

# Static large deviations for reaction-diffusion models

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# Motivation: Statistical Mechanics

## 1. Define a Gibbs measure

- $\mathbb{T}_N = \{0, 1_N, \dots, 1 - 1_N\}$      $1_N = 1/N$
- $\{0, 1\}^{\mathbb{T}_N}$      $\eta$      $\eta(x) \in \{0, 1\}$

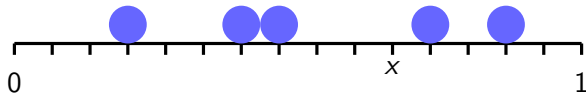
$$\mathbb{H}(\eta) = - \sum_{x \in \mathbb{T}_N} [1 - 2\eta(x + 1_N)][1 - 2\eta(x)]$$
$$\mu_\beta(\eta) = \frac{1}{Z_\beta} e^{-\beta \mathbb{H}(\eta)}$$

## 2. Define dynamics

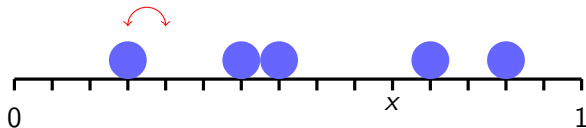
- Invariant for Gibbs measure (usually reversible)
- Here we start from the dynamics
- The invariant measure is not known explicitly

# Kawasaki dynamics

- $\mathbb{T}_N = \{0, 1/N, \dots, 1 - 1/N\}$   $\{0, 1\}^{\mathbb{T}_N}$   $\eta$   $\eta(x) \in \{0, 1\}$



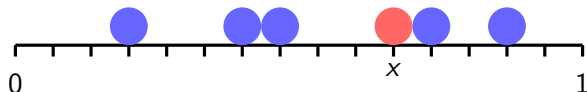
Kawasaki dynamics:



- Conservative dynamics
- Stationary states  $\nu_\alpha$ ,  $0 \leq \alpha \leq 1$
- Mixing time  $N^2$ .

# Glauber dynamics

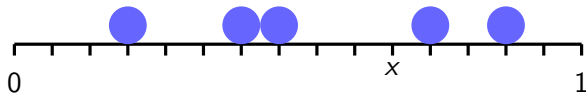
Glauber dynamics:  $c(\eta)$  local function  $c(\tau_x \eta)$   $\eta(x) \rightarrow 1 - \eta(x)$



- Non-conservative dynamics
- stationary states not known explicitly
- $\sigma(x) = 1 - 2\eta(x)$
- $c(\eta) = 1 - \gamma \sigma(0) [\sigma(1) + \sigma(-1)] + \gamma^2 \sigma(1) \sigma(-1)$   $0 \leq \gamma < 1$
- $\mathbb{H}(\sigma) = -\sum_x \sigma(x) \sigma(x+1)$
- $\mu_\gamma(\sigma) \sim \exp\{-\beta \mathbb{H}(\sigma)\}$   $\gamma = \tanh \beta$
- Mixing time 1

# The problem

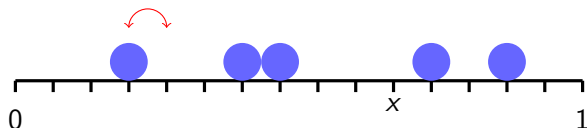
- $\mathcal{F} = \{u : \mathbb{T} \rightarrow [0, 1]\}$
- $u^{(k)} \rightarrow u \quad \int_{\mathbb{T}} G(x) u^{(k)}(x) dx \rightarrow \int_{\mathbb{T}} G(x) u(x) dx \quad G \in C(\mathbb{T})$
- $d(u, v)$



- $\eta \mapsto \rho_N \in \mathcal{F}$
- $\rho_N = \sum_{x \in \mathbb{T}_N} \eta(x) \chi_{(x, x+1_N]} \quad 1_N = 1/N$
- Glauber +  $N^2$ -Kawasaki
- $\mu_N$  stationary state not known
- Problem:  $\mathcal{A} \subset \mathcal{F}$  s.t.  $\mu_N(\mathcal{A}) \rightarrow 1$

