Bootstrapping $\mathcal{N} = 2$ SCFTs

Madalena Lemos



Iberian Strings Lisbon, Jan. 18 2017

Based on: 1312.5344 w/ C. Beem, P. Liendo, W. Peelaers, L. Rastelli and B. van Rees 1511.07449 w/ P. Liendo 1612.01536 w/ P. Liendo, C. Meneghelli and V. Mitev

1702.xxxx w/ M. Cornagliotto and V. Schomerus

Outline

- **1** The (Super)conformal Bootstrap Program
- **2** A solvable subsector
- **3** Constraining the space of $\mathcal{N} = 2$ SCFTs
- 4 $4d \mathcal{N} = 3 \text{ SCFTs}$
- **5** Summary and Outlook

Outline

1 The (Super)conformal Bootstrap Program

- **2** A solvable subsector
- **3** Constraining the space of $\mathcal{N} = 2$ SCFTs
- 4 $\mathcal{A} = 3$ SCFTs
- **5** Summary and Outlook

What is the space of consistent (S)CFTs?

What is the space of consistent (S)CFTs?

 $\rightarrow\,$ Maximally supersymmetric theories: well known list of theories

What is the space of consistent (S)CFTs?

 $\rightarrow\,$ Maximally supersymmetric theories: well known list of theories

 $\rightarrow \ \mathcal{N}=2$ theories: large known list of theories many lacking a Lagrangian description

What is the space of consistent (S)CFTs?

- $\rightarrow\,$ Maximally supersymmetric theories: well known list of theories
- $\rightarrow \ \mathcal{N}=$ 3 theories: not known to exist until García-Etxebarria and Regalado
- $\rightarrow~\mathcal{N}=2$ theories: large known list of theories many lacking a Lagrangian description

What is the space of consistent (S)CFTs?

- $\rightarrow\,$ Maximally supersymmetric theories: well known list of theories
- $\rightarrow \ \mathcal{N}=3$ theories: not known to exist until García-Etxebarria and Regalado
- $\rightarrow~\mathcal{N}=2$ theories: large known list of theories many lacking a Lagrangian description

Can we bootstrap specific theories?

What is the space of consistent (S)CFTs?

- $\rightarrow\,$ Maximally supersymmetric theories: well known list of theories
- $\rightarrow \ \mathcal{N}=3$ theories: not known to exist until García-Etxebarria and Regalado
- $\rightarrow~\mathcal{N}=2$ theories: large known list of theories many lacking a Lagrangian description

Can we bootstrap specific theories?

- $\rightarrow\,$ Particularly helpful if theory is uniquely fixed by a set of discrete data
- \rightarrow Only tool available for finite N non-Lagrangian theories

 Various conformal families related by action of supercharges

- Various conformal families related by action of supercharges
- ► Conformal blocks ~→ superconformal blocks

- Various conformal families related by action of supercharges
- ► Conformal blocks ~→ superconformal blocks
- Finite re-organization of an infinite amount of data

- Various conformal families related by action of supercharges
- ► Conformal blocks ~→ superconformal blocks
- Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

- Various conformal families related by action of supercharges
- ► Conformal blocks ~→ superconformal blocks
- Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

ightarrow Yes, for 4d $\mathcal{N}\geqslant 2$ [Beem, ML, Liendo, Peelaers, Rastelli, van Rees]

- Various conformal families related by action of supercharges
- ► Conformal blocks ~→ superconformal blocks
- Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

 \rightarrow Yes, for 4*d* $\mathcal{N} \ge 2$ [Beem, ML, Liendo, Peelaers, Rastelli, van Rees] (6*d* $\mathcal{N} = (2,0)$ and 2*d* $\mathcal{N} = (0,4)$ [Beem, Rastelli, van Rees])

- Various conformal families related by action of supercharges
- ► Conformal blocks ~→ superconformal blocks
- Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

- \rightarrow Yes, for 4*d* $N \ge 2$ [Beem, ML, Liendo, Peelaers, Rastelli, van Rees] (6*d* N = (2,0) and 2*d* N = (0,4) [Beem, Rastelli, van Rees])
 - Step 1: Solve this subsector

- Various conformal families related by action of supercharges
- ► Conformal blocks ~→ superconformal blocks
- Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

- \rightarrow Yes, for 4*d* $N \ge 2$ [Beem, ML, Liendo, Peelaers, Rastelli, van Rees] (6*d* N = (2,0) and 2*d* N = (0,4) [Beem, Rastelli, van Rees])
 - Step 1: Solve this subsector
 - Step 2: Full blown numerics for the rest

Outline

1 The (Super)conformal Bootstrap Program

2 A solvable subsector

3 Constraining the space of $\mathcal{N} = 2$ SCFTs

4 $\mathcal{A} = 3$ SCFTs

5 Summary and Outlook

Organize operators in representations of superconformal algebra

 $\{\mathcal{O}_{\Delta,(j_1,j_2),}\}$

Organize operators in representations of superconformal algebra

 $\{\mathcal{O}_{\Delta,(j_1,j_2),\underbrace{R,r}_{SU(2)_R},\underbrace{r}_{U(1)_r}\}$

