## A Momentum/Complexity Correspondence

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- Iberian Strings 2021, IST Lisboa

- Based on 1912.05996 / 2006.06607 / 2012.02603
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$$P_{\rho} = m \, \frac{\mathrm{d}\rho}{\mathrm{d}\tau} \sim m \, e^{\lambda t}$$

 $\lambda = \frac{2\pi}{R}$  (maximal Lyapunov exponent)



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$$P_{\rho} = m \, \frac{\mathrm{d}\rho}{\mathrm{d}\tau} \, \sim \, m \, e^{\lambda t}$$

 $\lambda = \frac{2\pi}{2}$ (maximal Lyapunov exponent)

$$P_{\rho} \leftrightarrow \text{size}$$



- $\bullet$ into a black hole and the size of the precursor in the fast scrambler picture.
- throat of a near extremal RN black hole  $\leftrightarrow$  Comparison with SYK. [Brown, Gharibyan, Streicher, Susskind, Thorlacius, Zhao '18] [Susskind '19] [Maldacena, Lin, Zhao '19]



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• The correspondence was refined studying the motion of a particle falling through the

[Roberts, Stanford, Streicher '18] [Qi, Streicher '19]



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 $\mathcal{H}^+$ 

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• At finite S, complexity keeps growing for a long time after scrambling. Can some notion of complexity keep track of the momentum of the particle in the black hole interior?

If the tendency for complexity to increase is related to the mechanism behind gravity, then there must be a more general correspondence outside the black hole context.

#### **Complexity = Volume**

- Consider a holographic CFT on  $\mathbf{S}^{d-1}$  and a state  $|\Psi_t\rangle$  with geometric dual. e.g.  $|\Psi_t\rangle = e^{-iHt} \mathcal{O}_{\mathrm{simple}} |0\rangle$
- The **complexity** of  $|\Psi_t\rangle$  is given by the volume (in units of  $G \ell$ ) of the extremal spatial hypersurface  $\Sigma_t$  anchored to the corresponding timeslice of  $\partial AdS$ .

$$\mathcal{C}\left(|\Psi_t\rangle\right) \propto rac{\mathrm{Vol}[\Sigma_t]}{G\ell}$$



#### **Variational Problem**

• Extremal volume rate is described by a Hamilton-Jacobi equation.



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• The (d-1)-form  $\boldsymbol{\omega}$  captures the asymptotic bending of  $\Sigma$ .

$$\boldsymbol{\omega} = (n_{\partial \Sigma} \cdot \partial_t) \boldsymbol{\epsilon}$$

$$\frac{\mathrm{d}}{\mathrm{d}t} \mathcal{C} = \int_{\partial \Sigma} \boldsymbol{\omega}$$



t

# • The strategy is to find an extension j on $\Sigma$ such that $\mathrm{d}j = \mathcal{P}$ . $rac{\mathrm{d}}{\mathrm{d}t}\,\mathcal{C} = \int_{\Sigma}\mathcal{P} = P_{\mathcal{C}}$

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- By extending onto  $\Sigma$  , we implicitly specified the bulk time variable to talk about

# $\frac{\mathrm{d}}{\mathrm{d}t} \mathcal{C} = \int_{\Sigma} \mathcal{P} = P_{\mathcal{C}}$

momentum. This time variable is given by foliating spacetime with extremal slices.

#### **Black Hole Interior**

• Extremal slices explore the black hole in respect to the asymptotic time t.



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 $\mathcal{H}^+$ 

	t
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#### **Black Hole Interior**

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 $P_{\mathcal{C}} \sim \text{constant}$ 

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Using a local notion of momentum that is smooth  $\,P_{\mathcal{C}}( au)$ , then the quantity  $\,P_{\mathcal{C}}(t)$ 

#### **Spatial component of Momentum**

- Theorem: Given a tangent vector field  $C^{\mu}$  which asymptotically is:
  - radial
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  - of magnitude r



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#### then the PVC correspondence holds



where

 $\mathcal{P}_C = -N_{\Sigma}^{\mu} T_{\mu\nu} C^{\nu} e$ 



$$oldsymbol{\mathcal{P}}_C \ + \ \int_{\Sigma} oldsymbol{R}_C \ oldsymbol{R}_C \ = -rac{1}{8\pi G} \, K_{\mu
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is a conformal Killing vector of the induced metric on  $\Sigma$  .

