ CFT $SO(3)_f$.
- An infinite tower of Kaluza-Klein gravitons arises in AdS_4 upon compactifying massive IIA on M_6 . The short gravitons should be dual to the BPS spin-2 operators in the CFT.
- We solved the spectrum of KK gravitons around the $\mathcal{N} = 3$, $AdS_4 \times M_6$ vacuum [Pang, Rong'15]. Some of them form short $\mathcal{N} = 3$ spin-2 multiplets (BPS graviton). Others form long $\mathcal{N} = 3$ spin-2 multiplets.

$$E_0 = j_r + 3, \quad j_r \in \mathbb{Z}^+ \cup \{0\}, \quad j_f = 0.$$

- The fact that the short gravitons are all singlets of $SO(3)_f$ is not obvious, since $SO(3)_f$ commutes with $SO(3)_r$. Our results match with the CFT computation [Minwalla, P. Narayan, T. Sharma, V. Umesh and X. Yin '11].

- Matching the free energies on both sides.
- Tree level free energy of the $\mathcal{N} = 3$ solution.

$$F_{\text{gravity}} = \frac{16\pi^3 \ell^2}{(2\pi\ell_s)^8 g^6} \text{Vol}(S^6), \quad \ell^2 = \frac{3\sqrt{3}}{16} g^{-7/3} (m/2)^{1/3},$$

- Holographic dictionary

Page charge \rightarrow rank of gauge group

$$-\frac{1}{(2\pi\ell_s)^5} \int_{M^6} e^{\frac{1}{2}\phi} \hat{*} F_{(4)} + A_{(2)} \wedge dA_{(3)} + \frac{1}{6} m A_{(2)} \wedge A_{(2)} \wedge A_{(2)} = N,$$

$$\mathcal{N} = 3 \quad \text{solution :} \quad \frac{1}{(2\pi\ell_s)^5 g^5} \frac{16\pi^3}{3} = N.$$

Chern-Simons level k is related to the Romans mass [[Gaiotto and Tomasiello '09](#)].

$$m = F_{(0)} = \frac{k}{2\pi\ell_s},$$



$$F_{\text{gravity}} = \frac{9\pi}{40} 3^{1/6} k^{1/3} N^{5/3}.$$

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- Free energy of SU(N) gauged Chern-Simons theory with 2 adjoint chiral multiplets on S^3 can be computed via localization. In the limit $N \gg k$

$$F_{\text{SCFT}}^{\mathcal{N}=3} = \frac{9\pi}{40} 3^{1/6} k^{1/3} N^{5/3} = F_{\text{gravity}},$$

- Is supergravity computation valid in the regime $N \gg k$?

In the string frame, curvature radius of bulk solution $R/\ell_s \sim (N/k)^{1/6} \gg 1$. String coupling $g_s \sim 1/(N^{1/6} k^{5/6})$. Yes, supergravity approximation is valid.

Conclusion

- We provided evidence to a new example of AdS_4/CFT_3 , in which the bulk solution is the $\mathcal{N} = 3$ $AdS_4 \times M_6$ solution in massive IIA. The boundary CFT is the so far simplest one, with one $SU(N)$ Chern-Simons gauge field and two adjoint chiral multiplets.
- More tests of this duality shall be carried out in future.

Thank You!