Organize operators in representations of superconformal algebra

7/21

$$\{\mathcal{O}_{\Delta,(j_1,j_2),\underbrace{R}_{SU(2)_R},\underbrace{r}_{U(1)_r}\}$$

Claim

 $\rightarrow~\mathsf{Pick}$ a plane $\mathbb{R}^2\in\mathbb{R}^4$

Organize operators in representations of superconformal algebra

◆□ > ◆□ > ◆三 > ◆三 > 三日 のへの

7/21

$$\{\mathcal{O}_{\Delta,(j_1,j_2),\underbrace{R,r}_{SU(2)_R},\underbrace{r}_{U(1)_r}\}$$

Claim

- $\rightarrow~\mathsf{Pick}$ a plane $\mathbb{R}^2\in\mathbb{R}^4$
- ightarrow Restrict to operators with $\Delta=2R+j_1+j_2$

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta,(j_1,j_2),\underbrace{R}_{SU(2)_R},\underbrace{r}_{U(1)_r}\}$$

Claim

- $\rightarrow~\mathsf{Pick}$ a plane $\mathbb{R}^2\in\mathbb{R}^4$
- ightarrow Restrict to operators with $\Delta=2R+j_1+j_2$

$$\langle \mathcal{O}_1^{I_1}(z_1, \overline{z}_1) \dots \mathcal{O}_n^{I_n}(z_n, \overline{z}_n) \rangle$$

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta,(j_1,j_2),\underbrace{R,r}_{SU(2)_R},\underbrace{r}_{U(1)_r}\}$$

Claim

- $\rightarrow~\mathsf{Pick}$ a plane $\mathbb{R}^2\in\mathbb{R}^4$
- ightarrow Restrict to operators with $\Delta=2R+j_1+j_2$

$$u_{I_1}(\overline{z}_1)\ldots u_{I_n}(\overline{z}_n)\langle \mathcal{O}_1^{I_1}(z_1,\overline{z}_1)\ldots \mathcal{O}_n^{I_n}(z_n,\overline{z}_n)\rangle$$

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta,(j_1,j_2),\underbrace{R}_{SU(2)_R},\underbrace{r}_{U(1)_r}\}$$

Claim

- $\rightarrow~\mathsf{Pick}$ a plane $\mathbb{R}^2\in\mathbb{R}^4$
- ightarrow Restrict to operators with $\Delta=2R+j_1+j_2$

$$u_{I_1}(\overline{z}_1)\ldots u_{I_n}(\overline{z}_n)\langle \mathcal{O}_1^{I_1}(z_1,\overline{z}_1)\ldots \mathcal{O}_n^{I_n}(z_n,\overline{z}_n)\rangle = f(z_i)$$

\rightarrow Meromorphic!

Why?

Subsector = Cohomology of nilpotent Q

Why?

 \blacktriangleright Subsector = Cohomology of nilpotent $\mathbb{Q}\sim \mathcal{Q}+\mathcal{S}$

Why?

- $\blacktriangleright \ \ Subsector = Cohomology \ of \ \ nilpotent \ \ \mathbb Q \sim \mathcal Q + \mathcal S$
- $\rightarrow~$ Cohomology at the origin $\Rightarrow~$ non-empty classes

Why?

- $\blacktriangleright \ \ Subsector = Cohomology \ of \ \ nilpotent \ \ \mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- $\rightarrow~\mbox{Cohomology}$ at the origin $\Rightarrow~\mbox{non-empty}$ classes

$$\Delta = 2R + j_1 + j_2$$

Why?

- $\blacktriangleright \ \ Subsector = Cohomology \ of \ \ nilpotent \ \ \mathbb Q \sim \mathcal Q + \mathcal S$
- $\rightarrow~\mbox{Cohomology}$ at the origin $\Rightarrow~\mbox{non-empty}$ classes

$$\Delta = 2R + j_1 + j_2$$

• On plane $\mathfrak{sl}_2 \times \mathfrak{sl}_2$

Why?

- $\blacktriangleright \ \ Subsector = Cohomology \ of \ \ nilpotent \ \ \mathbb Q \sim \mathcal Q + \mathcal S$
- $\rightarrow~\mbox{Cohomology}$ at the origin $\Rightarrow~\mbox{non-empty}$ classes
 - $\Delta = 2R + j_1 + j_2$



Why?

- $\blacktriangleright \ \ Subsector = Cohomology \ of \ \ nilpotent \ \ \mathbb Q \sim \mathcal Q + \mathcal S$
- $\rightarrow~\mbox{Cohomology}$ at the origin $\Rightarrow~\mbox{non-empty}$ classes
 - $\Delta = 2R + j_1 + j_2$



Why?

