

Anderson Localization in the presence of topologically protected channels

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Research



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- E. Khalaf, M. A. Skvortsov, and P. M. Ostrovsky, PRB 2016
- E. Khalaf, P. M. Ostrovsky, PRL 2017
- E. Khalaf, P. M. Ostrovsky, PRB 2017
- E. Khalaf, M. A. Skvortsov, and P. M. Ostrovsky, in preparation.
- E. Khalaf, PhD Thesis

Anderson Localization

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- Absence of diffusion in disorder one or two dimensional systems for any disorder strength and in three dimensional systems for sufficiently strong disorder (Anderson 58, Abraham *et al.* 80)

Anderson Localization

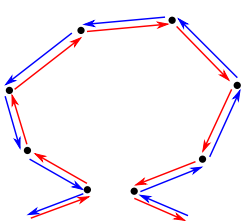
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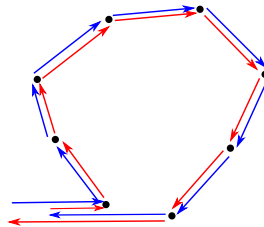
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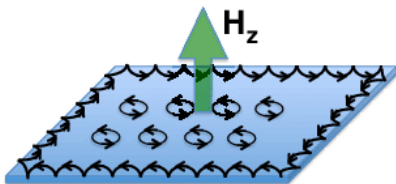


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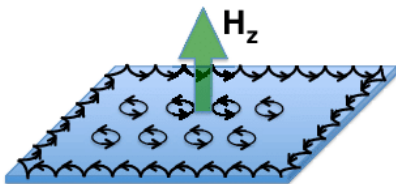


P' is the time-reverse of P

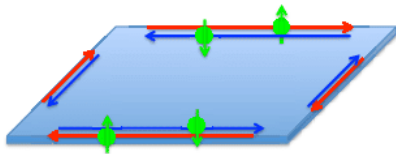
Topological protection



Quantum Hall effect
von Klitzing 81



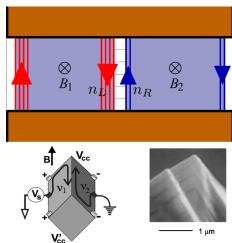
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Quantum Spin Hall effect
Kane Mele 05, Bernevig *et al.* 06, K'önig *et al* 07

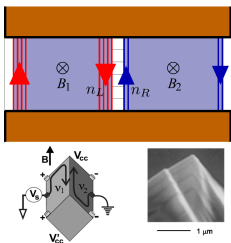
Localization + protected channels

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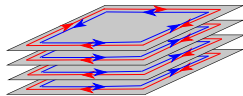
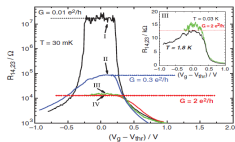


**Interface between two
quantum Hall systems**
(Grayson *et. al.* 07,08)

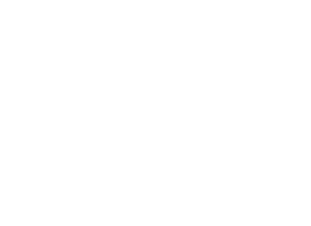
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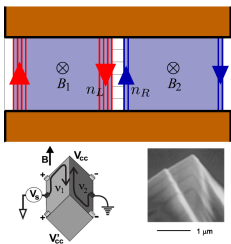
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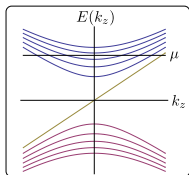
Stack of quantum spin-Hall systems (Koenig *et. al.* 07)



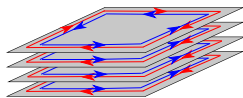
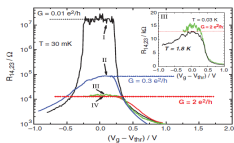
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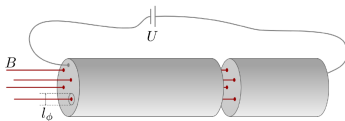
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Weyl semimetal in a magnetic field (Wan *et. al.* '11, Zyuzin & Burkov '12, Altland & Bagrets '15, Huang *et. al.* '15, Shekhar *et. al.* '15, ...)