• The condition  $\, oldsymbol{R}_C = 0 \,$  cannot be satisfied in general. A sufficient condition is that  $C^{\mu}$ 



- is a conformal Killing vector of the induced metric on  $\Sigma$  .
- For any spherically symmetric state, there is a candidate conformal Killing vector.

$$ds_{\Sigma}^{2} = dy^{2} + r^{2}(y)d\Omega_{d-1}^{2}$$
$$C = -r(y)\partial_{y}$$

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$$\mathrm{d}s_{\Sigma}^2 = e^{2\omega} \,\left(\mathrm{d}y^2 \,+\, \sinh^2(y)\,\mathrm{d}y^2\right)$$

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- is a conformal Killing vector of the induced metric on  $\Sigma$  .
- For any spherically symmetric state, there is a candidate conformal Killing vector.
- For any state in 2+1 dimensions, there is a candidate conformal Killing vector.
- $\bullet$ formalizes the idea that gravitational clumping of matter increases complexity.

$$\frac{\mathrm{d}}{\mathrm{d}t} \,\mathcal{C} \,=\, \int_{\Sigma} \,\mathcal{P}_C \,=\, P_{\mathrm{infall}}$$

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The relevant momentum in both cases measures **infall** towards the center of the box. This



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- There is a center of infall set by the box. (Mach's principle)

 $\bullet$ 



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- There is a center of infall set by the box.
- Complexity is a radial moment of inertia.

$$\mathcal{I}_{\text{clump}} = -\frac{1}{2\ell} \sum_{i} m_i \, \mathbf{x}_i^2$$

 $\mathcal{C} = \mathcal{C}_0 + \mathcal{I}_{clump}$ 

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- There is a center of infall set by the box.
- Complexity is a radial moment of inertia.
- For an unbound gravitational system

$$\frac{\mathrm{d}^2}{\mathrm{d}t^2}\,\mathcal{C}\,<0$$

#### **Obstruction for exact PVC: gravitational radiation**



• For vacuum one-sided geometries, the whole complexity rate is given by the remainder.

$$\frac{\mathrm{d}}{\mathrm{d}t} \,\mathcal{C} \,=\, \int_{\Sigma} \boldsymbol{R}_C$$

#### **Obstruction for exact PVC: gravitational radiation**



No candidate for exact momentum of gravity. Perturbatively, we expect that  $\bullet$ 

$$\frac{\mathrm{d}}{\mathrm{d}t} \, \mathcal{C} \, \approx \, \int_{\Sigma_0} \, \mathcal{P}_{\mathrm{grav}}$$

for some notion of conserved energy-momentum pseudotensor (Landau-Lifshitz or others).

• For vacuum one-sided geometries, the whole complexity rate is given by the remainder.

$$\frac{\mathrm{d}}{\mathrm{d}t} \,\mathcal{C} \,=\, \int_{\Sigma} \boldsymbol{R}_C$$

$$\boldsymbol{\mathcal{P}}_{\text{grav}} = -N^{\mu}_{\Sigma_0} t_{\mu\nu} C^{\nu}_0$$

#### Generalization

- Theorem: Given a tangent tensor field  $M^{\mu 
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 $M^{\mu\nu\rho} \sim h^{\mu\nu} C^{\rho} - h^{\mu\rho} C^{\nu}$  and same asymptotics for  $C^{\mu}$  as before,

#### Generalization

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then the generalized PVC correspondence holds

$$\frac{\mathrm{d}}{\mathrm{d}t} \,\mathcal{C} \,=\, \int_{\Sigma} \,\mathcal{P}_C \,+\, \int_{\Sigma} \,\mathcal{W}_M \,+\, \int_{\Sigma} \,\mathcal{R}_M$$

where  $C^{\mu} = h_{\nu\rho} M^{\nu\mu\rho}$  and

$$\mathcal{W}_M = -$$

 $\gamma \nu$ and same asymptotics for  $C^{\mu}$  as before,

$$\frac{1}{16\pi G} N^{\mu}_{\Sigma} W_{\mu\nu\rho\sigma} M^{\nu\rho\sigma}$$

the extra components of the infall tensor.

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- The idea is that now  $oldsymbol{R}_M=0\,$  is always possible in one-sided situations because of

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$$\mathrm{d}s^2 = -\mathrm{d}u\,\mathrm{d}v \,+\, L^2(u)\,\left(e^{2\beta(u)}\,\mathrm{d}u\right)\,$$

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 $\bullet$ 

$$ds^{2} = -du \, dv + L^{2}(u) \left(e^{2\beta(u)} dt - \frac{d}{dt} \mathcal{C}\right) = \int_{\Sigma} \mathcal{V}$$

and the "Weyl-momentum"
$$\mathcal{W}_M = -N_\Sigma^\mu$$
for  $t$  the equate root of the Pol Poh

for  $t_{\mu\nu}$  the square root of the Bel-Robinson tensor. [Bonilla, Senovilla '97]

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- attributed to matter infall nor to the gravitational radiation.
- we assume energy conditions for matter.
- $\bullet$ gravitationally clump together. For that we need to understand what VC really is.

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More speculatively, the second law of complexity is the statistical tendency for matter to

## **Obrigado** !

#### **Topological Obstruction**

- Consider two holographic CFTs with Hamiltonian  $H_L + H_R$ . For product states,  $\bullet$ everything works the same way since  $\mathcal{C} = \mathcal{C}_L + \mathcal{C}_R$ .
- For entangled states representing classical wormholes, left and right infall are incompatible for the conformal Killing vector.



• Exact PVC holds for the Killing Hamiltonian  $H_R - H_L$ .



• For  $H_L + H_R$ , one can introduce a topological defect that swaps the direction of infall. The only contribution to the rate of complexity comes from the remainder on the defect.

## 2M

J defect