- $\blacktriangleright \ \ Subsector = Cohomology \ of \ \ nilpotent \ \ \mathbb Q \sim \mathcal Q + \mathcal S$
- $\rightarrow\,$ Cohomology at the origin \Rightarrow non-empty classes
 - $\Delta = 2R + j_1 + j_2$
- On plane \mathfrak{sl}_2 $\times \mathfrak{sl}_2$ $\overset{\circ}{\mathfrak{sl}_2}$ $\overset{\circ}{\mathfrak{sl}$

< □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □ > < □

Why?

- $\blacktriangleright \ \ Subsector = Cohomology \ of \ \ nilpotent \ \ \mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- $\rightarrow~\mbox{Cohomology}$ at the origin $\Rightarrow~\mbox{non-empty}$ classes

$$\Delta = 2R + j_1 + j_2$$

- On plane \mathfrak{sl}_2 $\times \mathfrak{sl}_2$ does not does not
- ightarrow twisted translations $u_I(ar{z})$
- \hookrightarrow diagonal subalgebra $\bar{\mathfrak{sl}}_2 \times \mathfrak{su}(2)_R$ is Q exact

Why?

- $\blacktriangleright \ \ Subsector = Cohomology \ of \ \ nilpotent \ \ \mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- $\rightarrow~\mbox{Cohomology}$ at the origin $\Rightarrow~\mbox{non-empty}$ classes

$$\Delta = 2R + j_1 + j_2$$

- On plane $\mathfrak{sl}_2 \times \mathfrak{sl}_2$ commutes with \mathbb{Q} does not
- ightarrow twisted translations $u_I(ar{z})$
- \hookrightarrow diagonal subalgebra $\bar{\mathfrak{sl}}_2 \times \mathfrak{su}(2)_R$ is Q exact

$$\rightarrow h = R + j_1 + j_2$$

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology
Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

9/21

$$Q' = \begin{bmatrix} Q \ ilde{Q}^{\star} \end{bmatrix}, \qquad ilde{Q}' = \begin{bmatrix} ilde{Q} \ -Q^{\star} \end{bmatrix}$$

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = egin{bmatrix} Q \ ilde{Q}^{\star} \end{bmatrix}, \qquad ilde{Q}' = egin{bmatrix} ilde{Q} \ -Q^{\star} \end{bmatrix}$$

$$u_I = (1, \overline{z})$$

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \ ilde{Q}^{\star} \end{bmatrix}, \qquad ilde{Q}' = \begin{bmatrix} ilde{Q} \ -Q^{\star} \end{bmatrix}$$

$$u_I = (1, \overline{z})$$

$$q(z,\bar{z}) = u_I Q^I = Q(z,\bar{z}) + \bar{z} \tilde{Q}^*(z,\bar{z}),$$

$$\widetilde{q}(z,\overline{z}) = u_I \widetilde{Q}^I = \widetilde{Q}(z,\overline{z}) - \overline{z} Q^\star(z,\overline{z})$$

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \ ilde{Q}^{\star} \end{bmatrix}, \qquad ilde{Q}' = \begin{bmatrix} ilde{Q} \ -Q^{\star} \end{bmatrix}$$

$$u_I = (1, \overline{z})$$

$$q(z,\overline{z}) = u_I Q^I = Q(z,\overline{z}) + \overline{z} \widetilde{Q}^{\star}(z,\overline{z}),$$

$$ilde{q}(z,ar{z}) = u_I ilde{Q}^I = ilde{Q}(z,ar{z}) - ar{z} Q^\star(z,ar{z})$$

$$\rightarrow q(z,\bar{z})\tilde{q}(0) \sim \bar{z}\tilde{Q}^{\star}(z,\bar{z})\tilde{Q}(0) \sim \frac{\bar{z}}{z\bar{z}} = \frac{1}{z}$$

Which operators are in the cohomology?

ightarrow Stress tensor $T_{\mu
u}$

Which operators are in the cohomology?

ightarrow Stress tensor $\mathcal{T}_{\mu
u} \rightsquigarrow$ superdescendant

- \rightarrow Stress tensor $T_{\mu
 u} \rightsquigarrow$ superdescendant
- $\rightarrow~$ Stress tensor supermultiplet

- \rightarrow Stress tensor $T_{\mu
 u} \rightsquigarrow$ superdescendant
- ightarrow Stress tensor supermultiplet \Rightarrow 2*d* stress tensor

- \rightarrow Stress tensor $T_{\mu
 u} \rightsquigarrow$ superdescendant
- ightarrow Stress tensor supermultiplet \Rightarrow 2*d* stress tensor

$$T(z)T(0) \sim -12 \frac{c_{4d}/2}{z^4} + 2 \frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

Which operators are in the cohomology?

- \rightarrow Stress tensor $T_{\mu
 u} \rightsquigarrow$ superdescendant
- ightarrow Stress tensor supermultiplet \Rightarrow 2*d* stress tensor

$$T(z)T(0) \sim -12 \frac{c_{4d}/2}{z^4} + 2 \frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

 $\,\hookrightarrow\,$ Global \mathfrak{sl}_2 enhances to Virasoro

Which operators are in the cohomology?