Stack of quantum spin-Hall systems (Koenig *et. al.* 07)



Diffusion + Drift

Classical Diffusion

$$W(x, t) = \frac{e^{-\frac{x^2}{2Dt}}}{2\sqrt{\pi Dt}}$$

Chiral
channels \rightarrow

Classical Diffusion+Drift

$$W(x, t) = \frac{e^{-\frac{(x-mvt/N)^2}{2Dt}}}{2\sqrt{\pi Dt}}$$

Quantum
interference
 \downarrow

Anderson localization

$$W(x, t) = ?$$

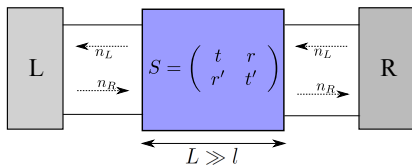
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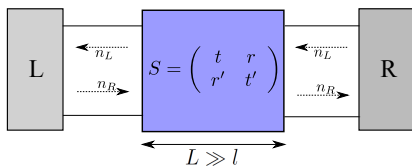
?

Transport in quasi-one-dimensional conductors

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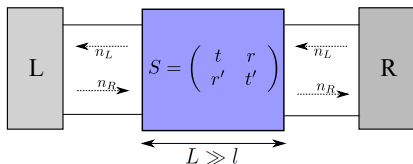


Transport in quasi-one-dimensional conductors



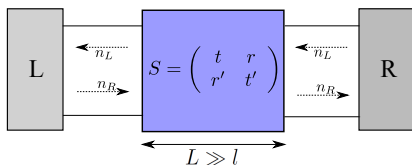
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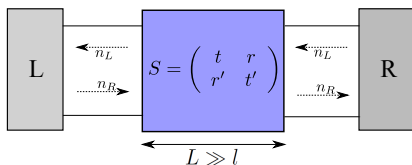
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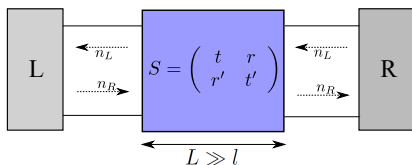
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$$\rho(T) = \langle \text{tr } \delta(t^\dagger t - T) \rangle, \quad 0 \leq T \leq 1$$

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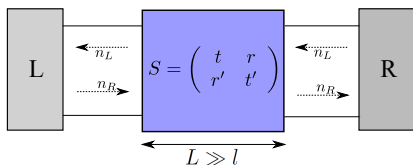
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- Lyapunov exponent $\lambda \geq 0$: $T = 1/\cosh^2 \lambda$
 - $\lambda = 0$: perfect transmission ($T = 1$)
 - $\lambda = \infty$: zero transmission ($T = 0$)

Effective field theory of disorder

Efetov '83

Effective field theory of disorder

- Moments in terms of Green's functions: Kubo formula

$$\text{tr}(t^\dagger t)^n = \text{tr}[\hat{v}(x_L)G_\epsilon^R(x_L, x_R)\hat{v}(x_R)G_\epsilon^A(x_R, x_L)]^n$$

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$$\langle e^{i\psi_i V_{ij} \psi_j} \rangle = \int dV e^{-\frac{1}{2N\tau} V_{ij} V_{ji} + i\psi_i V_{ij} \psi_j} = e^{-2N\tau (\psi_i^* \psi_j)^2}$$

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The tenfold way

Symmetry class Cartan label	Symmetry			NLSM compact sector	Spatial dimension	
	\mathcal{T}	\mathcal{C}	\mathcal{S}		1	2
A	0	0	0	$U(2n)/U(n) \times U(n)$	0	\mathbb{Z}
AIII	0	0	1	$U(n) \times U(n)/U(n)$	\mathbb{Z}	0
AI	1	0	1	$Sp(4n)/Sp(2n) \times Sp(2n)$	0	0
BDI	1	1	1	$U(2n)/Sp(2n)$	\mathbb{Z}	0
D	0	1	0	$O(2n)/U(n)$	\mathbb{Z}_2	\mathbb{Z}
DIII	-1	1	1	$O(n) \times O(n)/O(n)$	\mathbb{Z}_2	\mathbb{Z}_2
AII	-1	0	0	$O(2n)/O(n) \times O(n)$	0	\mathbb{Z}_2
CII	-1	-1	1	$U(n)/O(n)$	\mathbb{Z}	0
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Altland, Zirnbauer 97, Kitaev 09, Schnyder *et al.* 09