# Kawasaki hydrodynamics

- $u_0 : \mathbb{T} \rightarrow [0, 1]$      $\rho_N(\cdot) \sim u_0(\cdot)$
- $\rho_N(t) \sim u(t)$

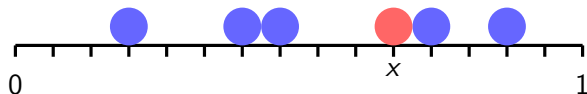


$$\begin{aligned} \frac{d}{dt} \eta_t(x) &= \eta_t(x - 1_N) [1 - \eta_t(x)] + \eta_t(x + 1_N) [1 - \eta_t(x)] \\ &\quad - \eta_t(x) [1 - \eta_t(x + 1_N)] - \eta_t(x) [1 - \eta_t(x - 1_N)] \end{aligned}$$

$$\frac{d}{dt} \frac{1}{N} \sum_{X \in \mathbb{T}_N} G(x) \eta_t(x) = \frac{1}{N^2} \frac{1}{N} \sum_{X \in \mathbb{T}_N} (\Delta_N G)(x) \eta_t(x)$$

$$(\Delta_N G)(x) = N^2 [G(x + 1_N) + G(x - 1_N) - 2G(x)]$$

# Glauber hydrodynamics

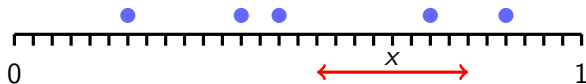


$$\frac{d}{dt} \eta_t(x) = c(\tau_x \eta_t) [1 - \eta_t(x)] - c(\tau_x \eta_t) \eta_t(x)$$

$$\frac{d}{dt} \frac{1}{N} \sum_{X \in \mathbb{T}_N} G(x) \eta_t(x) = \frac{1}{N} \sum_{X \in \mathbb{T}_N} G(x) c(\tau_x \eta_t) [1 - 2\eta_t(x)]$$

$$\langle \rho_N(t), G \rangle = \frac{1}{N} \sum_{X \in \mathbb{T}_N} G(x) \eta_t(x)$$

$$\frac{d}{dt} \langle \rho_N(t), G \rangle = \frac{1}{N} \sum_{x \in \mathbb{T}_N} \left\{ (\Delta_N G)(x) \eta_t(x) + G(x) c(\tau_x \eta_t) [1 - 2\eta_t(x)] \right\}$$



- Exclusion speeded-up  $N^2$   $\mu_N(t, x) \sim \nu_\alpha$   $\alpha \sim \rho_N(t, x)$
- $F(\alpha) = E_{\nu_\alpha} [c(\eta) [1 - 2\eta(0)]]$
- $c(\tau_x \eta_t) [1 - 2\eta_t(x)] \sim F(\rho_N(t, x))$

$$\partial_t \rho = \Delta \rho + F(\rho)$$



# Reaction-diffusion equations

- $F(\alpha) = -V'(\alpha)$

$$\begin{cases} \partial_t u = \Delta u - V'(u) \\ u(0, \cdot) = u_0(\cdot) \end{cases}$$

- $F(\alpha) = E_{\nu_\alpha} [c(\eta) [1 - 2\eta(0)]]$  polynomial
- Exists unique solution and  $u \in C^\infty((0, \infty) \times \mathbb{T})$
- $u \leq v \implies u(t) \leq v(t)$
- $V'(0) = -F(0) = -c(\mathbf{0}) < 0$      $V'(1) = -F(1) = c(\mathbf{1}) > 0$

$$0 = \Delta \rho - V'(\rho)$$

- $\mathcal{S}$  solutions

**Theorem:** [Chen, Matano \(89\)](#)  $\forall u_0, \exists \rho \in \mathcal{S}$  s.t.  $u(t) \longrightarrow \rho$ .

# Stationary state

- $\mathfrak{N}_\epsilon(\mathcal{S}) \quad \mu_N[\mathfrak{N}_\epsilon(\mathcal{S})^c] \rightarrow 0$

$$\mu_N[\rho_N \notin \mathfrak{N}_\epsilon(\mathcal{S})] = \mathbf{P}_{\mu_N}[\rho_N(0) \notin \mathfrak{N}_\epsilon(\mathcal{S})] = \mathbf{P}_{\mu_N}[\rho_N(t) \notin \mathfrak{N}_\epsilon(\mathcal{S})]$$