- \rightarrow Stress tensor $T_{\mu
 u} \rightsquigarrow$ superdescendant
- ightarrow Stress tensor supermultiplet \Rightarrow 2*d* stress tensor

$$T(z)T(0) \sim -12 \frac{c_{4d}/2}{z^4} + 2 \frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

 $\, \hookrightarrow \, \operatorname{\mathsf{Global}}\, \mathfrak{sl}_2 \text{ enhances to Virasoro}$

 \hookrightarrow $c_{2d} = -12c_{4d}$

Which operators are in the cohomology?

 $\rightarrow\,$ Flavor symmetries current multiplet

- $\rightarrow\,$ Flavor symmetries current multiplet
- $\, \hookrightarrow \, \operatorname{Affine} \, \operatorname{Kac} \, \operatorname{Moody} \, \operatorname{current} \, \operatorname{algebra}$

$$J^{a}(z)J^{b}(0)\sim -rac{k_{4d}/2\delta^{ab}}{z^{2}}+i\;f^{abc}rac{J^{c}(0)}{z}+\ldots\;,$$

Which operators are in the cohomology?

- $\rightarrow\,$ Flavor symmetries current multiplet
- $\, \hookrightarrow \, \operatorname{Affine} \, \operatorname{Kac} \, \operatorname{Moody} \, \operatorname{current} \, \operatorname{algebra}$

$$J^{a}(z)J^{b}(0)\sim -rac{k_{4d}/2\delta^{ab}}{z^{2}}+i\;f^{abc}rac{J^{c}(0)}{z}+\ldots\;,$$

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三回日 のの()

$$\hookrightarrow$$
 $k_{2d} = -\frac{k_{4d}}{2}$

- $\rightarrow\,$ Flavor symmetries current multiplet
- $\, \hookrightarrow \, \operatorname{Affine} \, \operatorname{Kac} \, \operatorname{Moody} \, \operatorname{current} \, \operatorname{algebra}$

$$J^{a}(z)J^{b}(0)\sim -rac{k_{4d}/2\delta^{ab}}{z^{2}}+i\;f^{abc}rac{J^{c}(0)}{z}+\ldots\;,$$

4*d* $\mathcal{N} = 2$ **SCFTs with a flavor symmetry** $\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$

4d $\mathcal{N} = 2$ SCFTs with a flavor symmetry $\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$

Block decomposition:



4*d* $\mathcal{N} = 2$ **SCFTs with a flavor symmetry** $\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TT J^a J^b \rangle$

Block decomposition:



4*d* $\mathcal{N} = 2$ **SCFTs with a flavor symmetry** $\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TT J^a J^b \rangle$

Block decomposition:



assumptions: interacting theory, unique stress tensor

4*d* $\mathcal{N} = 2$ **SCFTs with a flavor symmetry** $\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TT J^a J^b \rangle$

Block decomposition:



assumptions: interacting theory, unique stress tensor

4*d* $\mathcal{N} = 2$ **SCFTs with a flavor symmetry** $\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TT J^a J^b \rangle$

Block decomposition:



assumptions: interacting theory, unique stress tensor

Outline

1 The (Super)conformal Bootstrap Program

2 A solvable subsector

3 Constraining the space of $\mathcal{N} = 2$ SCFTs

- 4 $4d \mathcal{N} = 3 \text{ SCFTs}$
- **5** Summary and Outlook

$4d \mathcal{N} = 2$ SCFTs with E_6 flavor symmetry



[Beem, ML, Liendo, Peelaers, Rastelli, van Rees] [ML, Liendo], (문화 (문화 (문화)), (Endo)), (Endo)),

$4d \mathcal{N} = 2$ SCFTs with E_6 flavor symmetry



[Beem, ML, Liendo, Rastelli, van Rees]

Outline

- **1** The (Super)conformal Bootstrap Program
- **2** A solvable subsector
- **3** Constraining the space of $\mathcal{N} = 2$ SCFTs
- 4 $4d \mathcal{N} = 3 \text{ SCFTs}$
- **5** Summary and Outlook

$4d \ \mathcal{N} = 3 \ \text{SCFTs}$

$\rightarrow\,$ Non-trivial interacting theories

$4d \ \mathcal{N} = 3 \ \text{SCFTs}$

$\rightarrow\,$ Non-trivial interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

 $\rightarrow\,$ Non-Lagrangian, isolated

$4d \mathcal{N} = 3 \text{ SCFTs}$

$\rightarrow\,$ Non-trivial interacting theories

- $\rightarrow\,$ Non-Lagrangian, isolated
- $\rightarrow~$ No flavor symmetry

$4d \ \mathcal{N} = 3 \ \text{SCFTs}$

$\rightarrow\,$ Non-trivial interacting theories

- $\rightarrow\,$ Non-Lagrangian, isolated
- $\rightarrow~$ No flavor symmetry
- $\rightarrow c = a$