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- A, C, D: arbitrary number of protected (chiral) channels (\mathbb{Z})

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- A, C, D: arbitrary number of protected (chiral) channels (\mathbb{Z})
- AII, DIII: 0 or 1 protected (helical) channels (\mathbb{Z}_2).

The tenfold way

Symmetry class Cartan label	Symmetry			NLSM compact sector	Spatial dimension	
	\mathcal{T}	\mathcal{C}	\mathcal{S}		1	2
A	0	0	0	$U(2n)/U(n) \times U(n)$	0	\mathbb{Z}
AIII	0	0	1	$U(n) \times U(n)/U(n)$	\mathbb{Z}	0
AI	1	0	1	$Sp(4n)/Sp(2n) \times Sp(2n)$	0	0
BDI	1	1	1	$U(2n)/Sp(2n)$	\mathbb{Z}	0
D	0	1	0	$O(2n)/U(n)$	\mathbb{Z}_2	\mathbb{Z}
DIII	-1	1	1	$O(n) \times O(n)/O(n)$	\mathbb{Z}_2	\mathbb{Z}_2
AII	-1	0	0	$O(2n)/O(n) \times O(n)$	0	\mathbb{Z}_2
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Altland, Zirnbauer 97, Kitaev 09, Schnyder *et al.* 09

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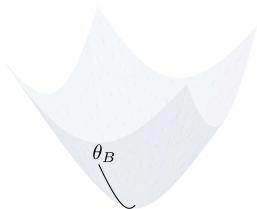
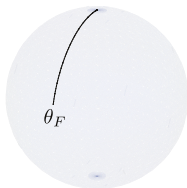
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Altland, Zirnbauer 97, Kitaev 09, Schnyder *et al.* 09

Unitary class

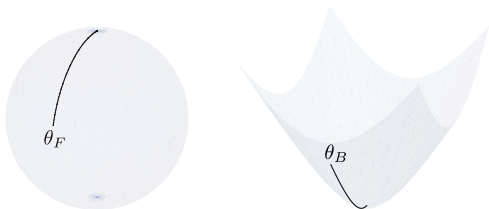
Unitary class

- Sigma model manifold: compact sector (Sphere) + non-compact sector (Hyperboloid)



Unitary class

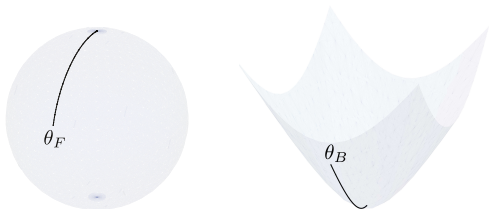
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- Boundary conditions at $x = 0$: $\theta_F = \theta_B = 0$

Unitary class

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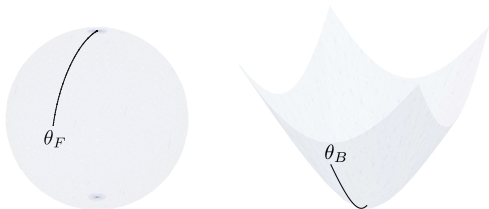


- Boundary conditions at $x = 0$: $\theta_F = \theta_B = 0$
- Transmission distribution functions

$$\rho(\lambda) = \frac{2}{\pi} \operatorname{Re} \left. \frac{\partial Z[\theta_F, \theta_B]}{\partial \theta_F} \right|_{\theta_F = i\theta_B = \pi - \epsilon - 2i\lambda}$$

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- Topological term: field of a magnetic monopole

$$S_{\text{top}} = \frac{im}{2} \int dx (1 - \cos \theta) \dot{\phi}, \quad m = n_L - n_R$$

Saddle-point approximation ($L \ll \xi$)