- $\rho_N(t) \rightarrow u(t)$

$$\begin{cases} \partial_t u = \Delta u - V'(u) \\ u(0, \cdot) = u_0(\cdot) \end{cases}$$

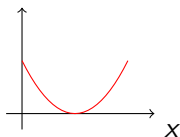
- $u(t) \rightarrow \gamma \in \mathcal{S}$

- $u(t) \in \mathfrak{N}_\epsilon(\mathcal{S}) \quad t \text{ large}$

# Solutions of the elliptic equation

$$0 = \Delta u - V'(u)$$

- $V'(\alpha) = 0$   $\gamma(x) = \alpha$  solution.
- $V$  polynomial degree 2.
- $V'(0) < 0$   $V'(1) > 0$
- $u_0(x) = a$   $u(t, x) = a(t)$   $\dot{a}(t) = -V'(a(t))$
- $u^{(a)}(x) = a$   $u^{(0)} \leq u \leq u^{(1)}$   $u^{(0)}(t) \leq u(t) \leq u^{(1)}(t)$
- $u(t) \rightarrow \mathfrak{m}$   $\mathfrak{m}$  minima of  $V$ .



# Solutions of the elliptic equation

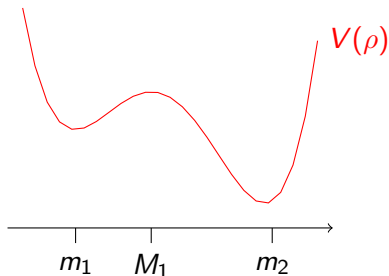
$$0 < m_1 < M_1 < m_2 < \cdots < m_{n-1} < M_{n-1} < m_n$$

- $\Delta \rho - V'(\rho) = 0$
- $\partial_t^2 \rho - V'(\rho) = 0$

Hamiltonian dynamics

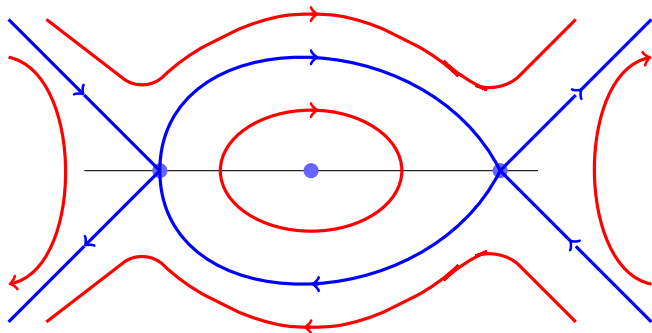
- $(q(t), p(t)) \quad q(t) = \rho(t)$

$$\begin{cases} \dot{q}(t) = p(t) \\ \dot{p}(t) = -V'(q(t)) \end{cases}$$



# Hamiltonian dynamics

- $\dot{q}(t) = p(t) \quad \dot{p}(t) = -V'(q(t))$



- Periodic solutions

$$\begin{cases} \dot{q}(t) = p(t) \\ \dot{p}(t) = -V'(q(t)) \end{cases} \quad E(p, q) = \frac{1}{2}p^2 - V(q)$$

$$\dot{E}(t) = p(t)\dot{p}(t) - V'(q(t))\dot{q}(t) = p(t)V'(q(t)) - V'(q(t))p(t) = 0$$

$$\frac{1}{2}p(t)^2 - V(q(t)) = C_0 \quad p(t) = \pm \sqrt{2V(q(t)) + C_0}$$

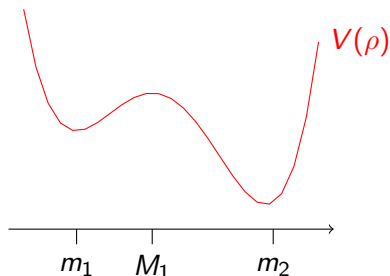
Theorem:

- Zeros of  $F$  are real numbers
- Critical points of  $V$  are maxima or minima
- Then, at most finite number of solutions [up to translations].