$4d \ \mathcal{N} = 3 \ \text{SCFTs}$

$\rightarrow\,$ Non-trivial interacting theories

- $\rightarrow\,$ Non-Lagrangian, isolated
- $\rightarrow~$ No flavor symmetry
- $\rightarrow c = a$
- \rightarrow Just another SCFT

$4d \ \mathcal{N} = 3 \ \text{SCFTs}$

$\rightarrow\,$ Non-trivial interacting theories

- $\rightarrow\,$ Non-Lagrangian, isolated
- $\rightarrow~$ No flavor symmetry
- $\rightarrow c = a$
- \rightarrow Just another SCFT

▶ 4 $d N \ge 3$: some of the extra supercharges commute with Q

▶ 4*d* $N \ge 3$: some of the extra supercharges commute with Q \hookrightarrow 4*d* $N = 4 \Rightarrow 2d$ "small" N = 4 chiral algebra

▶ 4d $N \ge 3$: some of the extra supercharges commute with Q \hookrightarrow 4d $N = 4 \Rightarrow 2d$ "small" N = 4 chiral algebra \hookrightarrow 4d $N = 3 \Rightarrow 2d$ N = 2 chiral algebra [Nishinaka, Tachikawa]

▶ 4d $N \ge 3$: some of the extra supercharges commute with \mathbb{Q} $\hookrightarrow 4d \mathcal{N} = 4 \Rightarrow 2d$ "small" $\mathcal{N} = 4$ chiral algebra $\hookrightarrow 4d \mathcal{N} = 3 \Rightarrow 2d \mathcal{N} = 2$ chiral algebra [Nishinaka, Tachikawa]

 $2d \mathcal{N} = 2$ Stress tensor $\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J}\rangle$

ightarrow Local, interacting $\mathcal{N}=3$ SCFT

▶ 4*d* $N \ge 3$: some of the extra supercharges commute with Q \hookrightarrow 4d $\mathcal{N} = 4 \Rightarrow 2d$ "small" $\mathcal{N} = 4$ chiral algebra \hookrightarrow 4d $\mathcal{N} = 3 \Rightarrow 2d \mathcal{N} = 2$ chiral algebra [Nishinaka, Tachikawa]

 $2d \mathcal{N} = 2$ Stress tensor $\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J}\rangle$

 \rightarrow Local, interacting $\mathcal{N} = 3$ SCFT


Chiral algebra

▶ 4d $N \ge 3$: some of the extra supercharges commute with Q \hookrightarrow 4d $N = 4 \Rightarrow 2d$ "small" N = 4 chiral algebra \hookrightarrow 4d $N = 3 \Rightarrow 2d$ N = 2 chiral algebra [Nishinaka, Tachikawa]

 $2d \ \mathcal{N} = 2 \ \text{Stress tensor} \ \langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J} \rangle$

 $\rightarrow\,$ Local, interacting $\mathcal{N}=3$ SCFT

$$c_{4d} \geqslant rac{13}{24}$$
 [Cornagliotto, ML, Schomerus]

◆□▶ ◆□▶ ◆ヨ▶ ◆ヨ▶ ●□□ ◇◇◇

18/21

 \hookrightarrow Not saturated by any known SCFT

Chiral algebra

▶ 4d $N \ge 3$: some of the extra supercharges commute with Q \hookrightarrow 4d $N = 4 \Rightarrow 2d$ "small" N = 4 chiral algebra \hookrightarrow 4d $N = 3 \Rightarrow 2d$ N = 2 chiral algebra [Nishinaka, Tachikawa]

 $2d \ \mathcal{N} = 2 \ \text{Stress tensor} \ \langle \mathcal{J} \mathcal{J} \mathcal{J} \mathcal{J} \rangle$

 $\rightarrow\,$ Local, interacting $\mathcal{N}=3$ SCFT

$$c_{4d} \geqslant rac{13}{24}$$
 [Cornagliotto, ML, Schomerus]

 \hookrightarrow Not saturated by any known SCFT smallest interacting $c_{4d} = \frac{15}{12}$

Chiral algebra

▶ 4d $N \ge 3$: some of the extra supercharges commute with Q \hookrightarrow 4d $N = 4 \Rightarrow 2d$ "small" N = 4 chiral algebra \hookrightarrow 4d $N = 3 \Rightarrow 2d$ N = 2 chiral algebra [Nishinaka, Tachikawa]

 $2d \ \mathcal{N} = 2 \ \text{Stress tensor} \ \langle \mathcal{J} \mathcal{J} \mathcal{J} \mathcal{J} \rangle$

 $\rightarrow\,$ Local, interacting $\mathcal{N}=3$ SCFT

$$c_{4d} \geqslant \frac{13}{24}$$
 [Cornagliotto, ML, Schomerus

- \hookrightarrow Not saturated by any known SCFT smallest interacting $c_{4d} = \frac{15}{12}$
- $\hookrightarrow \mbox{ Similar bounds in } \mathcal{N} = 4 \mbox{ and } \mathcal{N} = 2 \mbox{ saturated by known} \mbox{ SCFTs [Beem, Rastelli, van Rees] [Liendo, Ramirez, Seo]}$

Outline

- **1** The (Super)conformal Bootstrap Program
- **2** A solvable subsector
- **3** Constraining the space of $\mathcal{N} = 2$ SCFTs
- 4 $\mathcal{A} = 3$ SCFTs
- **5** Summary and Outlook

 $\rightarrow\,$ New constraints on the space of allowed $\mathcal{N} \geqslant 2 \; \text{SCFTs}$

- $\rightarrow~$ New constraints on the space of allowed $\mathcal{N} \geqslant 2~\text{SCFTs}$
- $\rightarrow\,$ Can the numerical bootstrap complement these?