Action in the compact sector

$$S_F = \int_0^L dx \left[\frac{\xi}{4} (\dot{\theta}^2 + \sin^2 \theta \dot{\phi}^2) + i \frac{m}{2} (1 - \cos \theta) \dot{\phi} \right], \quad \xi = (n_R + n_L)l = 2\pi\nu D$$

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$$m = 0$$

$$m \neq 0$$

$\rightarrow \theta_F$ close to π is classically unreachable

Saddle-point approximation: distribution function

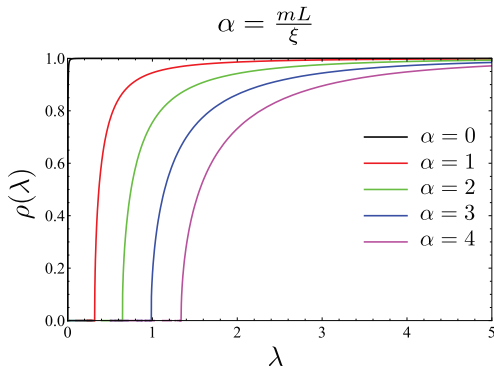
- Distribution function

$$\rho(\lambda) = \frac{2}{\pi} \operatorname{Re} \frac{\partial Z[\theta_F, \theta_B]}{\partial \theta_F} \Big|_{\theta_F = i\theta_B = \pi - \epsilon - 2i\lambda} = m\delta(\lambda) + \rho_{\text{ns}}(\lambda),$$

Saddle-point approximation: distribution function

- Distribution function

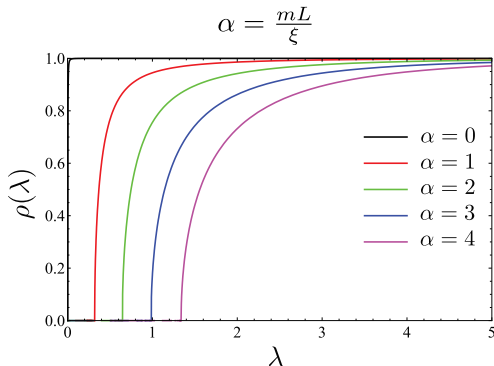
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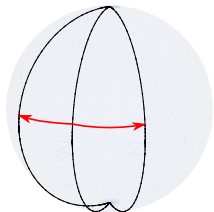
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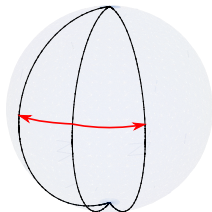
- Gap close to unit transmission \implies transport from the unprotected channels gets *suppressed*

EK, Skvortsov and Ostrovsky PRB 16

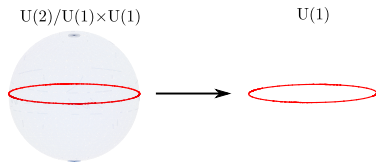


- $\theta_F \approx \pi \rightarrow$ many trajectories joining north pole to south pole \rightarrow ground state degeneracy \rightarrow some modes become very soft for $\pi - \theta \ll \sqrt{\frac{L}{\xi}}$

Soft modes



- $\theta_F \approx \pi \rightarrow$ many trajectories joining north pole to south pole \rightarrow ground state degeneracy \rightarrow some modes become very soft for $\pi - \theta \ll \sqrt{\frac{L}{\xi}}$
- Projection onto the manifold of soft modes \rightarrow trajectories labelled by polar angle ϕ "equator".



EK, Skvortsov and Ostrovsky PRB 16

H	“E”	Q_{FF}	$d=1$	$d=2$
A	AIII	$U(2n)/U(n) \times U(n)$	0	\mathbb{Z}
AIII	A	$U(n) \times U(n)/U(n)$	\mathbb{Z}	0
AI	CI	$Sp(4n)/Sp(2n) \times Sp(2n)$	0	0
BDI	AI	$U(2n)/Sp(2n)$	\mathbb{Z}	0
D	BDI	$O(2n)/U(n)$	\mathbb{Z}_2	\mathbb{Z}
DIII	D	$O(n) \times O(n)/O(n)$	\mathbb{Z}_2	\mathbb{Z}_2
AII	DIII	$O(2n)/O(n) \times O(n)$	0	\mathbb{Z}_2
CII	AII	$U(n)/O(n)$	\mathbb{Z}	0
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Mapping to 0D