# Unstable solutions

$\mathcal{S}$  solutions  $\mu_N(\mathfrak{N}_\epsilon(\mathcal{S})) \rightarrow 1$

- Unstable solutions
- $\partial_t u = \Delta u - V'(u)$
- $u(0) = a \quad u(t) = a(t)$
- $\dot{a}(t) = -V'(a(t))$
- $u(0) = M_j \pm \epsilon \quad u(t) \rightarrow m_{j\pm 1}$



$\mathcal{U}$  “unstable” solutions  $\mu_N(\mathfrak{N}_\epsilon(\mathcal{U})) \rightarrow 0$

# Dynamical large deviations

Kipnis, Olla, Varadhan (89)    Donsker, Varadhan (89)

Jona-Lasinio, L, Vares (93)    Bodineau, Lagouge (10, 12)    L, Tsunoda (15)

- initial state  $\eta^N$      $\rho_N(\eta^N) \rightarrow \gamma$

Hydrodynamical limit:

$$\mathbb{P}_{\eta^N}[\rho_N(t) \sim u(t), 0 \leq t \leq T] \rightarrow 1 \quad \begin{cases} \partial_t u = \Delta u + F(u) \\ u(0, \cdot) = \gamma(\cdot) \end{cases}$$

Large deviations:

- $u(t) = u(t, x) \quad 0 \leq t \leq T$

$$\mathbb{P}_{\eta^N}[\rho_N(t) \sim u(t), 0 \leq t \leq T] \approx e^{-N I_{[0, T]}(u)}$$



# Large Deviations rate function

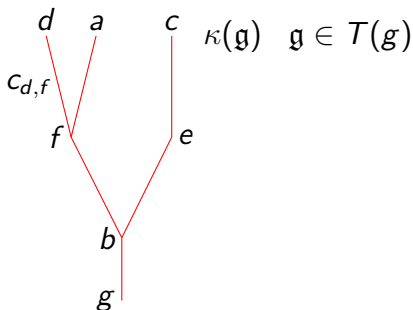
- $B(\alpha) = E_{\nu_\alpha} [\{1 - \eta(0)\} c(\eta)]$      $D(\alpha) = E_{\nu_\alpha} [\eta(0) c(\eta)]$
- $F(\alpha) = B(\alpha) - D(\alpha)$
- $\sigma(\alpha) = \alpha(1 - \alpha)$
- $B, D$  concave functions

$$\partial_t u = \Delta u + F(u) - \partial_x [\sigma(u) \partial_x H] + B(u)(e^H - 1) - D(u)(e^{-H} - 1)$$

$$I_T(u) = \frac{1}{2} \int_0^T dt \int_{\mathbb{T}} dx \sigma(u) (\nabla H)^2 + \int_0^T dt \int_{\mathbb{T}} dx \left\{ B(u) (1 - e^H + He^H) + D(u) (1 - e^{-H} - He^{-H}) \right\}$$