- $\rightarrow~$ New constraints on the space of allowed $\mathcal{N} \geqslant 2~\text{SCFTs}$
- $\rightarrow\,$ Can the numerical bootstrap complement these?
- $\,\hookrightarrow\,$ Need more superconformal blocks

- $\rightarrow~$ New constraints on the space of allowed $\mathcal{N} \geqslant 2~\text{SCFTs}$
- $\rightarrow\,$ Can the numerical bootstrap complement these?
- $\,\hookrightarrow\,$ Need more superconformal blocks

Numerically solving theories?

 $\rightarrow\,$ Bounds select special set of theories, can we solve them?

- $\rightarrow~$ New constraints on the space of allowed $\mathcal{N} \geqslant 2~\text{SCFTs}$
- $\rightarrow\,$ Can the numerical bootstrap complement these?
- $\,\hookrightarrow\,$ Need more superconformal blocks

- $\rightarrow\,$ Bounds select special set of theories, can we solve them?
- \rightarrow Is there an exotic "minimal" $\mathcal{N}=3$ SCFT with $c=rac{13}{24}?$

- $\rightarrow~$ New constraints on the space of allowed $\mathcal{N} \geqslant 2~\text{SCFTs}$
- $\rightarrow\,$ Can the numerical bootstrap complement these?
- $\,\hookrightarrow\,$ Need more superconformal blocks

- $\rightarrow\,$ Bounds select special set of theories, can we solve them?
- \rightarrow Is there an exotic "minimal" $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$?
- \hookrightarrow Need more superconformal blocks

- $\rightarrow~$ New constraints on the space of allowed $\mathcal{N} \geqslant 2~\text{SCFTs}$
- $\rightarrow\,$ Can the numerical bootstrap complement these?
- $\,\hookrightarrow\,$ Need more superconformal blocks

- $\rightarrow\,$ Bounds select special set of theories, can we solve them?
- \rightarrow Is there an exotic "minimal" $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$?
- $\,\hookrightarrow\,$ Need more superconformal blocks
- \rightarrow Can we reach the "minimal" known $\mathcal{N} = 3$ SCFT $c = \frac{15}{12}$?

- $\rightarrow~$ New constraints on the space of allowed $\mathcal{N} \geqslant 2~\text{SCFTs}$
- $\rightarrow\,$ Can the numerical bootstrap complement these?
- $\,\hookrightarrow\,$ Need more superconformal blocks

- $\rightarrow\,$ Bounds select special set of theories, can we solve them?
- \rightarrow Is there an exotic "minimal" $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$?
- \hookrightarrow Need more superconformal blocks
- \rightarrow Can we reach the "minimal" known $\mathcal{N} = 3$ SCFT $c = \frac{15}{12}$? [ML, Liendo, Meneghelli, Mitev]

- $\rightarrow~$ New constraints on the space of allowed $\mathcal{N} \geqslant 2~\text{SCFTs}$
- $\rightarrow\,$ Can the numerical bootstrap complement these?
- \hookrightarrow Need more superconformal blocks

Numerically solving theories?

- $\rightarrow\,$ Bounds select special set of theories, can we solve them?
- \rightarrow Is there an exotic "minimal" $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$?
- \hookrightarrow Need more superconformal blocks
- → Can we reach the "minimal" known $\mathcal{N} = 3$ SCFT $c = \frac{15}{12}$? [ML, Liendo, Meneghelli, Mitev]

Other analytic approaches?

- $\rightarrow~$ New constraints on the space of allowed $\mathcal{N} \geqslant 2~\text{SCFTs}$
- $\rightarrow\,$ Can the numerical bootstrap complement these?
- \hookrightarrow Need more superconformal blocks

Numerically solving theories?

- $\rightarrow\,$ Bounds select special set of theories, can we solve them?
- \rightarrow Is there an exotic "minimal" $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$?
- \hookrightarrow Need more superconformal blocks
- \rightarrow Can we reach the "minimal" known $\mathcal{N} = 3$ SCFT $c = \frac{15}{12}$? [ML, Liendo, Meneghelli, Mitev]

Other analytic approaches?

 \rightarrow Is *c* bounded by *k* for large *k*?

Thank you!