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1D edge with protected channels \leftrightarrow random matrix with zero eigenvalues

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 Quantum Hall classes A, C and D \leftrightarrow chiral random matrices AIII, CII, BDI

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1D edge with protected channels	\leftrightarrow	random matrix with zero eigenvalues
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\mathbb{Z}_2 classes AII and DIII	\leftrightarrow	Random matrix classes DIII and D

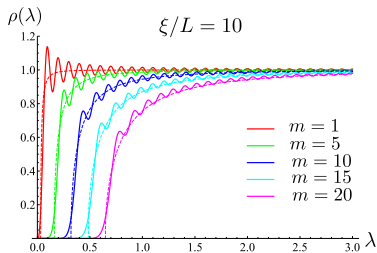
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Transmission distribution function	\leftrightarrow	Spectral density

Results

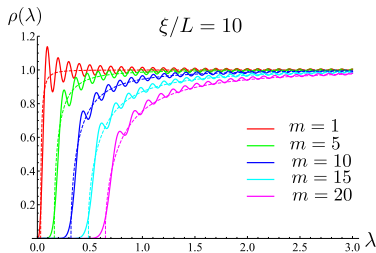
$$\rho_A(\lambda, m) = \frac{\pi \xi u}{2L} [J_m^2(u) - J_{m+1}(u)J_{m-1}(u)] + m\delta(\lambda), \quad u = \frac{\pi \xi \lambda}{L}$$



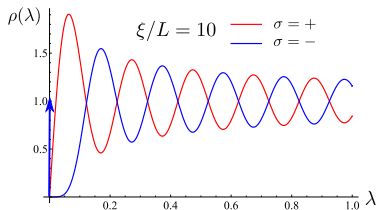
Ivanov 01, EK, Skvortsov, Ostrovsky PRB 16

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$$\rho_{\text{All}}(\lambda, \sigma = \pm) = \frac{\pi \xi u}{2L} [J_1^2(u) + J_0(u)J_1'(u)] \pm \frac{\pi \xi}{2L} J_1(u) + 2\delta_{\sigma,-}\delta(\lambda)$$



Ivanov 01, EK, Skvortsov, Ostrovsky PRB 16

Exact solution: transfer matrix method

- Transfer matrix method: 1D path integral \rightarrow Schrödinger equation

$$\partial_t \psi(Q, t) = -\mathcal{H}\psi(Q, t), \quad t = x/\xi, \quad \psi(Q, 0) = \delta(Q, \Lambda)$$

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$$S = \int dt \left[\frac{1}{4} \dot{y}^\alpha \dot{y}^\beta g_{\beta\alpha} + \dot{y}^\alpha A_\alpha \right]$$
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Solution

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- Canonical coordinates on \mathcal{G}/\mathcal{K} :
 - Spherical “Cartan” $(\mathbf{h}, K) \rightarrow$ spherical symmetry $\phi_s(Q) = \phi_s(\mathbf{h})$
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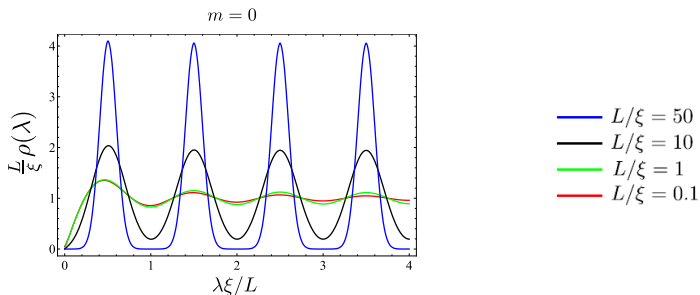
$$\phi_\nu(\mathbf{h}) = \int_{\mathcal{K}} dK e^{i\mathbf{p}_\nu \cdot \mathbf{a}(\mathbf{h}, K)}$$