# Quasi-potential

- $\bar{\rho}_1, \dots, \bar{\rho}_p$  stationary solutions  $\mathcal{M}_1, \dots, \mathcal{M}_p$
- $V_i(\gamma) = \inf_{T>0} \inf \{I_{[0,T]}(u) : u_0 \in \mathcal{M}_i, u(T) = \gamma\}$
- $c_{i,j} = V_i(\bar{\rho}_j)$
- $T(i)$  trees  $i$  root
- $\mathfrak{g} \in T(i) \quad \kappa(\mathfrak{g}) = \sum_{(a,b) \in \mathfrak{g}} c_{a,b}$
- $w_i = \min_{\mathfrak{g} \in T(i)} \kappa(\mathfrak{g})$
- $v_i = w_i - \min_k w_k$
- $\mu_N(\mathcal{N}_\epsilon(\mathcal{M}_i)) \approx e^{-Nv_i}$



$$W(\gamma) = \min_i \{v_i + V_i(\gamma)\}$$

# Static large deviations (Farfán, L, Tsunoda)

- $B(\rho) = E_{\nu_\rho} [[1 - \eta(0)] c(\eta)]$
- $D(\rho) = E_{\nu_\rho} [\eta(0) c(\eta)]$
- $-V'(\rho) = F(\rho) = B(\rho) - D(\rho)$
- $\Delta\rho = V'(\rho)$  has finite number of solutions
- $B$   $D$  concave functions

- $\mathcal{C}$  closed     $\mathcal{O}$  open

$$\limsup_{N \rightarrow \infty} \frac{1}{N} \log \mu_N(\rho_N \in \mathcal{C}) \leq - \inf_{\vartheta \in \mathcal{C}} W(\vartheta),$$

$$\liminf_{N \rightarrow \infty} \frac{1}{N} \log \mu_N(\rho_N \in \mathcal{O}) \geq - \inf_{\vartheta \in \mathcal{O}} W(\vartheta).$$

$W$  quasi-potential

# Static large deviations (Farfán, L, Tsunoda)

- $\Delta\rho = V'(\rho)$  has finite number of solutions
- $B$   $D$  concave functions

$$\mu_N(\rho_N \sim \gamma) \approx e^{-NW(\gamma)}$$

- $W$  bounded, lower semicontinuous, compact level sets
- $\mu_N(\mathcal{N}_\epsilon(\mathcal{M}_i)) \approx e^{-Nv_i}$
- $\mathcal{S}_{\text{unst}} = \{\bar{\rho}_j : v_j > 0\}$
- $\mathcal{S}_0 = \mathcal{S} \setminus \mathcal{S}_{\text{unst}}$

$$\mu_N(\mathfrak{N}_\epsilon(\mathcal{S}_0)) \rightarrow 1$$

- Difficult to characterize unstable solutions
- Analytical problem    optimal formulation

$$\mu_N(\mathcal{N}_\epsilon(\mathcal{M}_i)) \approx e^{-Nv_i}$$

- $\mathcal{M}_i$  is **stable** if  $\forall \epsilon > 0 \inf_{\gamma \notin \mathcal{N}_\epsilon(\mathcal{M}_i)} V_i(\gamma) > 0$
- If  $\mathcal{M}_i$  stabile, then  $c_{i,j} > 0 \forall j \neq i$
- $\mathcal{M}_i$  stabile does not imply that  $v_i = 0$
- All minima  $m_k$  are stabile

Unstable condition:

- Let  $j$  such that  $c_{j,k} = 0$  for some  $k$
- Suppose that  $c_{k,\ell} > 0$  for all  $\ell \neq k$ .
- Then,  $v_j > 0$

# Heteroclinic orbits

- $\phi \rightarrow \psi \quad \lim_{t \rightarrow +\infty} u(t) = \psi \quad \lim_{t \rightarrow -\infty} u(t) = \phi$
- $\rho_j \rightarrow \rho_k \implies c_{j,k} = 0$
- $M_j \rightarrow m_{j+1} \quad M_j \rightarrow m_j \implies v(M_j) > 0$

Fiedler, Rocha, Wolfrum (04)

- $V''(M_j) \neq 0$
- $\phi$  non-constant solution  $m_i < \phi(x) < m_{i+1}$
- There are heteroclinic orbits  $\phi \rightarrow \bar{\rho}_i, \phi \rightarrow \bar{\rho}_{i+1}$
- $v(\phi) > 0$

# Conclusion

- Remains to remove local minima