Backup slides

<□> <圕> <≧> <≧> <≧> ≥|≥ ∽੧< 1/11

Outline

Conformal bootstrap

Constraining the space of $\mathcal{N} = 2$ SCFTs Bootstrapping $\mathcal{N} = 3$ SCFTs

Conformal field theory defined by

Set of local operators and their correlation functions

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

 $\mathcal{O}_k(0)$

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

Operator Product Expansion $\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_k \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k}$

 $\mathcal{O}_k(0)$

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

Operator Product Expansion

 $\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_k \qquad \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k}$

 $\rightarrow\,$ Finite radius of convergence

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

Operator Product Expansion

 $\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k \text{prim.}} \lambda_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$

 $\rightarrow\,$ Finite radius of convergence

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k ext{prim.}} \lambda_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$$

- \rightarrow Finite radius of convergence
- $\rightarrow n-{\rm point}$ function by recursive use of the OPE until $\langle \mathbb{1}\rangle=1$

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\{\lambda_{\mathcal{O}_i\mathcal{O}_j\mathcal{O}_k}\}$

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k ext{prim.}} \lambda_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$$

- \rightarrow Finite radius of convergence
- $\rightarrow n-{\rm point}$ function by recursive use of the OPE until $\langle \mathbb{1}\rangle=1$

Conformal field theory defined by

Set of local operators and their correlation functions

CFT data $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\{\lambda_{\mathcal{O}_i\mathcal{O}_j\mathcal{O}_k}\}$

Operator Product Expansion $\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k \text{prim.}} \lambda_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$

- \rightarrow Finite radius of convergence
- $\rightarrow n-{\rm point}$ function by recursive use of the OPE until $\langle \mathbb{1} \rangle = 1$

CFT data strongly constrained

- Unitarity
- ► Associativity of the operator product algebra

Crossing Symmetry $\langle (\mathcal{O}_1(x_1) \ \mathcal{O}_2(x_2))\mathcal{O}_3(x_3) \ \mathcal{O}_4(x_4) \rangle =$







where
$$\Delta_{\mathcal{O}_i} = \Delta_{\mathcal{O}}$$
, $u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2} = z\bar{z}$, $v = \frac{x_{23}^2 x_{14}^2}{x_{13}^2 x_{24}^2} = (1-z)(1-\bar{z})$



where $\Delta_{\mathcal{O}_i} = \Delta_{\mathcal{O}}$, $u = \frac{x_{12}^2 x_{34}^2}{x_{13}^2 x_{24}^2} = z\overline{z}$, $v = \frac{x_{23}^2 x_{14}^2}{x_{13}^2 x_{24}^2} = (1-z)(1-\overline{z})$

 $\rightarrow\,$ Solve crossing equations for all four-point functions

 $\rightarrow\,$ Solve crossing equations for all four-point functions

- [Rattazzi, Rychkov, Tonni, Vichi]
 - Solving \Rightarrow constraining

 $\rightarrow\,$ Solve crossing equations for all four-point functions

[Rattazzi, Rychkov, Tonni, Vichi]

- Solving \Rightarrow constraining
 - ightarrow Guess for the spectrum

 $\rightarrow\,$ Solve crossing equations for all four-point functions

[Rattazzi, Rychkov, Tonni, Vichi]

- Solving \Rightarrow constraining
 - ightarrow Guess for the spectrum

 $\,\hookrightarrow\,$ there's a large gap in the spectrum

 $\rightarrow\,$ Solve crossing equations for all four-point functions

[Rattazzi, Rychkov, Tonni, Vichi]

- Solving \Rightarrow constraining
 - ightarrow Guess for the spectrum

 $\,\hookrightarrow\,$ there's a large gap in the spectrum

 \rightarrow Can it ever define a consistent CFT?

 $\rightarrow\,$ Solve crossing equations for all four-point functions

[Rattazzi, Rychkov, Tonni, Vichi]

- Solving \Rightarrow constraining
 - ightarrow Guess for the spectrum
 - $\,\hookrightarrow\,$ there's a large gap in the spectrum
 - \rightarrow Can it ever define a consistent CFT?

Sum rule: identical scalars ϕ

 $\rightarrow\,$ Solve crossing equations for all four-point functions

[Rattazzi, Rychkov, Tonni, Vichi]

- Solving \Rightarrow constraining
 - $\rightarrow~$ Guess for the spectrum

 $\,\hookrightarrow\,$ there's a large gap in the spectrum

 \rightarrow Can it ever define a consistent CFT?