- Vector potential \rightarrow Hamiltonian only simplifies in a special gauge
- Spherical eigenfunctions: gauge transformation + integration (EK, Skvortsov and Ostrovsky (to appear), EK (PhD thesis))

$$\phi_\nu(\mathbf{h}) = \int_{\mathcal{K}} dK e^{i\mathbf{p}_\nu \cdot \mathbf{a}(\mathbf{h}, K)} [\text{sdet } K_I(\mathbf{h}, K)]^m$$

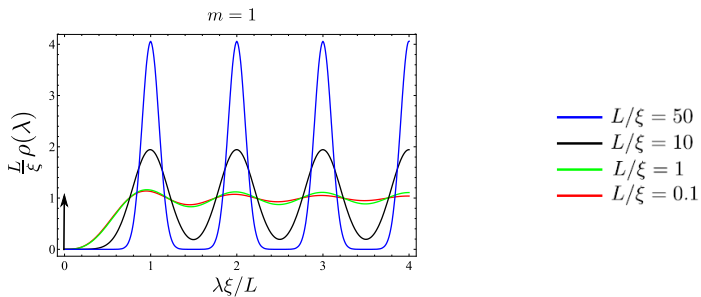
valid also without supersymmetry (compact or non-compact nism)

Effect of topology: Class A (unitary)



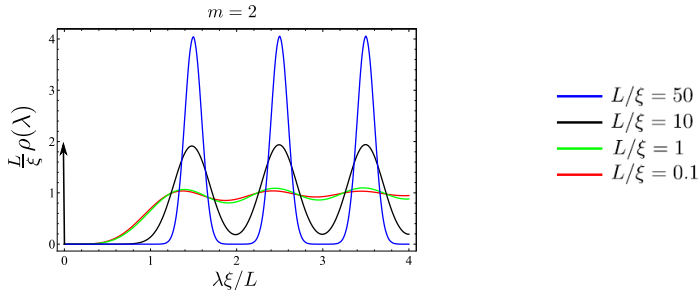
EK and Ostrovsky PRL 2017

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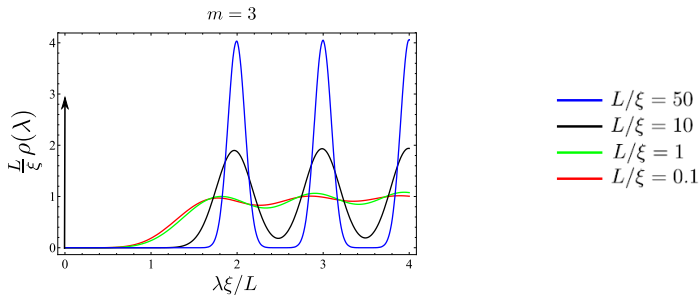
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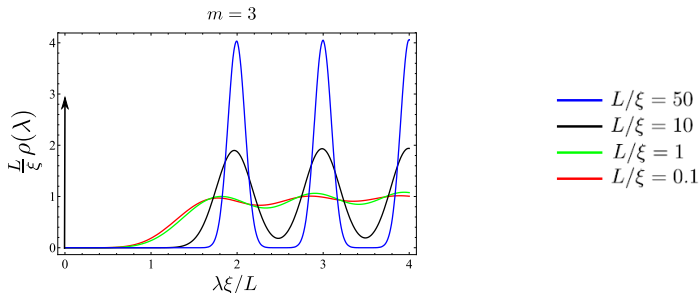
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$$G = \int_0^\infty d\lambda \frac{\rho(\lambda)}{\cosh^2 \lambda} \sim \int_0^\infty d\lambda e^{-2\lambda} \rho(\lambda)$$

- Localization length for the unprotected channels $\xi_m = \xi / (m + 1)$

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Diffusion + Drift

Classical Diffusion

$$W(x, t) = \frac{e^{-\frac{x^2}{2Dt}}}{2\sqrt{\pi Dt}}$$

Chiral
channels \rightarrow

Classical Diffusion+Drift

$$W(x, t) = \frac{e^{-\frac{(x-mvt/N)^2}{2Dt}}}{2\sqrt{\pi Dt}}$$

Quantum
interference
 \downarrow

Anderson localization

$$W(x, t) = ?$$

Chiral
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Quantum
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 \downarrow

?