Sum rule: identical scalars ϕ

$$ightarrow \,$$
 Identity operator $\lambda_{\mathcal{OOI}} = 1$

$$1 = \sum_{\substack{\mathcal{O}_{\Delta_{\ell}} \neq \mathbb{1} \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 \underbrace{\frac{u^{\Delta_{\phi}} g_{\Delta,\ell}(v,u) - v^{\Delta_{\phi}} g_{\Delta,\ell}(u,v)}{v^{\Delta_{\phi}} - u^{\Delta_{\phi}}}}_{F_{\Delta,\ell}}$$
Sum rule

 $\boxed{1 = \sum_{\substack{\mathcal{O}_{\Delta_{\ell}} \neq \mathbb{1} \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}}$

Sum rule

 $\boxed{1 = \sum_{\substack{\mathcal{O}_{\Delta_{\ell}} \neq \mathbb{1} \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}}$

Find Functional Ψ such that

$$\begin{array}{l} \hookrightarrow \ \psi \cdot 1 < 0 \ (1) \\ \hookrightarrow \ \psi \cdot \mathcal{F}_{\Delta,\ell}(u,v) \geq 0 \ \text{for all} \ \{\Delta,\ell\} \ \text{in spectrum} \end{array}$$

Sum rule

 $1 = \sum_{\substack{\mathcal{O}_{\Delta_\ell}
eq \mathbb{1} \ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}$

Find Functional Ψ such that

$$\begin{array}{l} \hookrightarrow \ \psi \cdot 1 < 0 \ (\mathbb{1}) \\ \hookrightarrow \ \psi \cdot F_{\Delta,\ell}(u,v) \geq 0 \ \text{for all} \ \{\Delta,\ell\} \ \text{in spectrum} \end{array}$$

 $\rightarrow\,$ Spectrum is inconsistent $\Rightarrow\,$ rule out CFT

Sum rule

$$1 = \sum_{\substack{\mathcal{O}_{\Delta_{\ell}} \neq \mathbb{1} \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}$$

- Find Functional Ψ such that
 - $\begin{array}{l} \hookrightarrow \ \psi \cdot 1 < 0 \ (\mathbb{1}) \\ \hookrightarrow \ \psi \cdot \mathcal{F}_{\Delta,\ell}(u,v) \geq 0 \ \text{for all} \ \{\Delta,\ell\} \ \text{in spectrum} \end{array}$
- $\rightarrow\,$ Spectrum is inconsistent \Rightarrow rule out CFT
 - Truncate

$$\psi = \sum_{m,n}^{m,n \leqslant \Lambda} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{z=\bar{z}=\frac{1}{2}}$$

Sum rule

 $1 = \sum_{\substack{\mathcal{O}_{\Delta_\ell}
eq \mathbb{1} \ \mathcal{O} \in \phi \phi}} \lambda_{\phi \phi \mathcal{O}}^2 \mathcal{F}_{\Delta, \ell}$

Find Functional Ψ such that

$$\begin{array}{l} \hookrightarrow \ \psi \cdot 1 < 0 \ (\mathbb{1}) \\ \hookrightarrow \ \psi \cdot \mathcal{F}_{\Delta,\ell}(u,v) \geq 0 \ \text{for all} \ \{\Delta,\ell\} \ \text{in spectrum} \end{array}$$

- $\rightarrow\,$ Spectrum is inconsistent \Rightarrow rule out CFT
 - Truncate

$$\psi = \sum_{m,n=1}^{m,n \leqslant \Lambda} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{z=\bar{z}=\frac{1}{2}}$$

 $\rightarrow\,$ Increase $\Lambda \Rightarrow$ bounds get stronger

Sum rule

 $1 = \sum_{\substack{\mathcal{O}_{\Delta_\ell}
eq \mathbb{1} \ \mathcal{O} \in \phi \phi}} \lambda_{\phi \phi \mathcal{O}}^2 F_{\Delta,\ell}$

Find Functional Ψ such that

$$\begin{array}{l} \hookrightarrow \ \psi \cdot 1 < 0 \ (\mathbb{1}) \\ \hookrightarrow \ \psi \cdot \mathcal{F}_{\Delta,\ell}(u,v) \geq 0 \ \text{for all} \ \{\Delta,\ell\} \ \text{in spectrum} \end{array}$$

- $\rightarrow\,$ Spectrum is inconsistent \Rightarrow rule out CFT
 - Truncate

$$\psi = \sum_{m,n=1}^{m,n \leqslant \Lambda} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{z=\bar{z}=\frac{1}{2}}$$

- $\rightarrow\,$ Increase $\Lambda \Rightarrow$ bounds get stronger
- \rightarrow Always true bounds

3d Ising Model

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, PRD 86 025022]



< □ > < @ > < 볼 > < 볼 > 볼 = 의익은 7/11

3d Ising Model

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, PRD 86 025022]



 \rightarrow Saturated by 3*d* Ising model

3d Ising Model

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, PRD 86 025022]



- \rightarrow Saturated by 3*d* Ising model
- ightarrow 3*d* Ising lives at "kink"

Outline

Conformal bootstrap Constraining the space of $\mathcal{N} = 2$ SCFTs Bootstrapping $\mathcal{N} = 3$ SCFTs

What is the space of consistent SCFTs?

4*d* $\mathcal{N} = 2$ **SCFTs with** SU(2) **flavor symmetry** $\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$



Outline

Conformal bootstrap Constraining the space of $\mathcal{N}=2$ SCFTs Bootstrapping $\mathcal{N}=3$ SCFTs

4d $\mathcal{N} = 3$ SCFTs with $c = \frac{15}{12}$



[ML, Liendo, Meneghelli, Mitev]