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- Sigma model at finite frequency $\omega + i0 \mapsto i\Omega$

$$S[Q] = - \int dx \operatorname{str} \left[\frac{\xi}{8} (\partial_x Q)^2 - \frac{\kappa^2}{16\xi} \Lambda Q + S_{\text{top}} \right], \quad \kappa = 2N \sqrt{\Omega \tau_e}$$

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$$W(t) = \int \frac{d\omega e^{-i\omega t}}{4\pi^2\nu} \langle G_{E+\omega}^R(x_1, x_2) G_E^A(x_2, x_1) \rangle \rightarrow \frac{\nu}{2\gamma} \int d\omega e^{-i\omega t} \langle M(Q) \rangle$$

$$M(Q) = \frac{1}{32\gamma} \operatorname{str}(k\Lambda Q)^2 - \operatorname{str}(kQ)^2.$$

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 - Eigenfunction: Coulomb spherical functions (no integral representation, no generalization to more variables)

Zero mode and local dynamical correlations

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Zero mode and local dynamical correlations

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$$|\Psi_0\rangle = \int_{\mathcal{K}} dK (\text{sdet } K_{R,A})^{\pm m} \exp \left[-\frac{\kappa}{2} \text{str } P_{\pm} (KT + T^{-1}K^{-1}) \right]$$
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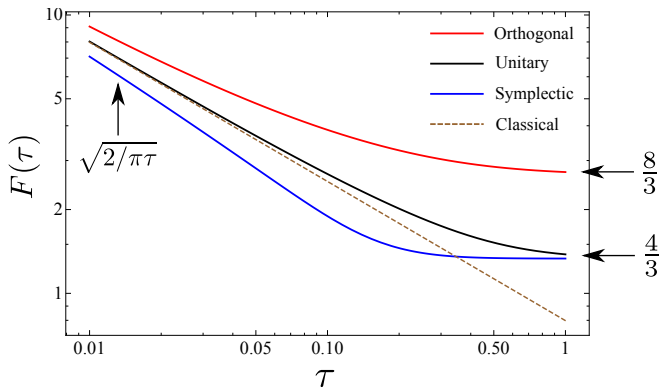
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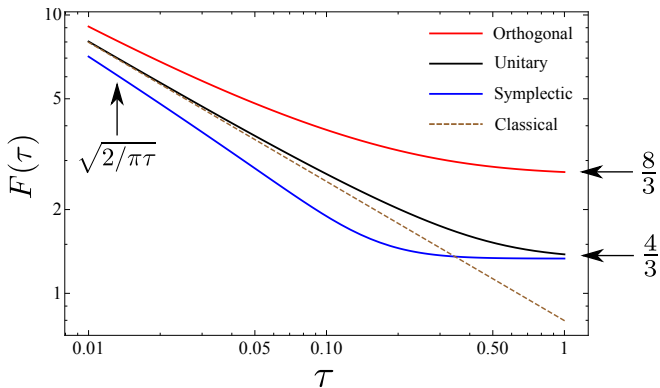
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Return probability: Different symmetry classes

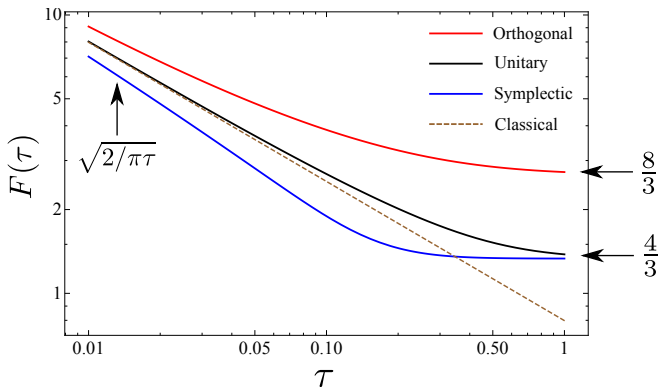


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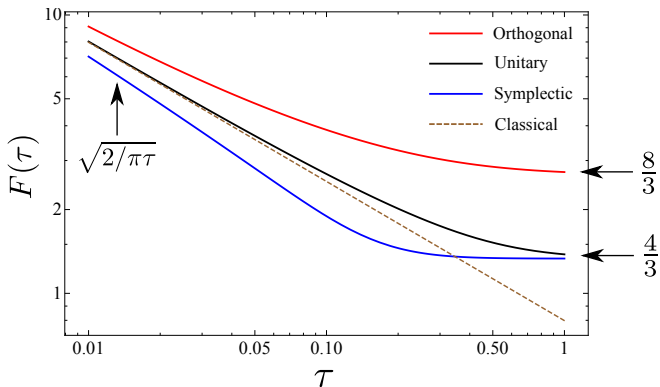
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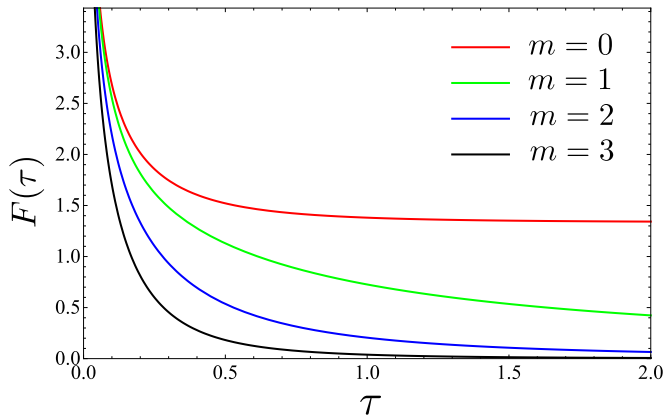
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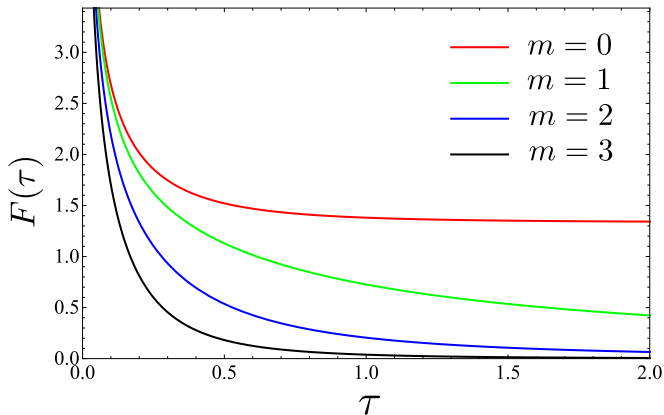
- All classes approach classical diffusion $\sqrt{\frac{2}{\pi\tau}}$ at short times.
- Unitary and symplectic classes saturate to the value $4/3$, while orthogonal class saturates to $8/3$.
- The saturation value is approached as a power-law: unitary $\sim 1/\tau^3$, orthogonal $\sim 1/\tau^2$ and symplectic $\sim 1/\tau^5$.

Return probability: chiral channels (unitary class)



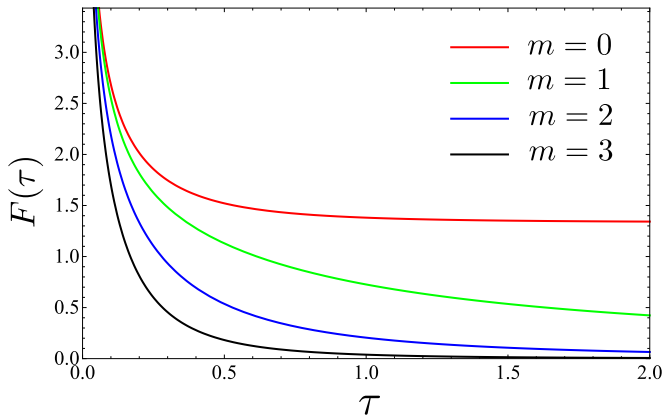
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- The wave packet leaves a tail behind due to localization.

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- Level statistics: connection to non-Hermitian systems? (Lee *et al.* PRL 2020)
- Wavefunction statistics: how does the wavefunctions look like in this system?